This lecture:

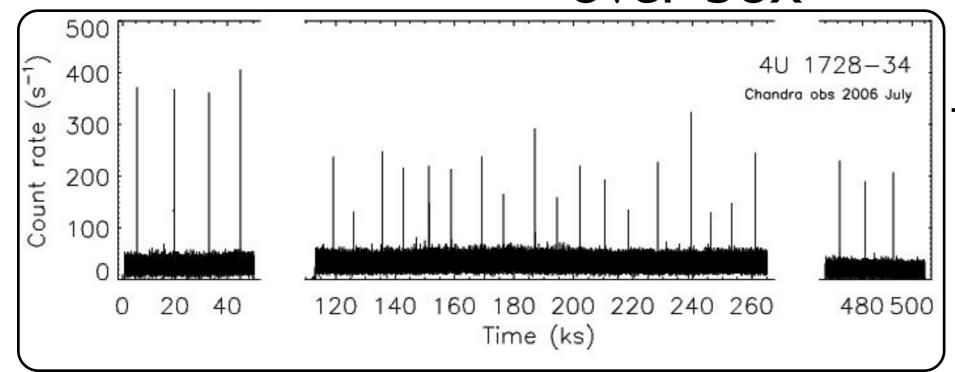
Total overview of a measurement procedure

Dealing with noise, e.g. optimal filtering

Signal-to-noise ratios

Is radiation indeed intrinsically stochastic?

One bin/measurement: average in time domain over box



Type I X-ray bursts

GALLOWAY PRIVATE COMMUNICATION

1 Measurement of x(t) in a time T => windowing and averaging over time Δ T

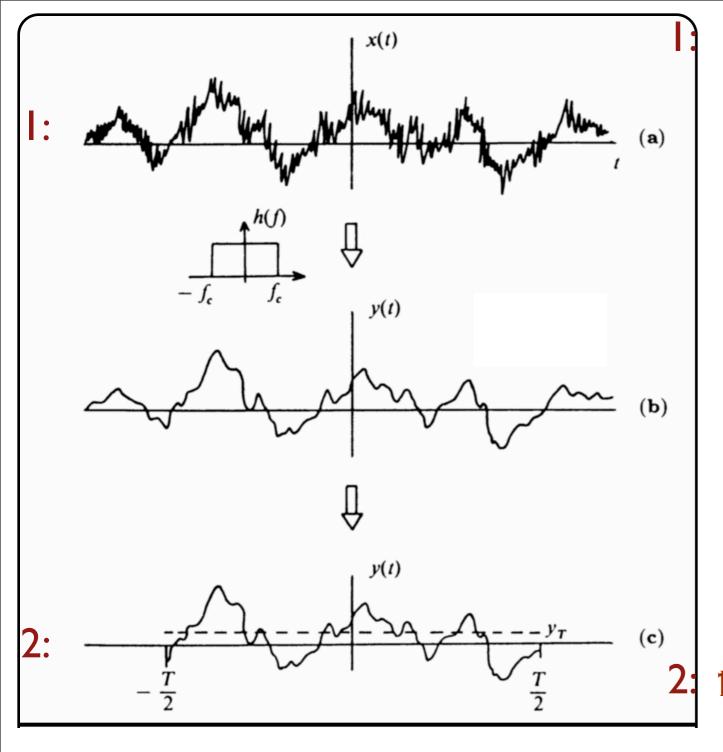
windowing:
$$y(t) = \Pi(\frac{t}{T})x(t)$$

This is a convolution!

AVERAGING:

$$z(t) \equiv y_{\Delta T}(t) = \frac{1}{\Delta T} \int_{t-\Delta T/2}^{t+\Delta T/2} y(t') dt' = \frac{1}{\Delta T} \int_{-\infty}^{\infty} \Pi(\frac{t-t'}{\Delta T}) y(t') dt'$$

= LOW-PASS FILTER, REMEMBER NYQUIST THEOREM



RESPONSE

REMOVE HIGH FREQUENCIES DUE TO FINITE RESPONSE OF INSTRUMENT AND DETECTOR

Remember convolution goes from $-\infty$ to ∞

SAMPLE -AVERAGING

Average in time domain over length T, i.e. convolve in time domain with box function

Note figure used here different from Lecture notes p25!

$$z(t) = \frac{1}{\Delta T} \prod \left(\frac{t}{\Delta T}\right) * y(t)$$

$$z(t) = \frac{1}{\Delta T} \prod \left(\frac{t}{\Delta T}\right) * \prod \left(\frac{t}{T}\right) x(t)$$

What is the combined effect of finite sample and finite bin size?

Consider example on pages 25-27 of Lecture notes

$$\sigma_{x_T}^2 = \sigma^2/N \text{ for } T >> \tau_0$$

ESTIMATING σ IF THE PROBABILITY DENSITY FUNCTION OF THE S.P. IS KNOWN



PROPAGATION OF ERRORS CHAPTER 5.2

ESTIMATING σ if the probability density function of the s.p. is not known



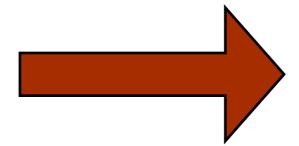
BOOTSTRAP METHOD

JACKKNIFE METHOD

MORE ON THIS LATER

SIGNAL DETECTION INVOLVES:

- LIMITED TIME INTERVAL → WINDOWING
- NOT CONTINUOUS → SAMPLING
- SAMPLES NOT INSTANTANEOUS → AVERAGING
- DEALING WITH NOISE → FILTERING
- RESPONSE OF THE DETECTION SYSTEM



THUS THE DETECTED SIGNAL WILL ONLY APPROXIMATE THE SOURCE SIGNAL

Recap

SAMPLING: HIGH FREQUENCIES ARE FILTERED OUT WINDOW: LOW FREQUENCIES ARE FILTERED OUT



LEADS TO BAND LIMITED DATA



MEASURE A PROCESS X(T) OVER INTERVAL T ASSUMED
ZERO OUTSIDE T

$$\equiv y(t) = \Pi(\frac{t}{T})x(t)$$

$$Y(f) = X(f) * Tsinc(Tf)$$

ALL INFORMATION ABOUT FREQUENCIES < 1/T IS LOST!

Recap

SAMPLING: HIGH FREQUENCIES ARE FILTERED OUT WINDOW: LOW FREQUENCIES ARE FILTERED OUT



LEADS TO BAND LIMITED DATA

→ FREQUENCY

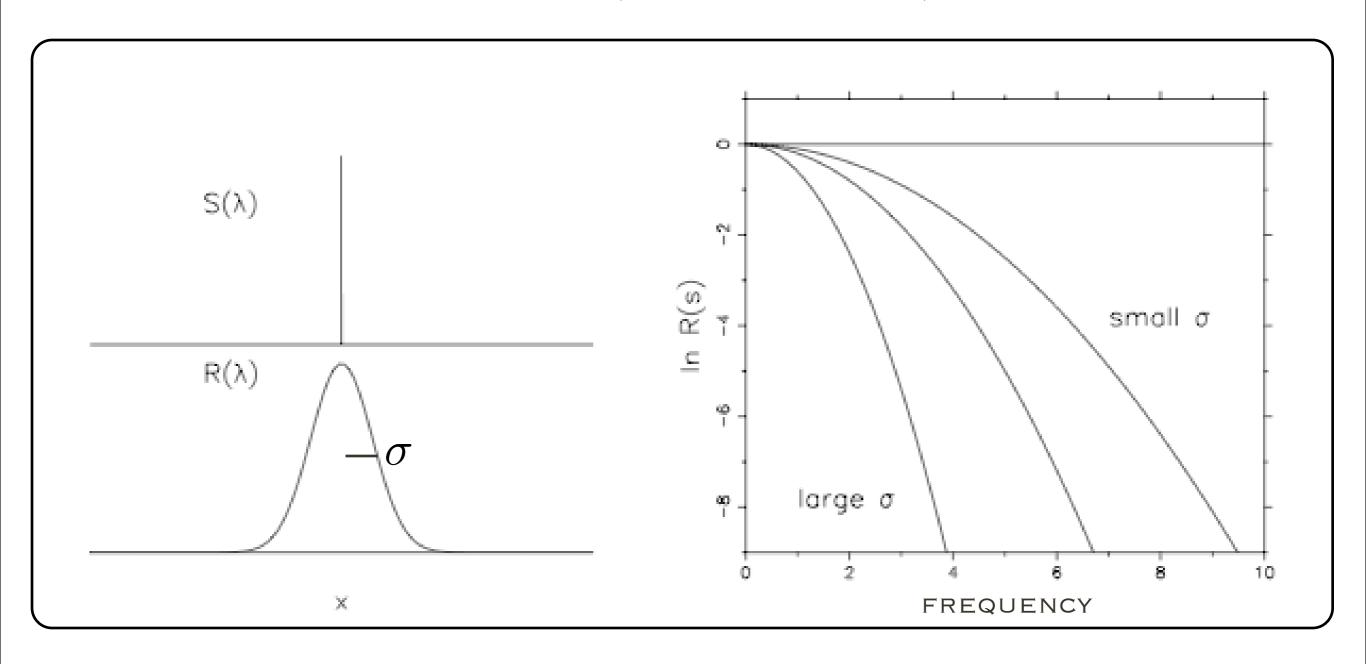
$$Y(f) = X(f)H(f)$$

$$y(t) = \int_{-\infty}^{\infty} x(t - \theta)h(\theta)d\theta$$
$$y(t) = x(t) * h(t)$$

FILTERING OF PROCESS X WITH FILTER H

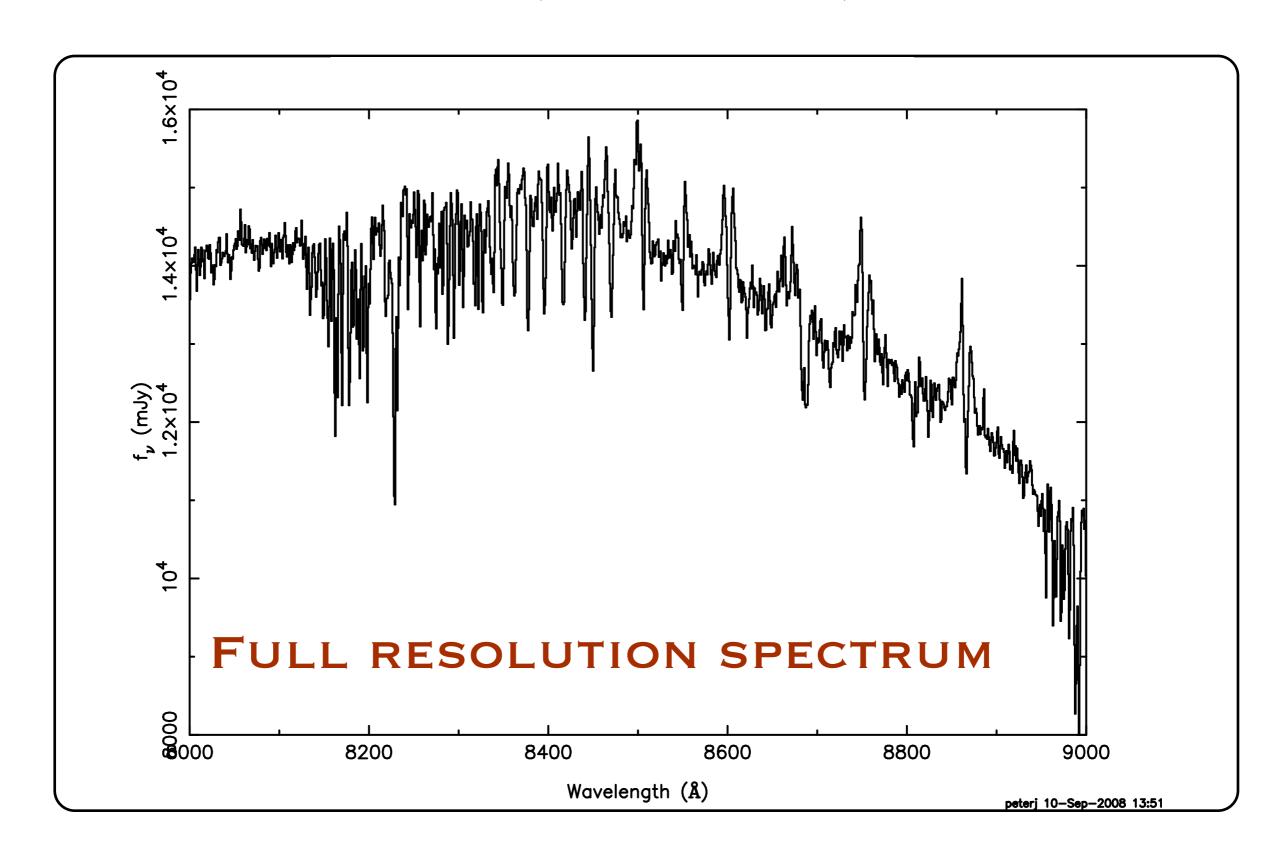
Recap: detector/instrument system response

GAUSSIAN RESPONSE FUNCTION

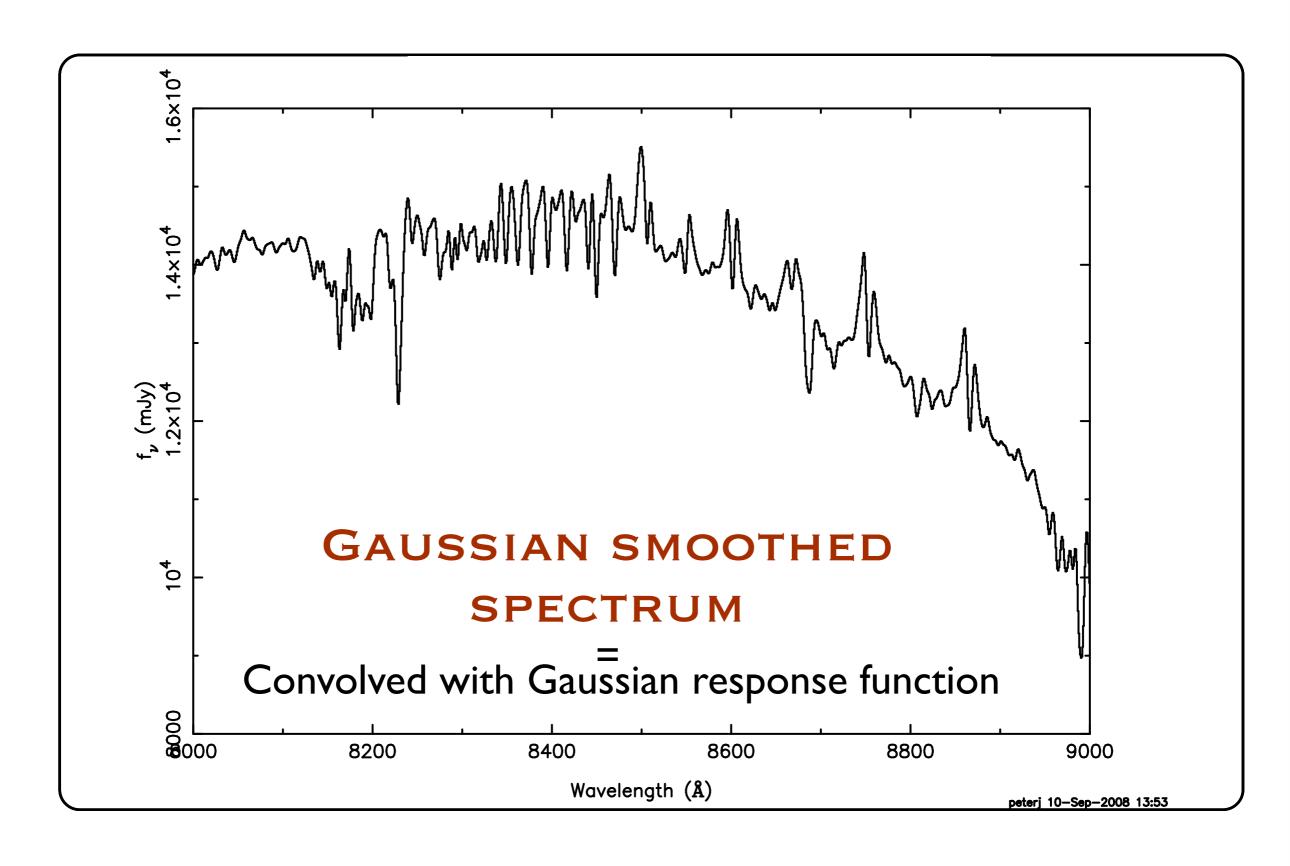


$$R(\lambda) \frac{1}{\sqrt{(2\pi)\sigma}} exp - (\frac{\lambda^2}{2\sigma^2})$$

GAUSSIAN RESPONSE FUNCTION



GAUSSIAN RESPONSE FUNCTION



Mazur: Peer Instruction

Identify in a real astrophysical measurement all the steps discussed above

Response instrument/detector
Finite duration measurement
Sampling
Average of I sample
Noise contributions?

RECAP

DATA SAMPLING: IDEAL CASE NYQUIST CRITERIUM IS FULFILLED -> SAMPLING DOES NOT LEAD TO LOSS OF INFORMATION RV completely described by the samples CONDITIONS either:

- →BAND-LIMITED RESPONSE OF THE DETECTOR

 REMOVES HIGHEST NOISE POWERS AND THE

 SAMPLING IS FAST ENOUGH TO COVER THE BAND

 LIMIT OF THE DETECTOR

 Or:
- \rightarrow SIGNAL IS BAND-LIMITED and $\nu_{\rm sampling} > \nu_{\rm max, detector} > \nu_{\rm max, signal}$

however, noise is often white and thus contains high frequencies!

Noise Removal by Optimal

FILTERING Num. Res 13.3

$$c(t) = s(t) + n(t)$$

s(t) is the smeared signal i.e. true * response s(t) = u(t) * r(t) DESIGN AN OPTIMAL FILTER $\varphi(\mathsf{T})$ THAT PRODUCES A SIGNAL $\widetilde{U}(T)$ AS CLOSE AS POSSIBLE TO U(T) or U(f)

$$\widetilde{U(f)} = \frac{C(f)\phi(f)}{R(f)}$$

CLOSE IN LEAST SQUARE SENSE

$$\int_{-\infty}^{\infty} |\widetilde{U(f)} - U(f)|^2 df \text{ is minimised}$$

NOISE REMOVAL BY OPTIMAL FILTERING

$$\int_{-\infty}^{\infty} |\frac{[S(f) + N(f)]\phi(f)}{R(f)} - \frac{S(f)}{R(f)}|^2 df = 0$$

 $\int S(f)N(f)df \ \ \text{Terms are Zero since noise}$ and signal are uncorrelated

$$\int_{-\infty}^{\infty} |R(f)|^{-2} \{ |S(f)|^2 |1 - \phi(f)|^2 + |N(f)|^2 |\phi(f)|^2 \} df = 0$$

 Θ minimised with respect to ϕ

Noise removal by optimal filtering

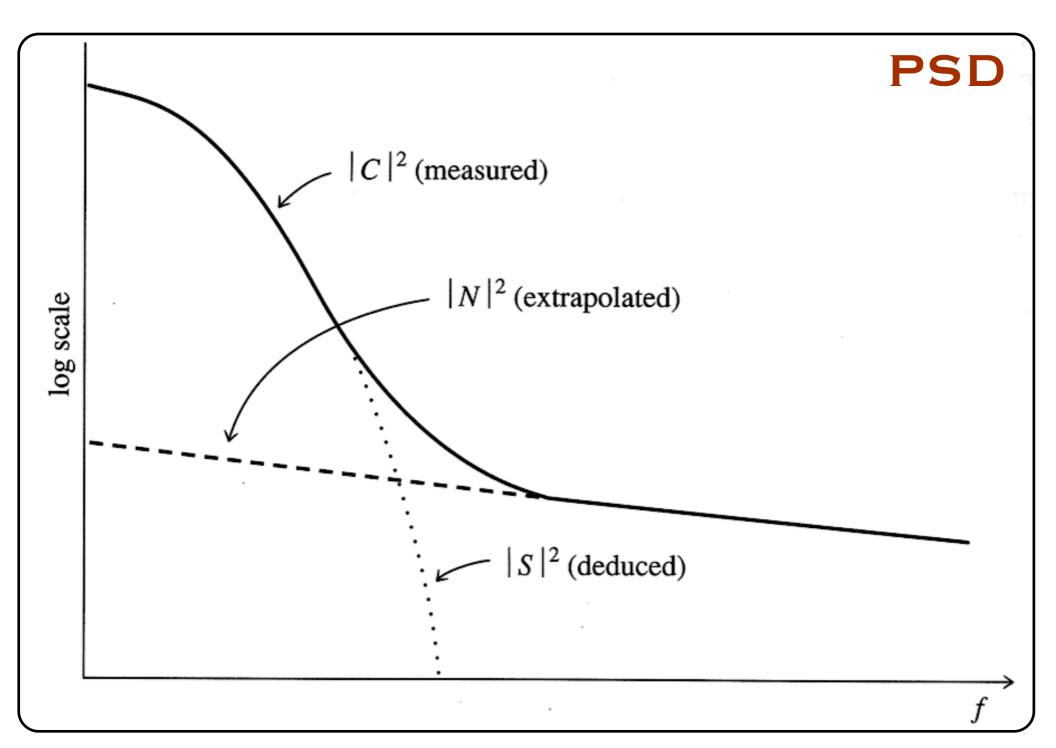
$$\frac{d\theta}{d\phi} = 0$$
$$-2S^2(1-\phi) + 2N^2\phi = 0$$

$$\phi = \frac{S^2}{S^2 + N^2}$$

Does not contain true signal directly! U(f)

$$|S(f)|^2 + |N(f)|^2 = PDS(f) = |C(f)|^2$$

NOISE REMOVAL BY OPTIMAL FILTERING



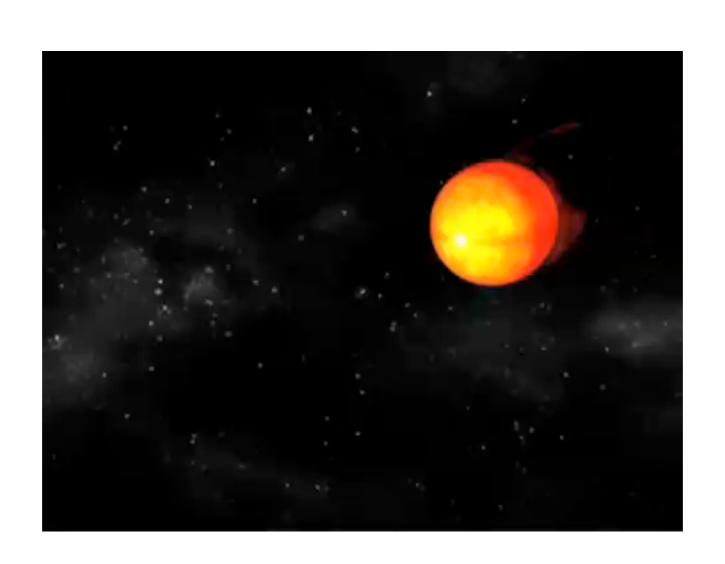
SOME APPLICATIONS OF FILTERING

ON OPTIMAL DETECTION OF POINT SOURCES IN CMB MAPS VIO ET AL, 2002, A&A, 391, 789

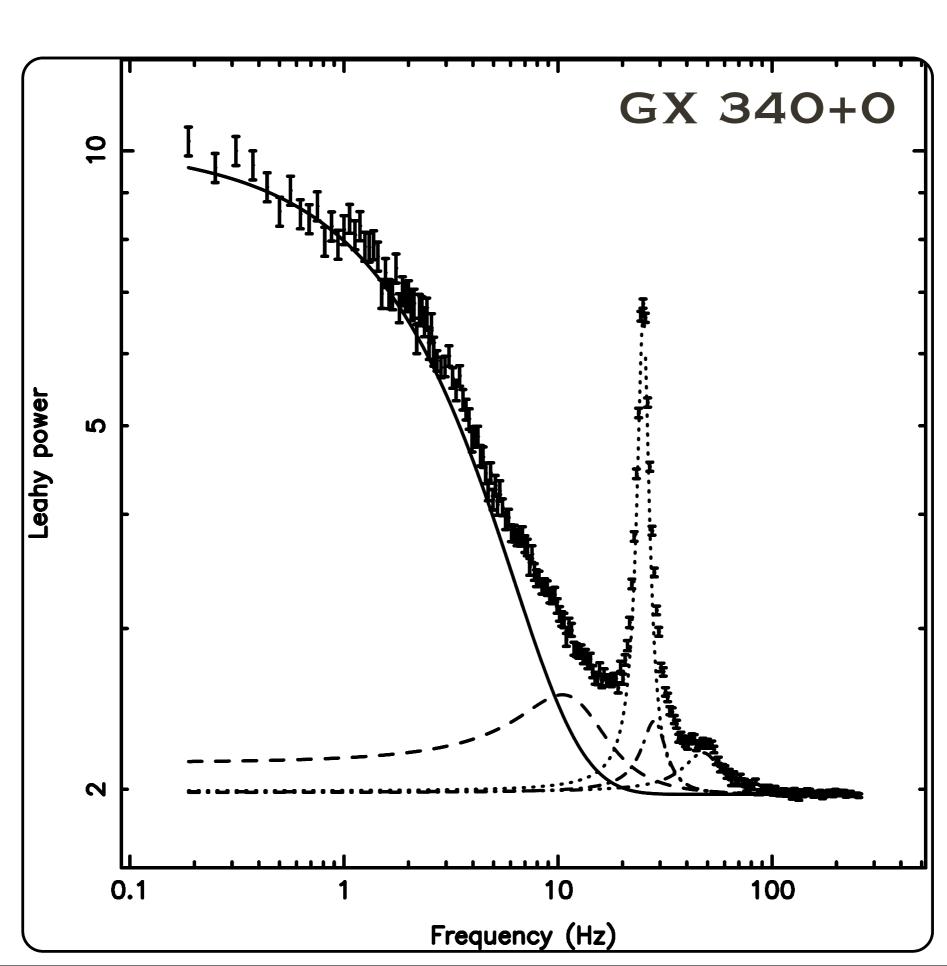
AN OPTIMAL FILTER FOR THE DETECTION OF GALAXY CLUSTERS THROUGH WEAK LENSING MATURI, ET AL. 2005, A&A, 442, 851

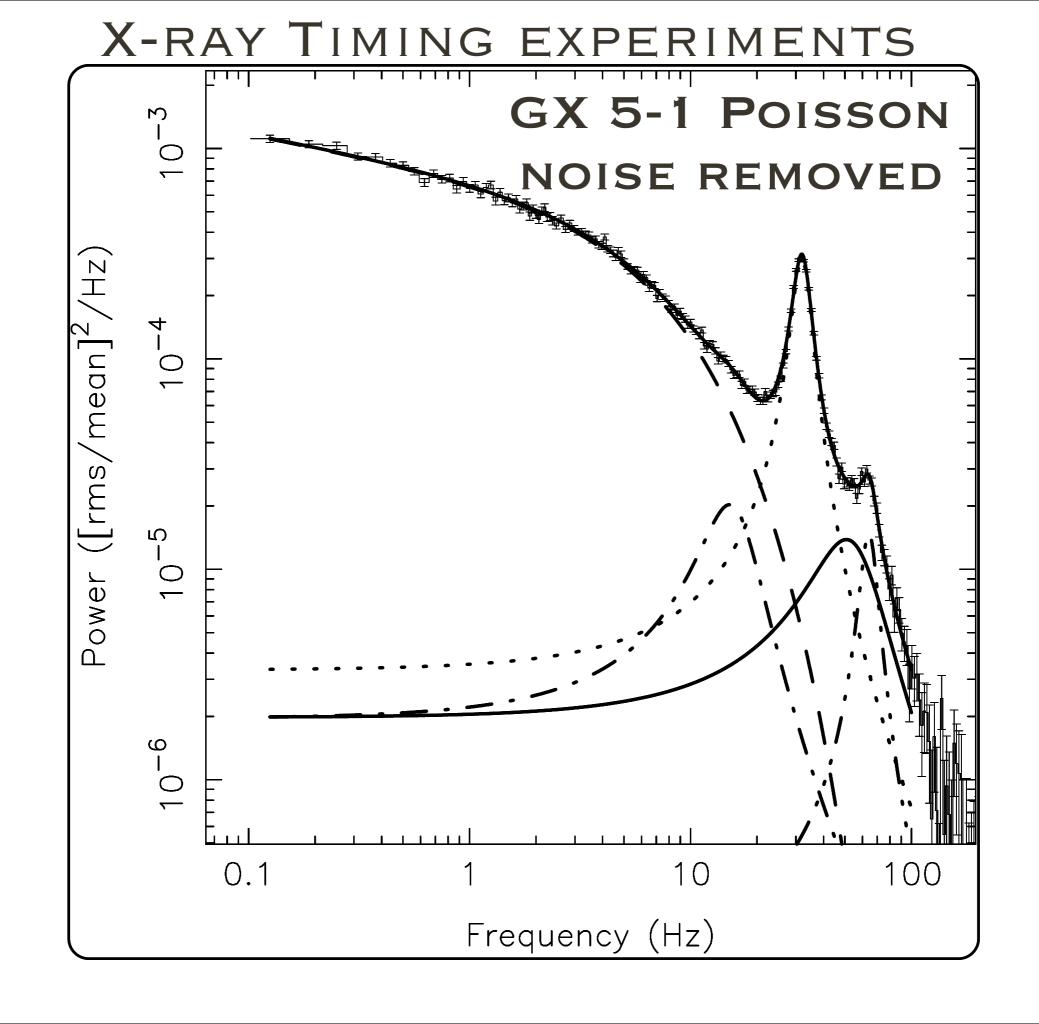
THE LARGEST SCALE PERTURBATIONS: A WINDOW ON THE PHYSICS OF THE BEGINNING
WANDELT, NEW ASTRONOMY REVIEW, 2006, 11, 900

Application to timing experiments in Low-mass X-ray binaries

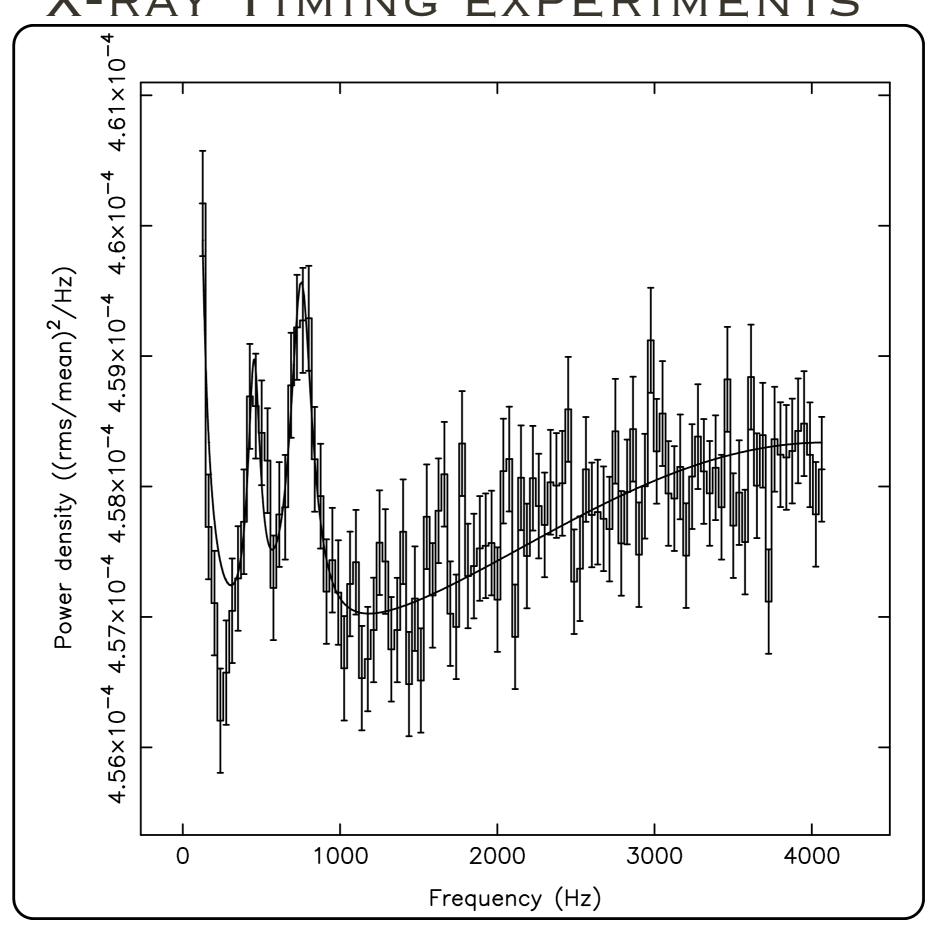


C.F. X-RAY TIMING EXPERIMENTS





X-RAY TIMING EXPERIMENTS



RECAP FILTERING:

ONE CAN DESIGN AN OPTIMAL FILTER SUCH THAT THE FILTERED MEASURED DATA-SET IS AS CLOSE AS POSSIBLE (IN LEAST-SQUARE SENSE) TO THE UNCORRUPTED SIGNAL

The Signal-to-Noise Ratio & Limiting Sensitivity

- 1. The level of noise determines the detectability of a weak (source) signal
- 2. Various noise contributions from:
 - a) other sources in FoV (point-source / diffuse) → sky-noise
 - b) operational environment of instrument (radiation belts, scatter light from Earth, Sun, Moon, atmosphere etc.)
 - c) constituting parts of instrument e.g. noise from amplifiers in electronics, induced radio-activity → instrumental noise
- 3. Quality of an observation is determined by the signal-to-noise ratio SNR

$$SNR = SNR(T_{exp}, \Delta E, A_{eff}, B)$$

4. For photon energies \geq eV (i.e. optical, UV, X-rays, γ -rays) quantum characterisation applicable: statistical treatment of the accumulated quanta (Poissonian statistics)



- \triangleright Consider point source embedded in uniform sky background noise with photon flux n_s (ph/m²s) i.e. the source-region
- Sky background noise n_{bq} (ph/m²s sr)
- Instrumental noise n_{det} (ph/m²s)

Number of counts registered in source region within observation time T_{obs}:

$$N_1 = ((n_s + n_{bg} \Delta \Omega) \bar{A}_{eff} + n_{det} A_{pix}) T_{obs}$$

 $\Delta\Omega$ - solid angle subtended by source region

 \bar{A}_{eff} - effective area (averaged over bandwidth ΔE : $\int A_{eff}(E) dE / \Delta E$)

A_{pix} - pixel surface

> Similar expression holds for a region excluding the source: the background region

$$N_2 = (n_{bg} \Delta \Omega \bar{A}_{eff} + n_{det} A_{pix}) T_{obs}$$

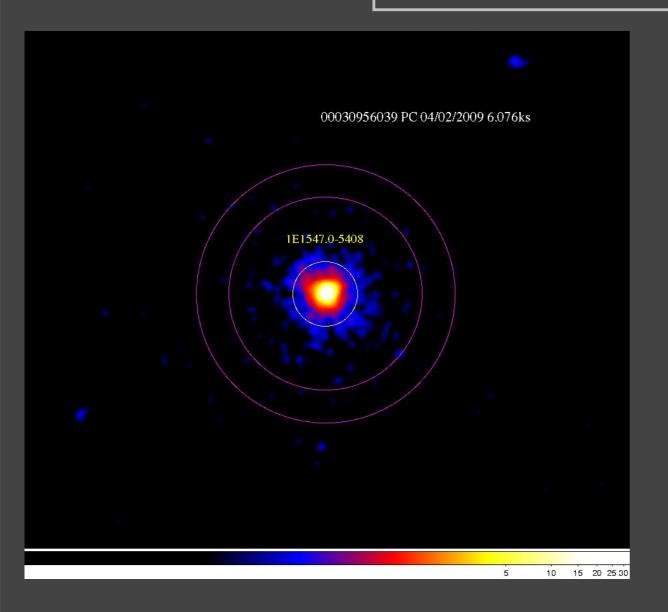
collected during same integration time simultaneously

 \triangleright Fluctuations in N1 and N2 are governed by Poissonian stat. : $\sqrt{N_1}$, $\sqrt{N_2}$



Example:

Swift-XRT observation of anomalous X-ray pulsar 1E1547.0-5408 during monitoring campaign after its Jan. 2009 outburst (in PC-mode)



Inner (yellow) ring: 60" source region

Purple annulus : 180"-240" bakground region

Note: Background region 7x larger than source region

SNR can be defined by:

$$SNR = (N_1 - N_2) / \sqrt{(N_1 + N_2)}$$

Signal strength is expressed in terms of the <u>fluctuation</u> of the noise term(s)!

We speak of a 3σ measurement, if SNR= k_{min} =3

(k_{min} is minimum confidence set by observer)

 5σ measurement, if SNR= k_{min} =5



Limiting sensitivity

We consider now 2 limiting cases:

Source signal dominates the noise: $N_1 \gg N_2$

$$SNR \sim \sqrt{N_1} = \sqrt{(n_s \bar{A}_{eff} T_{obs})}$$

demanding a minimum SNR of k_{min} (typically 3 / 5 to speak of a detection) \rightarrow

$$n_{smin} = k^2_{min} / (\bar{A}_{eff} T_{obs})$$
 \Rightarrow the detection is signal-photon-noise limited

Source signal drowned in the noise: $N_1 \approx N_2$

SNR =
$$k_{min} \sim (n_s \bar{A}_{eff} T_{obs}) / \sqrt{(2N_2)} \Rightarrow$$

$$n_{smin} = k_{min} \sqrt{(2N_2)/(\bar{A}_{eff} T_{obs})} \sim k_{min} \sqrt{(2\bar{A}_{eff} T_{obs} [n_{bg} \Delta \Omega + n_{det} \delta])/\bar{A}_{eff} T_{obs}} \propto (\bar{A}_{eff} T_{obs})^{-\frac{1}{2}}$$

$$(\delta = A_{pix}/\bar{A}_{eff})$$

n_{smin} improves only with the square root of the observation time and effective area...

If $n_{\text{det}} \delta \ll n_{\text{bg}} \Delta \Omega$ (i.e. detector noise negligible) \Rightarrow the detection is <u>background-photon-noise</u> limited



In summary:

- a) <u>signal-photon-noise</u> limited observation: $n_{smin} \sim (\bar{A}_{eff} T_{obs})^{-1}$
- b) <u>background-photon-noise</u> limited obs. : $n_{smin} \sim (\bar{A}_{eff} T_{obs})^{-1/2}$

Note: bandwidth ΔE is implicitely incorporated $n_s = \int n_s(E) dE = \check{n}_s \Delta E$

Similar expressions for n_{bg} and n_{det} :

<u>signal-photon-noise</u> limited: $\check{n}_{smin} \sim (\bar{A}_{eff} T_{obs} \Delta E)^{-1}$

<u>background-photon-noise</u> limited: $\check{n}_{smin} \sim (\bar{A}_{eff} T_{obs} \Delta E)^{-1/2}$

