## Lecture 2: Electromagnetic Waves in Isotropic Media 1

## Outline

- Electromagnetic Waves
(2) Quasi-Monochromatic Light
( Electromagnetic Waves Across Interfaces
( - Snell's law


## Fundamentals of Polarized Light

Electromagnetic Waves in Matter

- Maxwell's equations $\Rightarrow$ electromagnetic waves
optics: interaction of electromagnetic waves with matter as
described by material equations
- nolarization of alectromacnotic waves are integral part of optics


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## Maxwell's Equations in Matter

$$
\begin{aligned}
\nabla \cdot \vec{D} & =4 \pi \rho \\
\nabla \times \vec{H}-\frac{1}{c} \frac{\partial \vec{D}}{\partial t} & =\frac{4 \pi}{c} \vec{j} \\
\nabla \times \vec{E}+\frac{1}{c} \frac{\partial \vec{B}}{\partial t} & =0 \\
\nabla \cdot \vec{B} & =0
\end{aligned}
$$

## Symbols

$\vec{D}$ electric displacement
$\rho$ electric charge density magnetic field
c speed of light in vacuum electric current density electric field
magnetic induction
$t$ time

## Linear Material Equations

$$
\begin{gathered}
\vec{D}=\epsilon \vec{E} \\
\vec{B}=\mu \vec{H} \\
\vec{j}=\sigma \vec{E}
\end{gathered}
$$

## Symbols

$\epsilon$ dielectric constant
$\mu$ magnetic permeability $\sigma$ electrical conductivity

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## Symbols

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## Isotropic and Anisotropic Media

- isotropic media: $\epsilon$ and $\mu$ are scalars
- anisotropic media: $\epsilon$ and $\mu$ are tensors of rank 2
- isotropy of medium broken by
- anisotropy of material itself (e.g. crystals)
- external fields (e.g. Kerr effect)


## Wave Equation in Matter

- for most materials: $\rho=0, \mu=1$
- combine Maxwell, material equations $\Rightarrow$ differential equations for damped (vector) wave
- damping controlled by conductivity $\sigma$ - $\vec{E}$ and $\vec{H}$ are equivalent $\Rightarrow$ sufficient to consider $\vec{E}$


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& \nabla^{2} \vec{H}-\frac{\epsilon}{c^{2}} \frac{\partial^{2} \vec{H}}{\partial t^{2}}-\frac{4 \pi \sigma}{c^{2}} \frac{\partial \vec{H}}{\partial t}=0
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\end{aligned}
$$

- damping controlled by conductivity $\sigma$
- $\vec{E}$ and $\vec{H}$ are equivalent $\Rightarrow$ sufficient to consider $\vec{E}$
- interaction with matter almost always through $\vec{E}$
- but: at interfaces, boundary conditions for $\vec{H}$ are crucial


## Plane-Wave Solutions

- Plane Vector Wave ansatz $\vec{E}=\vec{E}_{0} e^{i(\vec{k} \cdot \vec{x}-\omega t)}$
$\vec{k}$ spatially and temporally constant wave vector
$\vec{k}$ normal to surfaces of constant phase
$|\vec{k}|$ wave number
$\vec{x}$ spatial location
$\omega$ angular frequency ( $2 \pi \times$ frequency)
$t$ time
$\vec{E}_{0}$ (generally complex) vector independent of time and space


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- could also use $\vec{E}=\vec{E}_{0} e^{-i(\vec{k} \cdot \vec{x}-\omega t)}$
- damping if $\vec{k}$ is complex
- real electric field vector given by real part of $\vec{E}$


## Complex Index of Refraction

- temporal derivatives $\Rightarrow$ Helmholtz equation

$$
\nabla^{2} \vec{E}+\frac{\omega^{2}}{c^{2}}\left(\epsilon+i \frac{4 \pi \sigma}{\omega}\right) \vec{E}=0
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- dispersion relation between $\vec{k}$ and $\omega$

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$$

- split into real ( $n$ : index of refraction) and imaginary parts ( $k$ : extinction coefficient)

$$
\tilde{n}=n+i k
$$

## Transverse Waves

## ${ }^{1}{ }^{1}$ A A An

- plane-wave solution must also fulfill Maxwell's equations

$$
\vec{E}_{0} \cdot \vec{k}=0, \quad \vec{H}_{0} \cdot \vec{k}=0, \quad \vec{H}_{0}=\tilde{n} \frac{\vec{k}}{|\vec{k}|} \times \vec{E}_{0}
$$

- isotropic media: electric, magnetic field vectors normal to wave
vector $\Rightarrow$ transverse waves
- $\vec{E}_{0}, \vec{H}_{0}$, and $\vec{k}$ orthogonal to each other, right-handed vector-triple


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- $\vec{E}_{0}, \vec{H}_{0}$, and $\vec{k}$ orthogonal to each other, right-handed vector-triple - conductive medium $\Rightarrow$ complex $\tilde{n}, E_{0}$ and $H_{0}$ out of phase


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- conductive medium $\Rightarrow$ complex $\tilde{n}, \vec{E}_{0}$ and $\vec{H}_{0}$ out of phase
- $\vec{E}_{0}$ and $\vec{H}_{0}$ have constant relationship $\Rightarrow$ consider only $\vec{E}$


## Energy Propagation in Isotropic Media

- Poynting vector

$$
\vec{S}=\frac{c}{4 \pi}(\vec{E} \times \vec{H})
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- $|\vec{S}|$ : energy through unit area perpendicular to $\vec{S}$ per unit time
- direction of $\vec{S}$ is direction of energy flow
$\square$


## Energy Propagation in Isotropic Media

- Poynting vector

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\vec{S}=\frac{c}{4 \pi}(\vec{E} \times \vec{H})
$$

- |S|: energy through unit area perpendic

$$
\langle\vec{S}\rangle=\frac{c}{8 \pi} \operatorname{Re}\left(\vec{E}_{0} \times \vec{H}_{0}^{*}\right)
$$

Re real part of complex expression

* complex conjugate

〈.) time average

- energy flow parallel to wave vector (in isotropic media)


## Energy Propagation in Isotropic Media

- $|\vec{S}|$ : energy through unit area perpendicular to $\vec{S}$ per unit time
- time-averaged Poynting vector given by

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$$
\langle\vec{S}\rangle=\frac{c}{8 \pi}|\tilde{n}|\left|E_{0}\right|^{2} \frac{\vec{k}}{|\vec{k}|}
$$

## Quasi-Monochromatic Light

- monochromatic light: purely theoretical concept
- monochromatic light wave always fully polarized

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- monochromatic light: purely theoretical concept
- monochromatic light wave always fully polarized
- real life: light includes range of wavelengths $\Rightarrow$ quasi-monochromatic light



## Quasi-Monochromatic Light

- real life: light includes range of wavelengths $\Rightarrow$ quasi-monochromatic light
- quasi-monochromatic: superposition of mutually incoherent monochromatic light beams whose wavelengths vary in narrow range $\delta \lambda$ around central wavelength $\lambda_{0}$

$$
\frac{\delta \lambda}{\lambda} \ll 1
$$

measurement of quasi-monochromatic light: integral over measurement time $t_{m}$
amplitude. phase (slow) functions of time for given spatial locatior slow: variations occur on time scales much longer than the mean

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- measurement of quasi-monochromatic light: integral over measurement time $t_{m}$
- amplitude, phase (slow) functions of time for given spatial location
- slow: variations occur on time scales much longer than the mean period of the wave


## Polarization of Quasi-Monochromatic Light

- electric field vector for quasi-monochromatic plane wave is sum of electric field vectors of all monochromatic beams

$$
\vec{E}(t)=\vec{E}_{0}(t) e^{i(\vec{k} \cdot \vec{x}-\omega t)}
$$

- can write this way because $\delta \lambda \ll \lambda_{0}$
- measured intensity of quasi-monochromatic beam

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$$
\left\langle\vec{E}_{x} \vec{E}_{x}^{*}\right\rangle+\left\langle\vec{E}_{y} \vec{E}_{y}^{*}\right\rangle=\lim _{t_{m} \rightarrow>\infty} \frac{1}{t_{m}} \int_{-t_{m} / 2}^{t_{m} / 2} \vec{E}_{x}(t) \vec{E}_{x}^{*}(t)+\vec{E}_{y}(t) \vec{E}_{y}^{*}(t) d t
$$

$\langle\cdots\rangle$ : averaging over measurement time $t_{m}$

- measured intensity independent of time
- quasi-monochromatic: frequency-dependent material properties (e.g. index of refraction) are constant within $\Delta \lambda$


## Polychromatic Light or White Light

- wavelength range comparable wavelength ( $\frac{\delta \lambda}{\lambda} \sim 1$ )
- incoherent sum of quasi-monochromatic beams that have large variations in wavelength
cannot write electric field vector in a plane-wave form
must take into account frequency-dependent material
characteristics


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## Electromagnetic Waves Across Interfaces

## Fields at Interfaces

- classical optics due to interfaces between 2 different media
medium 1 to medium 2



## Electromagnetic Waves Across Interfaces

## Fields at Interfaces

- from Maxwell's equations in integral form at interface from medium 1 to medium 2

$$
\begin{aligned}
\left(\vec{D}_{2}-\vec{D}_{1}\right) \cdot \vec{n} & =0 \\
\left(\vec{B}_{2}-\vec{B}_{1}\right) \cdot \vec{n} & =0 \\
\left(\vec{E}_{2}-\vec{E}_{1}\right) \times \vec{n} & =0 \\
\left(\vec{H}_{2}-\vec{H}_{1}\right) \times \vec{n} & =0
\end{aligned}
$$

$\vec{n}$ normal on interface, points from medium 1 to medium 2 - normal components of $D$ and $B$ are continuous across interface - tangential components of $E$ and $H$ are continuous across interface

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## Plane of Incidence

- plane wave onto interface
- incident ( ${ }^{( }$), reflected $\left({ }^{r}\right)$, and transmitted $\left({ }^{t}\right)$ waves

$$
\begin{aligned}
& \vec{E}^{i, r, t}=\vec{E}_{0}^{i, r, t} e^{i\left(\vec{k}^{i} r, t, \vec{x}-\omega t\right)} \\
& \vec{H}^{i, r, t}=\frac{c}{\omega} \vec{k}^{i, r, t} \times \vec{E}^{i, r, t}
\end{aligned}
$$

- interface normal $\vec{n}|\mid z$-axis

- spatial, temporal behavior at interface the same for all 3 waves - valid for all $\vec{x}$ in interface $\Rightarrow$ all 3 wave vectors in one plane, plane of incidence


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$$
\left(\vec{k}^{i} \cdot \vec{x}\right)_{z=0}=\left(\vec{k}^{r} \cdot \vec{x}\right)_{z=0}=\left(\vec{k}^{t} \cdot \vec{x}\right)_{z=0}
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## Snell's Law

- spatial, temporal behavior the same for all three waves
$\left(\vec{k}^{i} \cdot \vec{x}\right)_{z=0}=\left(\vec{k}^{r} \cdot \vec{x}\right)_{z=0}=\left(\vec{k}^{t} \cdot \vec{x}\right)_{z=0}$
- $|\vec{k}|=\frac{\omega}{c} \tilde{n}$
- $\omega, c$ the same for all 3 waves
- Snell's law



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- $|\vec{k}|=\frac{\omega}{c} \tilde{n}$
- $\omega, c$ the same for all 3 waves Snell's law $\tilde{n}_{1} \sin \theta_{i}=\tilde{n}_{1} \sin \theta_{r}=\tilde{n}_{2} \sin \theta_{t}$



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## Monochromatic Wave at Interface

$$
\vec{H}_{0}^{i, r, t}=\frac{c}{\omega} \vec{k}^{i, r, t} \times \vec{E}_{0}^{i, r, t}, \quad \vec{B}_{0}^{i, r, t}=\frac{c}{\omega} \vec{k}^{i, r, t} \times \vec{E}_{0}^{i, r, t}
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- 4 equations are not independent


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$$

- boundary conditions for monochromatic plane wave:

$$
\begin{array}{r}
\left(\tilde{n}_{1}^{2} \vec{E}_{0}^{i}+\tilde{n}_{1}^{2} \vec{E}_{0}^{r}-\tilde{n}_{2}^{2} \vec{E}_{0}^{t}\right) \cdot \vec{n}=0 \\
\left(\vec{k}^{i} \times \vec{E}_{0}^{i}+\vec{k}^{r} \times \vec{E}_{0}^{r}-\vec{k}^{t} \times \vec{E}_{0}^{t}\right) \cdot \vec{n}=0 \\
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\left(\vec{k}^{i} \times \vec{E}_{0}^{i}+\vec{k}^{r} \times \vec{E}_{0}^{r}-\vec{k}^{t} \times \vec{E}_{0}^{t}\right) \times \vec{n}=0
\end{array}
$$

- 4 equations are not independent
- only need to consider last two equations (tangential components of $\vec{E}_{0}$ and $\vec{H}_{0}$ are continuous)


## Two Special Cases


(1) electric field parallel to plane of incidence $\Rightarrow$ magnetic field is transverse to plane of incidence (TM)
(2) electric field particular (German: senkrecht) or transverse to plane of incidence (TE)

- general solution as (coherent) superposition of two cases
- choose direction of magnetic field vector such that Poynting vector
parallel, same direction as corresponding wave vector


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