

Imagers and Detectors

ATI 2015 Lecture 09
M. Kenworthy // Leiden Observatory

Observations

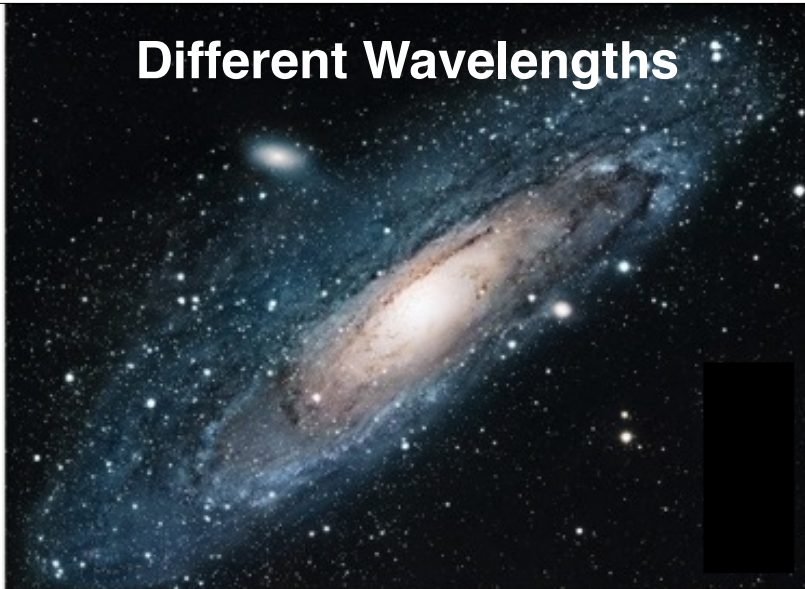
Astronomical observations are:

Expensive

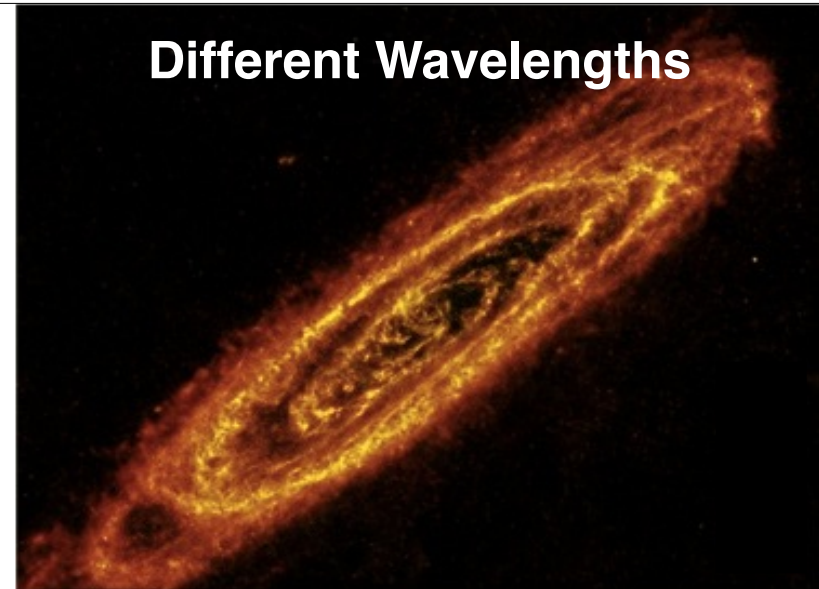
Impossible to repeat in a controlled way

An **OBSERVATION** is a permanent record of what is seen
at the focal plane of a telescope.

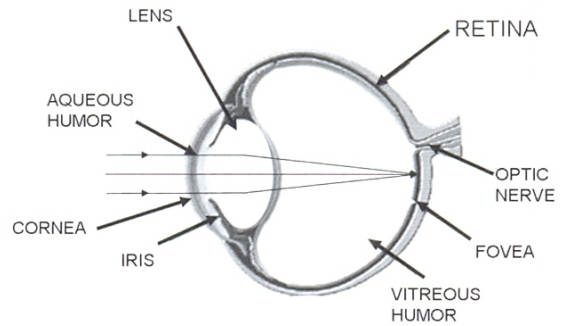
Different Wavelengths



Different Wavelengths



The Human Eye

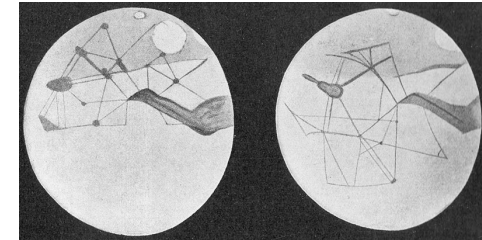
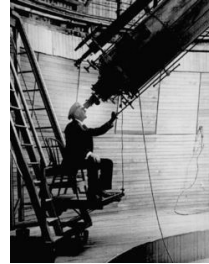


Theoretical: $\theta \sim \lambda/D \sim 0.5\mu m/7mm \sim 14''$
In practice: $\theta \sim 1$ arcminute

The Eye's Computer



Percival Lowell



Canals on Mars!!!

The Eye's Computer

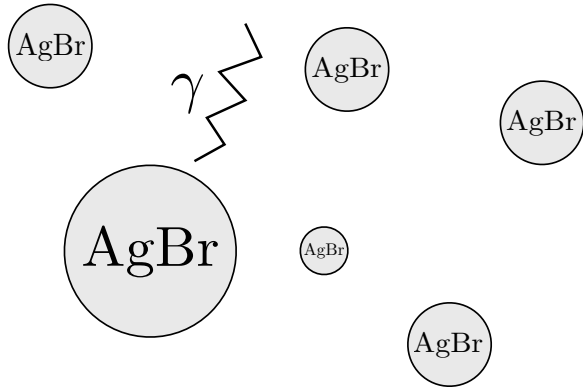


Photographic Plates



Photographic Plate Principle

Expose grains of slightly soluble Silver Halide salts to light:



Photographic Plate Principle

Chemical fixing - remove stray ions and develop ALL silver in a grain:



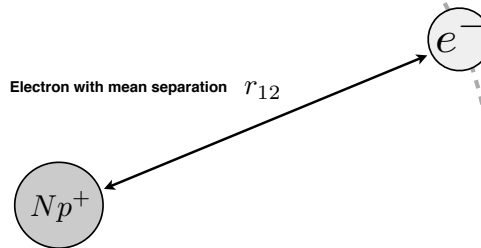
Photodetectors

Classical Mechanics treat electric charges as point particles interacting with electric fields

Electric Potential Energy between two charges:

$$U_E = \frac{1}{4\pi\epsilon_0} \frac{q_1 q_2}{r_{12}}$$

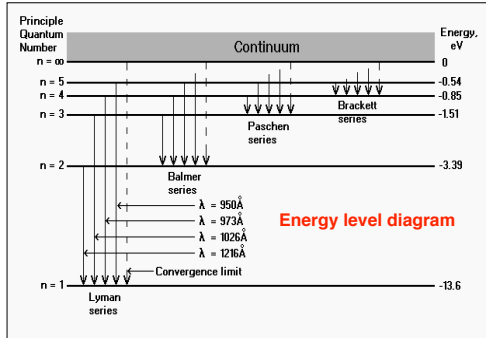
Electron with mean separation r_{12}



Atomic Nucleus with N protons

Electrons can absorb or emit photons and change to a different allowed orbital

e.g. the Hydrogen atom with one electron



Photons of only specific energies can be absorbed or emitted

In the **BOHR MODEL**, the orbital angular momentum of the electron is **quantized** in units of \hbar

$$p_{\theta} = n\hbar$$

...where $n = 1, 2, 3, \dots$

$m = 2$ to $n = 3, 4, 5, 6$



The QM properties of electrons lead to atomic lines and semiconductor bands

Multiple electrons around a positively charged nucleus have four quantum numbers:

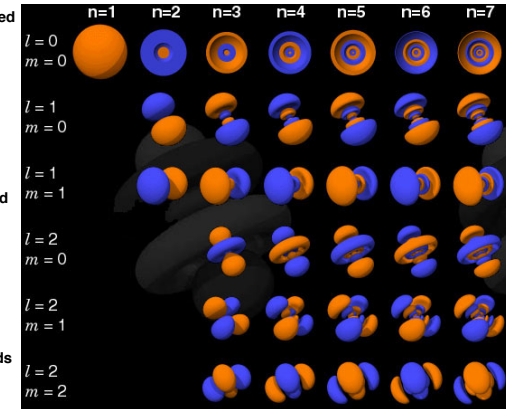
$$n, l, m_l, m_s$$

Only **ONE FERMION** can have one set of quantum numbers!

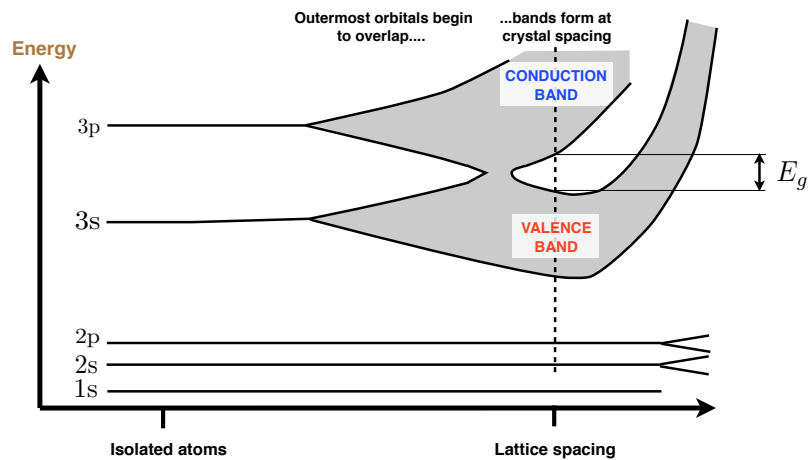
Electrons (and other particles) are described with Schrodinger's Wave Equation:

$$i\hbar \frac{\partial}{\partial t} \Psi(x, t) = \hat{H} \Psi(x, t)$$

Electrons are described by probability clouds called **ORBITALS** with specific energies.



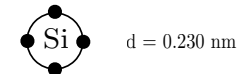
Atomic orbitals overlap in a crystal to form electronic bands



Decreasing atomic separation

Incomplete orbitals provide electrons for bonding

Silicon and Germanium have 4 electrons in their outermost (n=2) orbital:

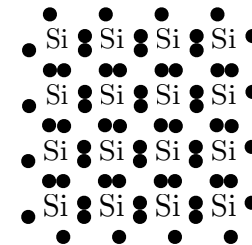


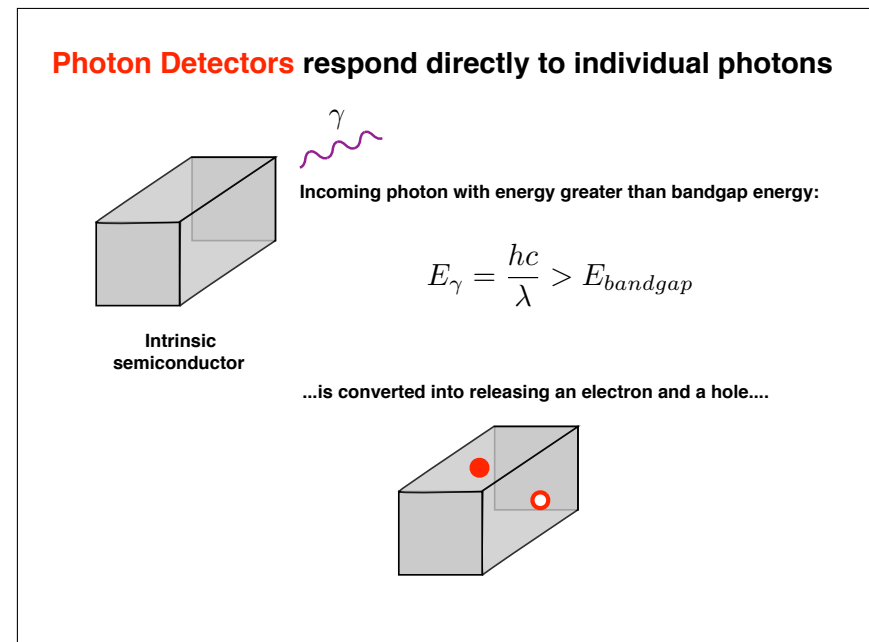
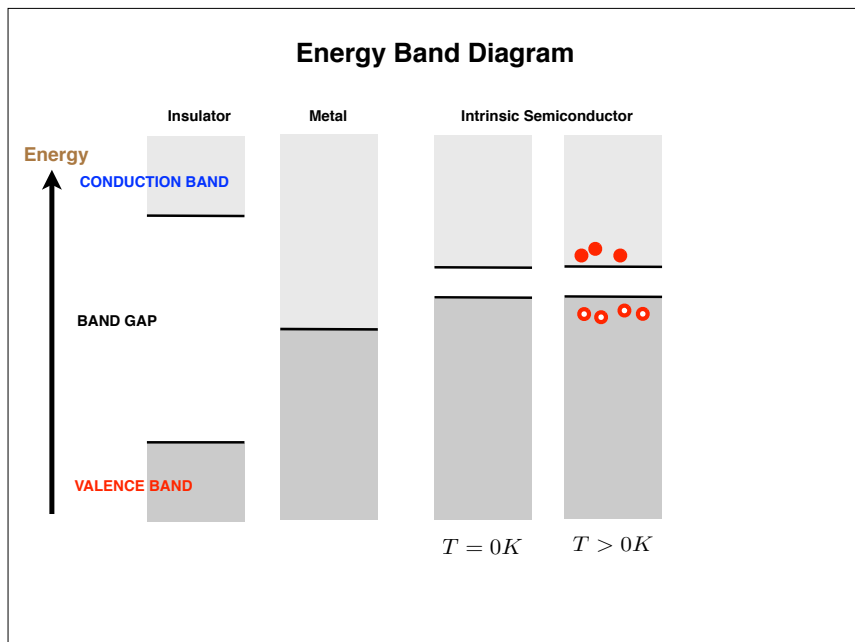
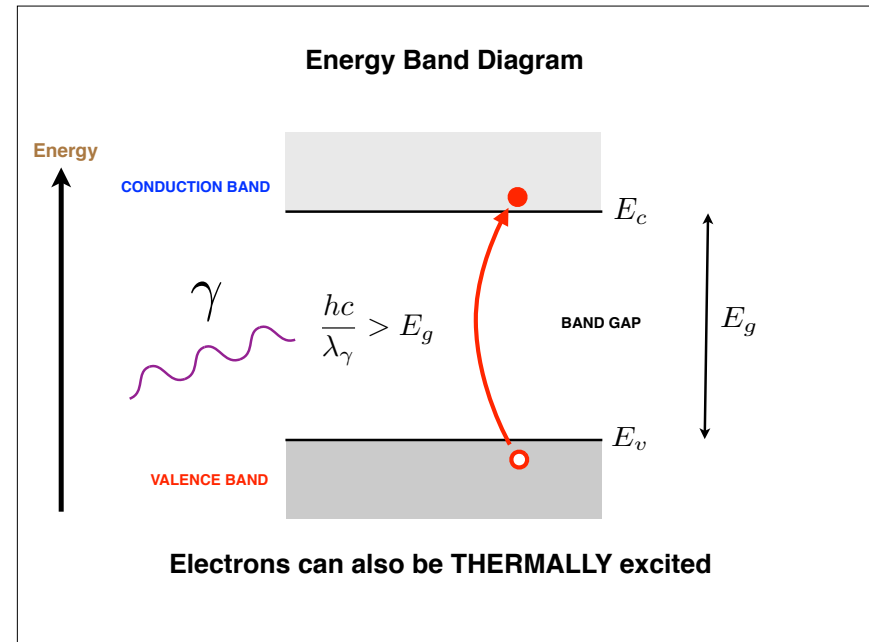
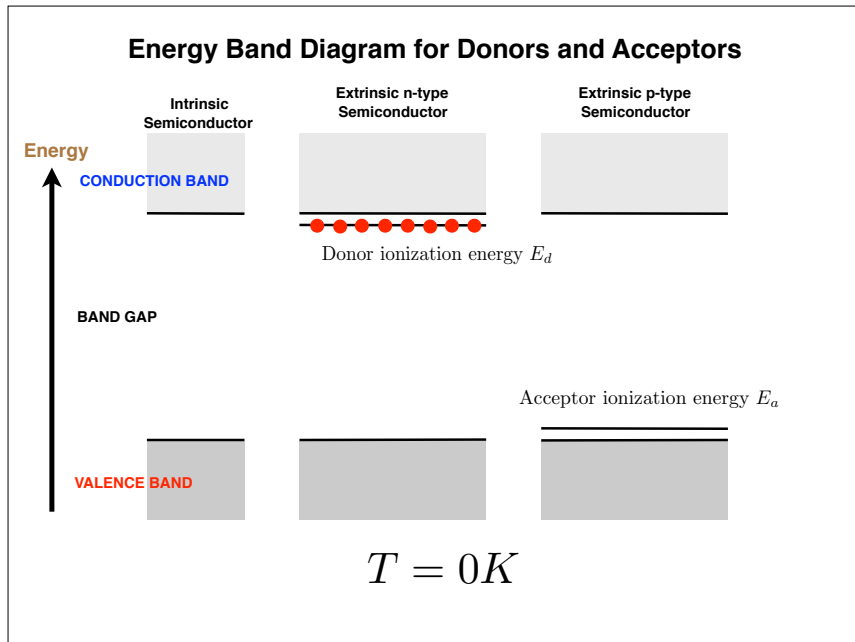
(In the Periodic Table these are **GROUP IV** elements)

Energetically they want to have 8 electrons to form a stable configuration:

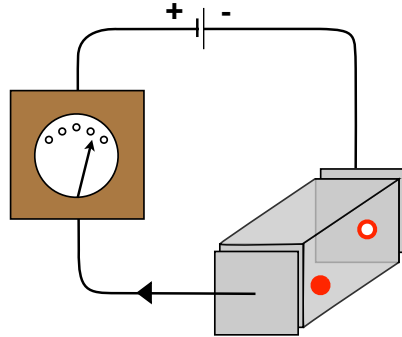


Forming a crystal sharing electrons with other Si atoms forms a stable **LATTICE**:





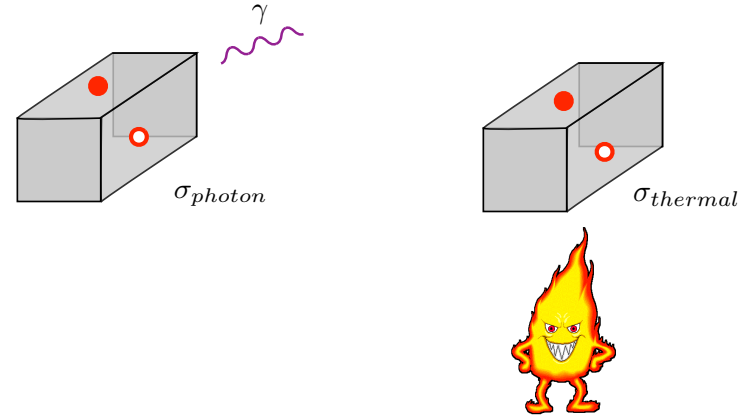
Charge carriers move out of semiconductor and register as a signal



Applying an electric field causes electric charges to move in the material and register a signal as an **electric current**

Charge carriers are generated with both **photons** and **thermal excitation**

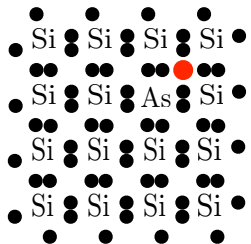
We measure the electrical conductivity!



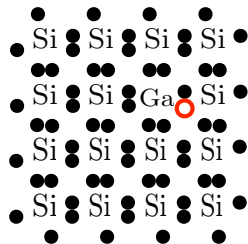
Dopants in Silicon

We can dope a pure silicon crystal with small amounts of **Group V** or **Group III** elements

Adding a **Group V** element introduces conduction electrons and creates **n-type** silicon, called a donor.



Adding a **Group III** element introduces an electron hole and creates **p-type** silicon, called an acceptor.



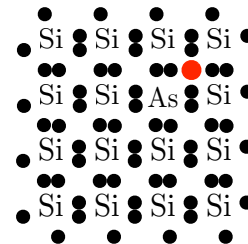
Pure semiconductors are **INTRINSIC**, doped semiconductors are **EXTRINSIC**

Why is a donor easily ionised?



As atom looks "hydrogen-like" with covalent bonds shielding large nuclear charge:

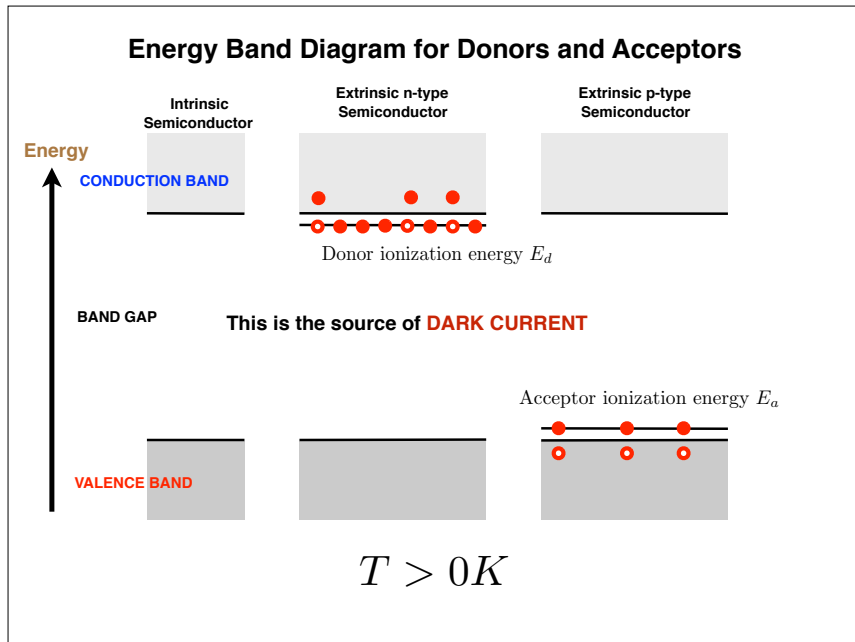
$$E_{\text{bohr}} = \frac{mq^4}{2K^2\hbar^2}$$



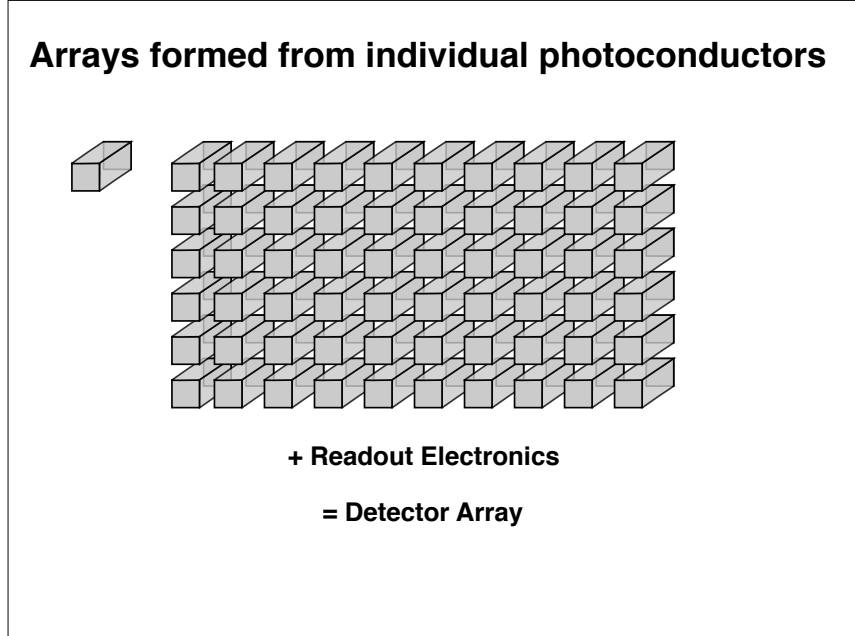
where $K = 4\pi\epsilon_0\epsilon_r$

$$\epsilon_r = 11.8$$

Electron is **REALLY** easy to ionise!



Detector Arrays



Two types of arrays in the Optical/IR

<p>IR Arrays ($1\mu m - 40\mu m$)</p>	<p>Charge Coupled Devices (CCDs) ($0.1\mu m - 1\mu m$)</p>
<p>+ directly access individual pixels</p> <p>- complex and expensive</p>	<p>+ monolithic structure built in Si wafer</p> <p>- charge transfer inefficiencies</p>

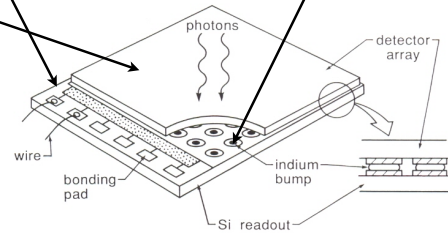
Production of IR Arrays

Make a grid of readout amplifiers in Silicon

Make a matching image of detector pixels

Squeeze them together to make a **hybrid array**

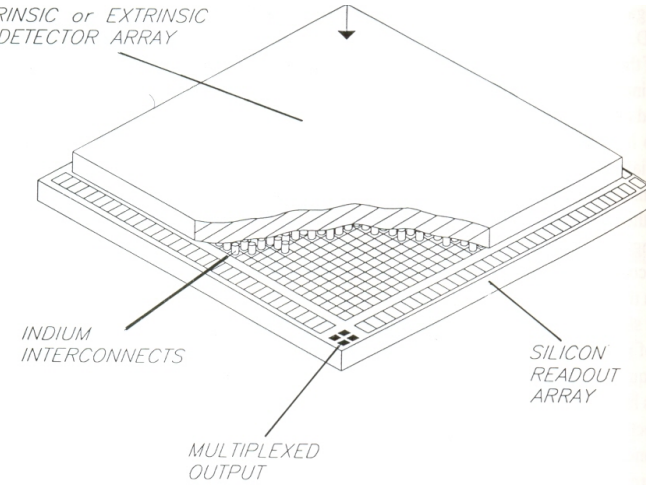
Deposit Indium bumps on both sides



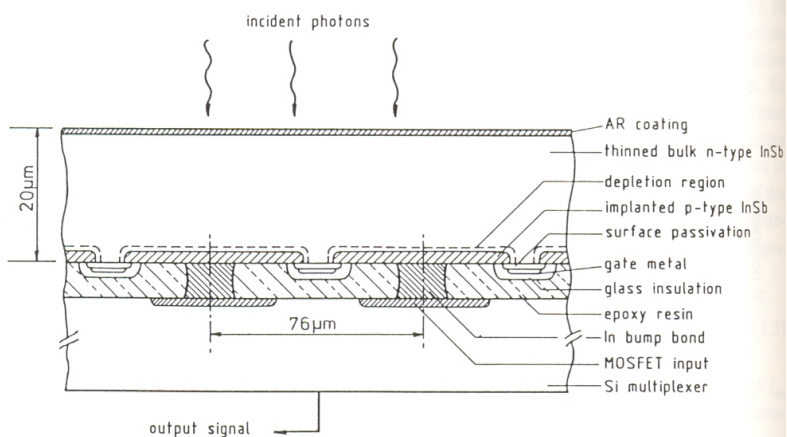
Why Indium? It's a **soft** metal and will still be ductile at cryogenic temperatures!

Production of IR Arrays

INTRINSIC or EXTRINSIC DETECTOR ARRAY

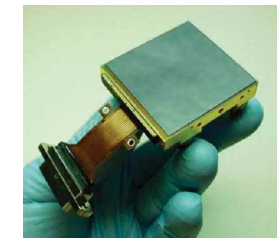


Detailed Bonding Structure of IR Array



The Teledyne 2k x 2k Hawaii-2RG detector

Parameter	Specification
Detector technology	HgCdTe or Si PIN
Detector input circuit	SFD
Readout mode	Ripple
Pixel readout rate	100 kHz to 5MHz (continuously adjustable)
Total pixels	2048 x 2048
Pixel pitch	18 μm
Fill factor	≥ 98%
Output ports	Signal: 1, 4, 32 selectable guide window and reference
Spectral range	0.3 - 5.3μm
Operating temperature	≥ 30K
Quantum efficiency (array mean)	≥ 65%
Charge storage capacity	≥ 100,000e ⁻
Pixel operability	≥ 95%
Dark current (array mean)	≤ 0.1 e ⁻ /sec (77K, 2.5 μm)
Read noise (array mean)	≤ 15 e ⁻ CDS @ 100 kHz
Power dissipation	≤ 4 mW @ 100 kHz



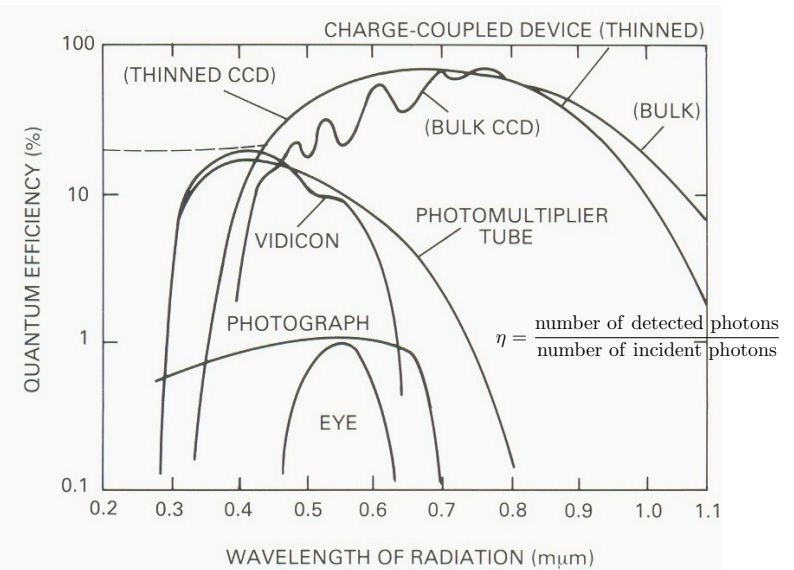
Can also be combined to a 2x2 mosaic



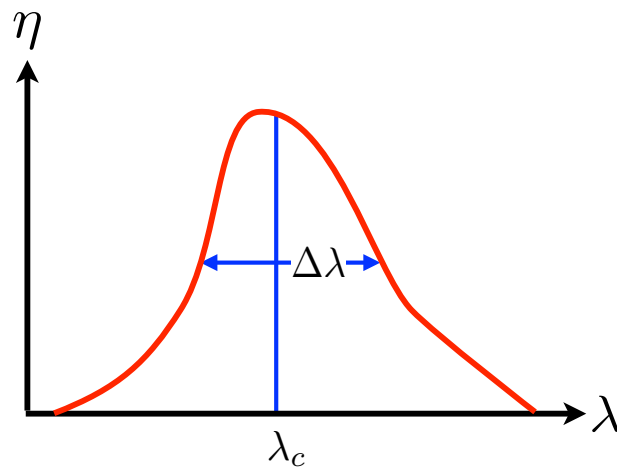
Some Performance Aspects of Detectors

- Spectral response and bandwidth
- Linearity / saturation
- Dynamic range
- Quantum efficiency
- Noise
- Geometric properties
- Time response
- Polarization
- Operational aspects

Spectral Response and Bandwidth



Spectral Response and Bandwidth



Linearity and Dynamic Range



<http://www.luckymanpress.com>

Commercial cameras: 8 to 12 bits $10^{2.4}$ to $10^{3.6}$

Astronomical cameras: 16 bits ++ $10^{4.2}$

Noise

Most important: $\sigma = \frac{\text{Signal}}{\text{Noise}}$ measured as $(S+B) - \text{mean}\{B\}$
 Total noise $= \sqrt{\sum (N_i)^2}$ if statist. independent

Most relevant noise sources:

Photon noise follows Poisson statistics: $P(m) = \frac{e^{-n} n^m}{m!}$

(= probability to detect m photons in a given time interval where, on average, n photons $S/N = \sqrt{n}$)

G-R noise: statistics of the generated and recombined holes and electrons, related to the Poisson statistics of the incoming photons.

Johnson, kTC or reset noise: thermodynamic noise due to the thermal motion of the charge carriers.

1/f noise (increased noise at low frequencies) due to bad electrical contacts, temperature fluctuations, surface effects (damage), crystal defects, JFETs, ...

Noise

Signal: $(S + B) - \text{mean}(B)$

$$\sigma = \frac{\text{Signal}}{\text{Noise}}$$

Noise: can be added as $\sqrt{\sum (N_i)^2}$

Photon noise follows Poisson statistics:

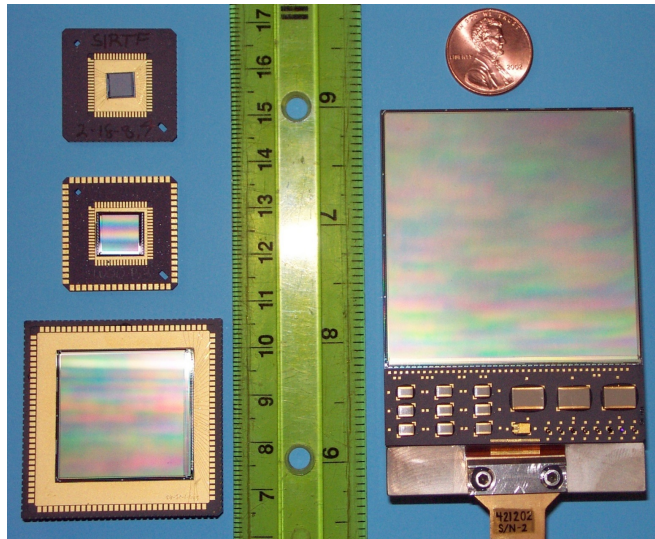
$$P(m) = \frac{e^{-n} n^m}{m!} \quad S/N = \sqrt{n}$$

where $P(m)$ the probability to detect m photons over a time interval and where the mean rate of photons is n

Geometrical Properties

Geometrical dimension and pixel number $x \times y$

4 Generations of Raytheon Infrared Detectors



Calibrating a CCD image

For each **SCIENCE** image S (exposure time t_s)

Subtract off a **BIAS** image B to remove ADC offset (zero time integration)

Subtract off a **DARK** image D to remove dark current offset (exposure time t_d)

Divide by a **FLAT FIELD** image F to remove gain variations (exposure time t_f)

$$S' = \frac{S - \frac{t_s}{t_d}(D - B) - B}{F - \frac{t_f}{t_d}(D - B) - B}$$

- $F - \frac{t_f}{t_d}(D - B) - B$ often normalized such that mean of $S' =$ mean of S

Gain, Read Noise, Saturation limit

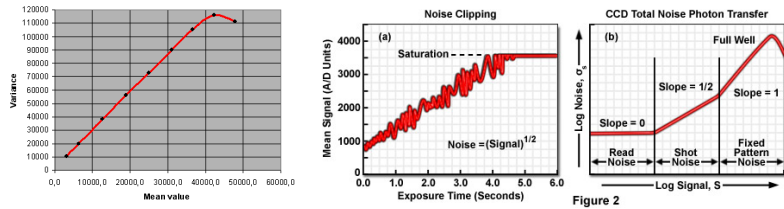
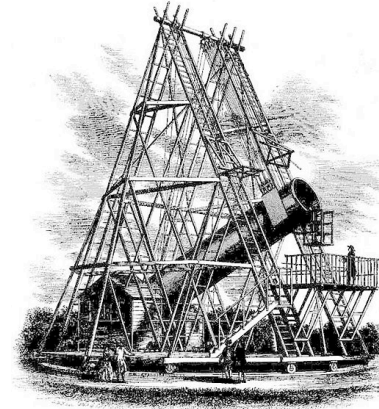


Figure 2

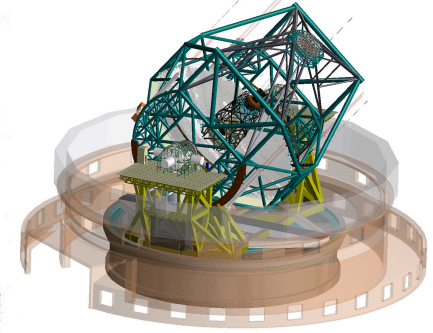
- gain (G) between arbitrary digital units (ADU, A) and number of photo-electrons (e): $A = G \cdot e$
- noise in e is given by $\sigma_e^2 = e$
- and therefore $\sigma_A^2 = G^2 \sigma_e^2 = G^2 e$
- gain G determined from $G = \frac{\sigma_A^2}{A}$

Accurately recording what you see

Looking at the sky without recording it is just tourism



Herschel 1789



E-ELT (2026)

The Atmosphere

Recap: The Atmosphere

Atmosphere is modeled with:

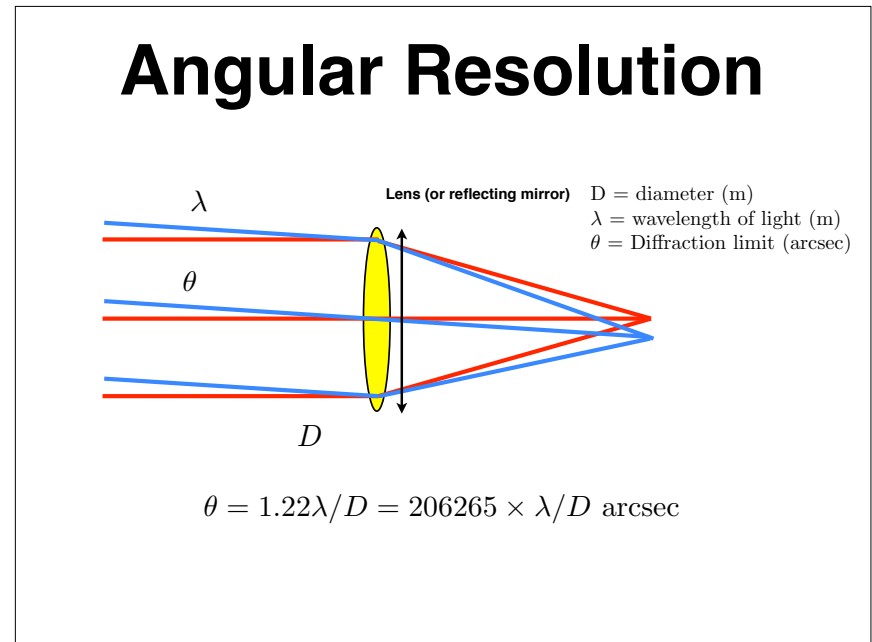
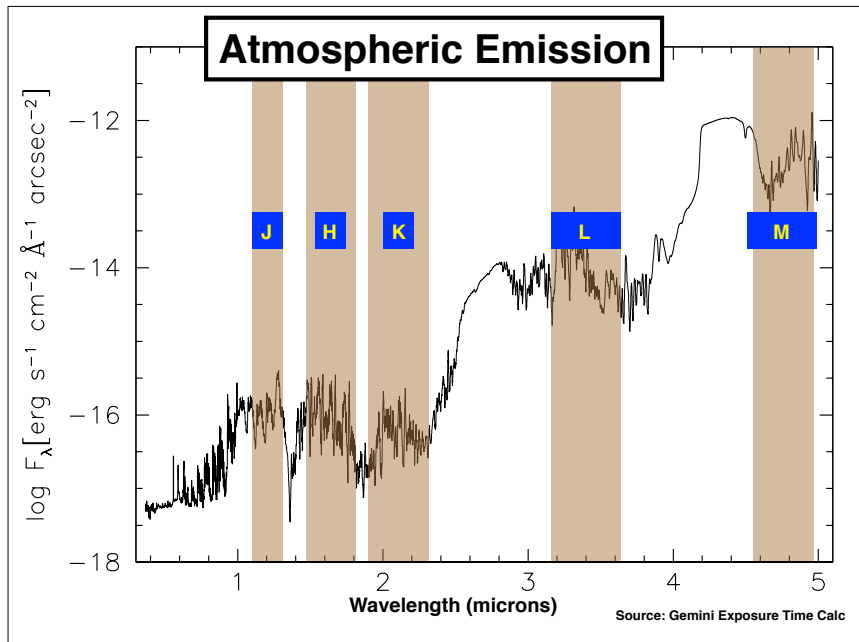
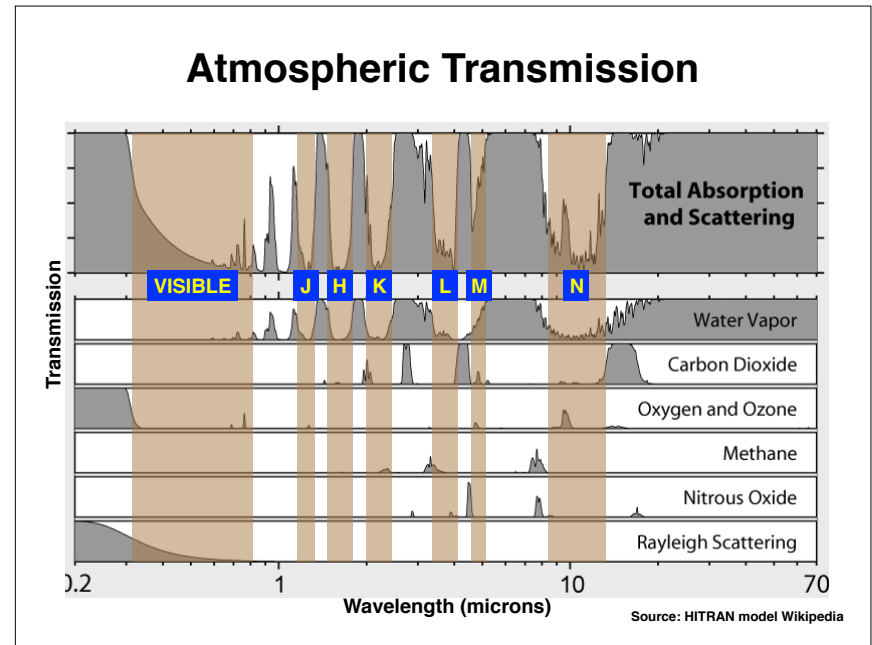
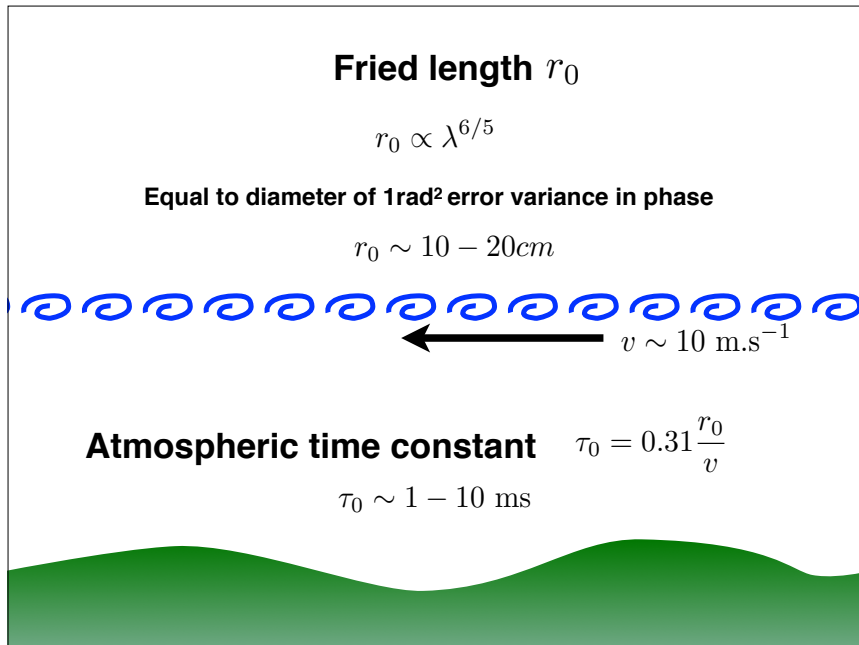
An outer and inner scale length, and a power spectrum of index fluctuations between them

Thin layers of frozen turbulence at 2 to 5 different altitudes

Described with three parameters:

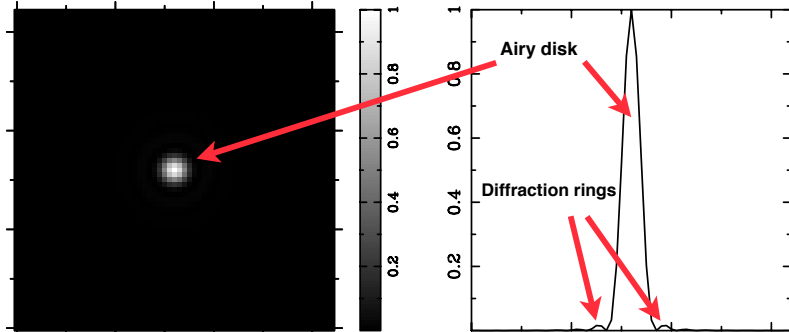
$$r_0, \tau_0 \text{ and } \theta_0$$



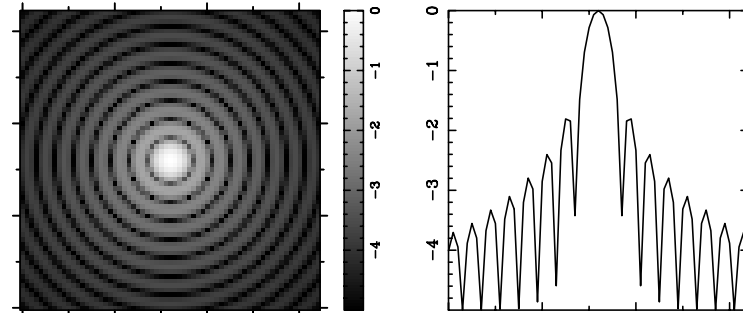


Full Width at Half Maximum (FWHM) and Airy Disk

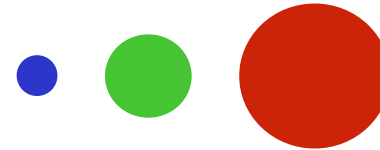
Imaging a point source with a telescope shows a diffraction pattern due to finite size of telescope aperture and wavelike nature of light



FWHM and Airy Disk (logarithmic)



$$\text{FWHM} = 1.22 \frac{\lambda}{D}$$



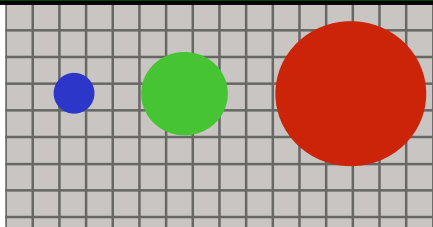
The Goldilocks Detector

There is an OPTIMUM pixel scale for a given wavelength

From Shannon and Nyquist Sampling theorem: $\sim 2.5\text{pix}/\text{FWHM}$

Most AO imagers have a plate scale that matches 2.5 pixels FWHM at the shortest wavelength

< 2 : undersampled



> 4 : oversampled

Astronomers want as much spatial resolution as possible

Diffraction limited by the telescope's primary mirror: $\approx \frac{\lambda}{D_{tel}}$

for the Hubble Space Telescope

$$\approx \frac{0.5\mu m}{2.4m} = 0.2\mu rad$$

$$\approx 43 \text{ milliarcsec}$$



Hubble Space Telescope Credit: NASA

The atmosphere limits diffraction limited imaging

Diffraction limited by the turbulent atmosphere: $\approx \frac{\lambda}{r_0}$

Typically for professional observatories:

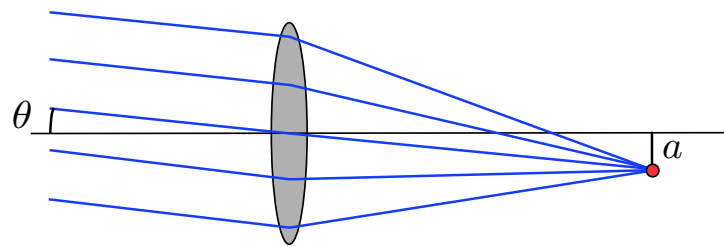
$$\approx \frac{0.5 \mu m}{10 cm} = 5 \mu rad$$

$$\approx 1 \text{ arcsec}$$

Imagers

The plate scale

Simple lens with diameter D and focal length f



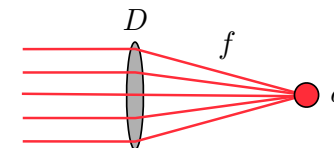
Arcseconds / mm

$$\text{Plate Scale} = \frac{\theta}{a} = \frac{206}{f}$$

metres

NOTE! Plate scale does **NOT** depend on diameter D , only the focal length

Take a seeing limited telescope and double it in size



Focal length doubled - focal ratio kept constant

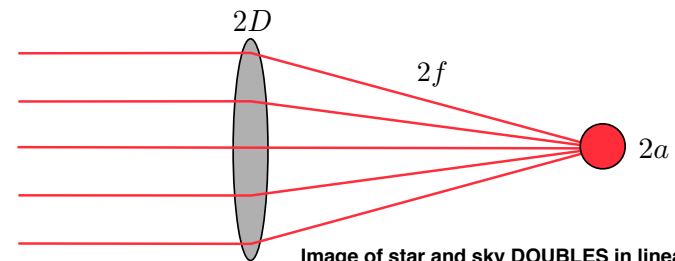
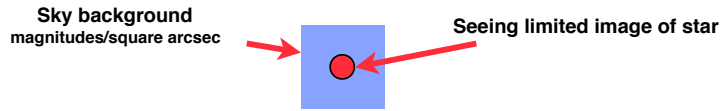


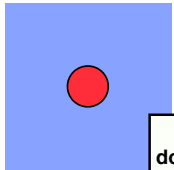
Image of star and sky **DOUBLES** in linear size

Seeing limited telescope sensitivity

$$\text{Signal to noise} = \frac{S}{N} \propto D^2$$



Double all sizes



Focal length constant



4x brighter

Flux/square arcsecond does not change in either case

Diffraction limited telescope sensitivity

Diffraction limited star image is SMALLER than seeing star image

Doubling the diameter



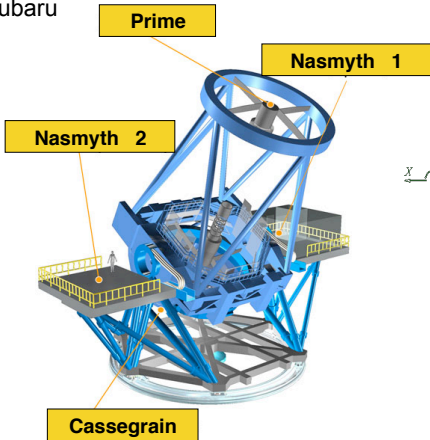
$$A_{PSF} = \pi d_{PSF} \propto \left(\frac{1}{D}\right)^2$$

Area DECREASES for the PSF, so the noise contribution goes down

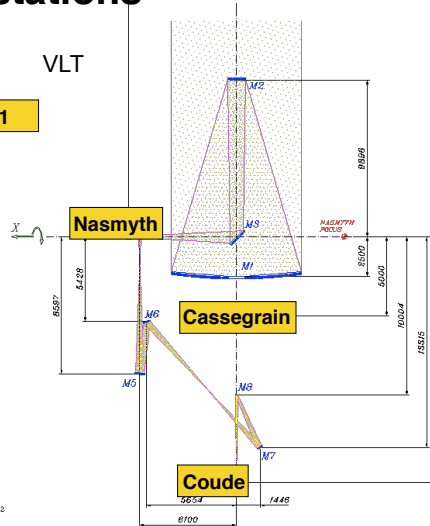
$$\text{Signal to noise} = \frac{S}{N} \propto D^4$$

Focal stations

Subaru



VLT



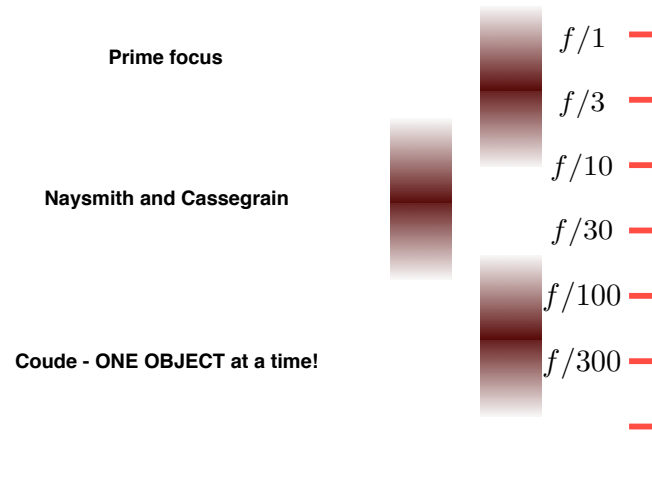
© M&BTA Corporation Japan #150132

Focal stations

Sky rotates with hour angle



Focal ratios for the various foci



Prime Focus Correctors

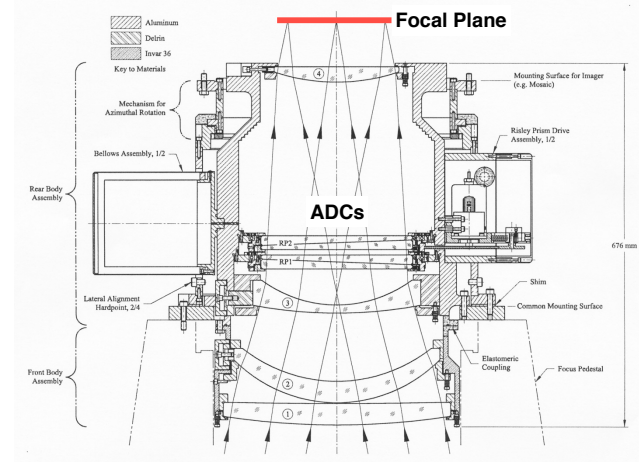


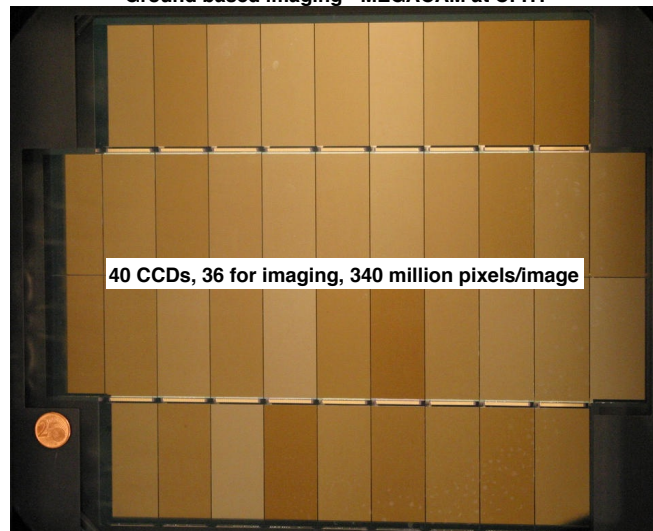
Figure 2. The opto-mechanical assembly for the new 4-m prime focus corrector. Elements #1, #2, #3, and #4 are SiO₂. The ADC materials comprising RP1 and RP2 are UBK7 (rear) and LLF6 (front), where "rear" is closer to the detector

4m Mayall telescope on Kitt Peak

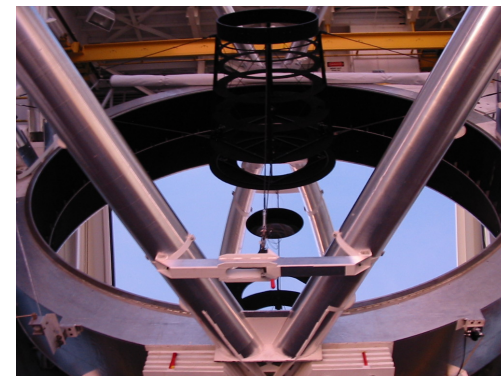
http://www-kpno.kpno.noao.edu/glaspey/mayall_params.html

Prime Focus Imagers

Ground based imaging - MEGACAM at CFHT



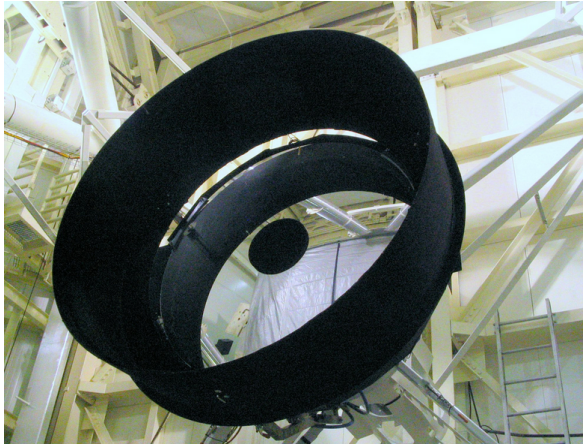
Wide Field Imagers



Telescope baffling above Cassegrain at MMT 6.5m telescope

<https://www.cfa.harvard.edu/~mlacasse/>

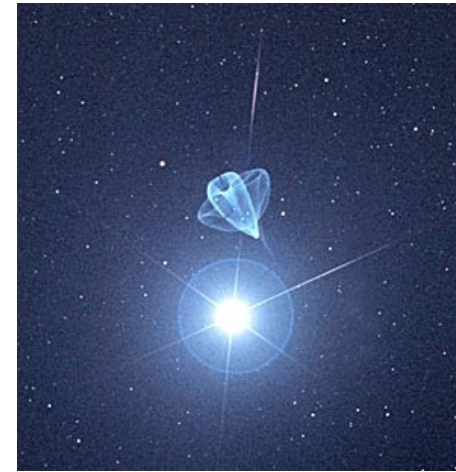
Wide Field Imagers



Telescope baffling for f/5 mirror at MMT0 6.5m telescope

<https://www.cfa.harvard.edu/~mlacasse/>

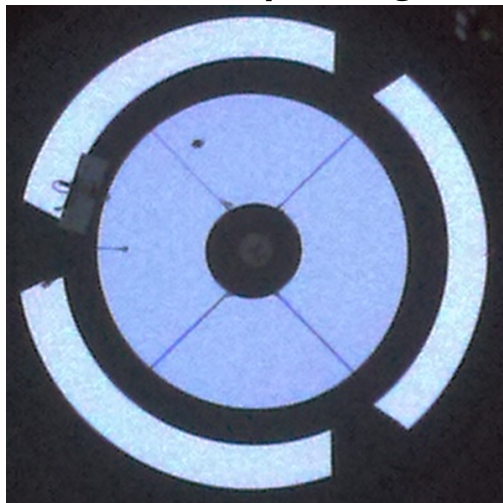
Wide Field Imagers



Internal reflections from a Schmidt camera

<http://www.robertreeves.com/repair1.htm>

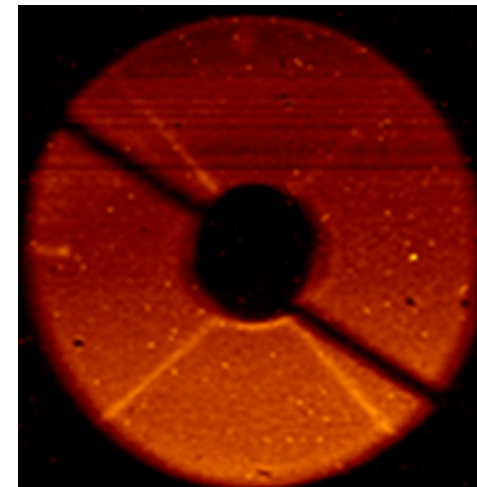
Visible Pupil image



<https://visao.as.arizona.edu/simulations/magao-pupils-and-fourier-optics/>

Kate Morzinski

Infrared 3.4 micron Pupil image

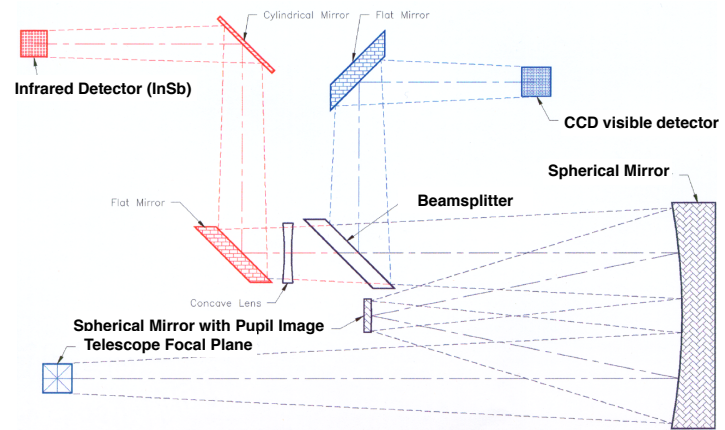


<https://visao.as.arizona.edu/simulations/magao-pupils-and-fourier-optics/>

Kate Morzinski

Offner Relay

Used to make cold stops in IR cameras



<http://www.astronomy.ohio-state.edu/~depoj/research/instrumentation/andicam/andicam.html>