# Lecture 2: Electromagnetic Waves in Isotropic Media 1

#### Outline

- Electromagnetic Waves
- Quasi-Monochromatic Light
- Electromagnetic Waves Across Interfaces
- Snell's law

## Electromagnetic Waves in Matter

- *Maxwell's equations* ⇒ electromagnetic waves
- optics: interaction of electromagnetic waves with matter as described by *material equations*
- polarization of electromagnetic waves are integral part of optics

#### Electromagnetic Waves in Matter

- *Maxwell's equations* ⇒ electromagnetic waves
- optics: interaction of electromagnetic waves with matter as described by *material equations*

polarization of electromagnetic waves are integral part of optics

$\nabla \cdot \vec{D} = 4\pi\rho$ $\nabla \times \vec{H} - \frac{1}{c} \frac{\partial \vec{D}}{\partial t} = \frac{4\pi}{c} \vec{j}$ $\nabla \times \vec{E} + \frac{1}{c} \frac{\partial \vec{B}}{\partial t} = 0$ $\nabla \cdot \vec{B} = 0$	

#### Electromagnetic Waves in Matter

- *Maxwell's equations* ⇒ electromagnetic waves
- optics: interaction of electromagnetic waves with matter as described by material equations
- polarization of electromagnetic waves are integral part of optics

$\nabla \cdot \vec{D} = 4\pi\rho$ $\nabla \times \vec{H} - \frac{1}{c} \frac{\partial \vec{D}}{\partial t} = \frac{4\pi}{c} \vec{j}$ $\nabla \times \vec{E} + \frac{1}{c} \frac{\partial \vec{B}}{\partial t} = 0$ $\nabla \cdot \vec{B} = 0$	

#### Electromagnetic Waves in Matter

- *Maxwell's equations* ⇒ electromagnetic waves
- optics: interaction of electromagnetic waves with matter as described by material equations
- polarization of electromagnetic waves are integral part of optics

Maxwell's Equations in Matter	Symbols
$\nabla \cdot \vec{D} = 4\pi\rho$ $\nabla \times \vec{H} - \frac{1}{c} \frac{\partial \vec{D}}{\partial t} = \frac{4\pi}{c} \vec{j}$ $\nabla \times \vec{E} + \frac{1}{c} \frac{\partial \vec{B}}{\partial t} = 0$ $\nabla \cdot \vec{B} = 0$	$\vec{D}$ electric displacement $\rho$ electric charge density $\vec{H}$ magnetic field c speed of light in vacuum $\vec{j}$ electric current density $\vec{E}$ electric field $\vec{B}$ magnetic induction t time

#### Linear Material Equations

$$\vec{D} = \epsilon \vec{E}$$
$$\vec{B} = \mu \vec{H}$$
$$\vec{j} = \sigma \vec{E}$$

#### Symbols

- € dielectric constant
- $\mu$  magnetic permeability
- $\sigma$  electrical conductivity

#### Isotropic and Anisotropic Media

- isotropic media:  $\epsilon$  and  $\mu$  are scalars
- anisotropic media:  $\epsilon$  and  $\mu$  are tensors of rank 2
- isotropy of medium broken by
  - anisotropy of material itself (e.g. crystals)
  - external fields (e.g. Kerr effect)

#### Linear Material Equations

$$\vec{D} = \epsilon \vec{E}$$
$$\vec{B} = \mu \vec{H}$$
$$\vec{j} = \sigma \vec{E}$$

#### Symbols

- $\epsilon$  dielectric constant
- $\mu$  magnetic permeability
- $\sigma$  electrical conductivity

#### Isotropic and Anisotropic Media

- isotropic media:  $\epsilon$  and  $\mu$  are scalars
- anisotropic media:  $\epsilon$  and  $\mu$  are tensors of rank 2
- isotropy of medium broken by
  - anisotropy of material itself (e.g. crystals)
  - external fields (e.g. Kerr effect)

## • for most materials: $\rho = 0$ , $\mu = 1$

 combine Maxwell, material equations ⇒ differential equations for damped (vector) wave

$$\nabla^2 \vec{E} - \frac{\epsilon}{c^2} \frac{\partial^2 \vec{E}}{\partial t^2} - \frac{4\pi\sigma}{c^2} \frac{\partial \vec{E}}{\partial t} = 0$$
$$\nabla^2 \vec{H} - \frac{\epsilon}{c^2} \frac{\partial^2 \vec{H}}{\partial t^2} - \frac{4\pi\sigma}{c^2} \frac{\partial \vec{H}}{\partial t} = 0$$

- damping controlled by conductivity  $\sigma$
- $\vec{E}$  and  $\vec{H}$  are equivalent  $\Rightarrow$  sufficient to consider  $\vec{E}$
- interaction with matter almost always through  $\vec{E}$
- but: at interfaces, boundary conditions for H
   are crucial

- for most materials:  $\rho = 0, \mu = 1$
- combine Maxwell, material equations ⇒ differential equations for damped (vector) wave

$$\nabla^{2}\vec{E} - \frac{\epsilon}{c^{2}}\frac{\partial^{2}\vec{E}}{\partial t^{2}} - \frac{4\pi\sigma}{c^{2}}\frac{\partial\vec{E}}{\partial t} = 0$$
$$\nabla^{2}\vec{H} - \frac{\epsilon}{c^{2}}\frac{\partial^{2}\vec{H}}{\partial t^{2}} - \frac{4\pi\sigma}{c^{2}}\frac{\partial\vec{H}}{\partial t} = 0$$

- damping controlled by conductivity  $\sigma$
- $\vec{E}$  and  $\vec{H}$  are equivalent  $\Rightarrow$  sufficient to consider  $\vec{E}$
- interaction with matter almost always through  $\vec{E}$
- but: at interfaces, boundary conditions for  $\vec{H}$  are crucial

#### • for most materials: $\rho = 0$ , $\mu = 1$

 combine Maxwell, material equations ⇒ differential equations for damped (vector) wave

$$\nabla^{2}\vec{E} - \frac{\epsilon}{c^{2}}\frac{\partial^{2}\vec{E}}{\partial t^{2}} - \frac{4\pi\sigma}{c^{2}}\frac{\partial\vec{E}}{\partial t} = 0$$
$$\nabla^{2}\vec{H} - \frac{\epsilon}{c^{2}}\frac{\partial^{2}\vec{H}}{\partial t^{2}} - \frac{4\pi\sigma}{c^{2}}\frac{\partial\vec{H}}{\partial t} = 0$$

- damping controlled by conductivity  $\sigma$
- $\vec{E}$  and  $\vec{H}$  are equivalent  $\Rightarrow$  sufficient to consider  $\vec{E}$
- interaction with matter almost always through  $\vec{E}$
- but: at interfaces, boundary conditions for H
   are crucial

#### • for most materials: $\rho = 0, \mu = 1$

 combine Maxwell, material equations ⇒ differential equations for damped (vector) wave

$$\nabla^{2}\vec{E} - \frac{\epsilon}{c^{2}}\frac{\partial^{2}\vec{E}}{\partial t^{2}} - \frac{4\pi\sigma}{c^{2}}\frac{\partial\vec{E}}{\partial t} = 0$$
$$\nabla^{2}\vec{H} - \frac{\epsilon}{c^{2}}\frac{\partial^{2}\vec{H}}{\partial t^{2}} - \frac{4\pi\sigma}{c^{2}}\frac{\partial\vec{H}}{\partial t} = 0$$

- damping controlled by conductivity  $\sigma$
- $\vec{E}$  and  $\vec{H}$  are equivalent  $\Rightarrow$  sufficient to consider  $\vec{E}$
- interaction with matter almost always through  $\vec{E}$
- but: at interfaces, boundary conditions for  $\vec{H}$  are crucial

- Plane Vector Wave ansatz  $\vec{E} = \vec{E}_0 e^{i(\vec{k}\cdot\vec{x}-\omega t)}$ 
  - $\vec{k}$  spatially and temporally constant wave vector
  - $\vec{k}$  normal to surfaces of constant phase
  - $\vec{k}$  wave number
  - $\vec{x}$  spatial location
  - $\omega$  angular frequency (2 $\pi$ × frequency)
  - t time
  - $\vec{E}_0$  (generally complex) vector independent of time and space
- could also use  $ec{E}=ec{E}_0e^{-i\left(ec{k}\cdotec{x}-\omega t
  ight)}$
- damping if  $\vec{k}$  is complex
- real electric field vector given by real part of *E*

- Plane Vector Wave ansatz  $\vec{E} = \vec{E}_0 e^{i(\vec{k}\cdot\vec{x}-\omega t)}$ 
  - $\vec{k}$  spatially and temporally constant wave vector
  - $\vec{k}$  normal to surfaces of constant phase
  - $\vec{k}$  wave number
  - $\vec{x}$  spatial location
  - $\omega$  angular frequency (2 $\pi$ × frequency)
  - t time
  - $\vec{E}_0$  (generally complex) vector independent of time and space

• could also use 
$$\vec{E} = \vec{E}_0 e^{-i(\vec{k}\cdot\vec{x}-\omega t)}$$

• damping if  $\vec{k}$  is complex

real electric field vector given by real part of *E*

- Plane Vector Wave ansatz  $\vec{E} = \vec{E}_0 e^{i(\vec{k}\cdot\vec{x}-\omega t)}$ 
  - $\vec{k}$  spatially and temporally constant wave vector
  - $\vec{k}$  normal to surfaces of constant phase
  - $\vec{k}$  wave number
  - $\vec{x}$  spatial location
  - $\omega$  angular frequency (2 $\pi$ × frequency)
  - t time
  - $\vec{E}_0$  (generally complex) vector independent of time and space
- could also use  $\vec{E} = \vec{E}_0 e^{-i(\vec{k}\cdot\vec{x}-\omega t)}$
- damping if  $\vec{k}$  is complex

real electric field vector given by real part of É

- Plane Vector Wave ansatz  $\vec{E} = \vec{E}_0 e^{i(\vec{k}\cdot\vec{x}-\omega t)}$ 
  - $\vec{k}$  spatially and temporally constant wave vector
  - $\vec{k}$  normal to surfaces of constant phase
  - $\vec{k}$  wave number
  - $\vec{x}$  spatial location
  - $\omega$  angular frequency (2 $\pi$ × frequency)
  - t time
  - $\vec{E}_0$  (generally complex) vector independent of time and space
- could also use  $\vec{E} = \vec{E}_0 e^{-i(\vec{k}\cdot\vec{x}-\omega t)}$
- damping if  $\vec{k}$  is complex
- real electric field vector given by real part of  $\vec{E}$

#### temporal derivatives ⇒ Helmholtz equation

$$\nabla^2 \vec{E} + \frac{\omega^2}{c^2} \left( \epsilon + i \frac{4\pi\sigma}{\omega} \right) \vec{E} = 0$$

• dispersion relation between k and  $\omega$ 

$$\vec{k} \cdot \vec{k} = rac{\omega^2}{c^2} \left( \epsilon + i rac{4\pi\sigma}{\omega} 
ight)$$

complex index of refraction

$$\tilde{n}^2 = \epsilon + i \frac{4\pi\sigma}{\omega}, \ \vec{k} \cdot \vec{k} = \frac{\omega^2}{c^2} \tilde{n}^2$$

 split into real (n: index of refraction) and imaginary parts (k: extinction coefficient)

$$\tilde{n} = n + ik$$

#### Christoph U. Keller, C.U.Keller@astro-uu.nl

temporal derivatives ⇒ Helmholtz equation

$$abla^2 ec{E} + rac{\omega^2}{c^2} \left(\epsilon + i rac{4\pi\sigma}{\omega}
ight) ec{E} = 0$$

• dispersion relation between  $\vec{k}$  and  $\omega$ 

$$\vec{k} \cdot \vec{k} = \frac{\omega^2}{c^2} \left( \epsilon + i \frac{4\pi\sigma}{\omega} \right)$$

complex index of refraction

$$\tilde{n}^2 = \epsilon + i \frac{4\pi\sigma}{\omega}, \ \vec{k} \cdot \vec{k} = \frac{\omega^2}{c^2} \tilde{n}^2$$

 split into real (n: index of refraction) and imaginary parts (k: extinction coefficient)

$$\tilde{n} = n + ik$$

temporal derivatives ⇒ Helmholtz equation

$$abla^2 ec{E} + rac{\omega^2}{c^2} \left(\epsilon + i rac{4\pi\sigma}{\omega}
ight) ec{E} = 0$$

• dispersion relation between  $\vec{k}$  and  $\omega$ 

$$\vec{k} \cdot \vec{k} = \frac{\omega^2}{c^2} \left( \epsilon + i \frac{4\pi\sigma}{\omega} \right)$$

• complex index of refraction

$$\tilde{n}^2 = \epsilon + i \frac{4\pi\sigma}{\omega}, \ \vec{k} \cdot \vec{k} = \frac{\omega^2}{c^2} \tilde{n}^2$$

 split into real (n: index of refraction) and imaginary parts (k: extinction coefficient)

$$\tilde{n} = n + ik$$

temporal derivatives ⇒ Helmholtz equation

$$abla^2 ec{E} + rac{\omega^2}{c^2} \left(\epsilon + i rac{4\pi\sigma}{\omega}\right) ec{E} = 0$$

• dispersion relation between  $\vec{k}$  and  $\omega$ 

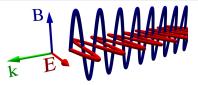
$$\vec{k} \cdot \vec{k} = \frac{\omega^2}{c^2} \left( \epsilon + i \frac{4\pi\sigma}{\omega} \right)$$

• complex index of refraction

$$\tilde{n}^2 = \epsilon + i \frac{4\pi\sigma}{\omega}, \ \vec{k} \cdot \vec{k} = \frac{\omega^2}{c^2} \tilde{n}^2$$

 split into real (n: index of refraction) and imaginary parts (k: extinction coefficient)

$$\tilde{n} = n + ik$$



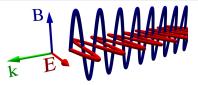
plane-wave solution must also fulfill Maxwell's equations

$$ec{E}_0\cdotec{k}=0,\ ec{H}_0\cdotec{k}=0,\ ec{H}_0=ec{n}rac{ec{k}}{ec{k}ec{k}} imesec{E}_0$$

 isotropic media: electric, magnetic field vectors normal to wave vector ⇒ transverse waves

•  $\vec{E}_0$ ,  $\vec{H}_0$ , and  $\vec{k}$  orthogonal to each other, right-handed vector-triple

- conductive medium  $\Rightarrow$  complex  $\tilde{n}$ ,  $E_0$  and  $H_0$  out of phase
- $E_0$  and  $H_0$  have constant relationship  $\Rightarrow$  consider only E



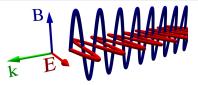
plane-wave solution must also fulfill Maxwell's equations

$$ec{E}_0\cdotec{k}=0,\ ec{H}_0\cdotec{k}=0,\ ec{H}_0=ec{n}rac{ec{k}}{ec{k}ec{k}} imesec{E}_0$$

 isotropic media: electric, magnetic field vectors normal to wave vector ⇒ transverse waves

•  $E_0$ ,  $H_0$ , and k orthogonal to each other, right-handed vector-triple • conductive medium  $\Rightarrow$  complex  $\tilde{n}$ ,  $\vec{E}_0$  and  $\vec{H}_0$  out of phase

•  $E_0$  and  $H_0$  have constant relationship  $\Rightarrow$  consider only E



plane-wave solution must also fulfill Maxwell's equations

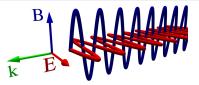
$$ec{E}_0\cdotec{k}=0,\ ec{H}_0\cdotec{k}=0,\ ec{H}_0=ec{n}rac{ec{k}}{ec{k}ec{k}} imesec{E}_0$$

 isotropic media: electric, magnetic field vectors normal to wave vector ⇒ transverse waves

•  $\vec{E}_0$ ,  $\vec{H}_0$ , and  $\vec{k}$  orthogonal to each other, right-handed vector-triple

• conductive medium  $\Rightarrow$  complex  $\tilde{n}$ ,  $E_0$  and  $H_0$  out of phase

•  $E_0$  and  $H_0$  have constant relationship  $\Rightarrow$  consider only E



plane-wave solution must also fulfill Maxwell's equations

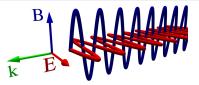
$$ec{E}_0\cdotec{k}=0,\ ec{H}_0\cdotec{k}=0,\ ec{H}_0=ec{n}rac{ec{k}}{ec{k}ec{k}} imesec{E}_0$$

 isotropic media: electric, magnetic field vectors normal to wave vector ⇒ transverse waves

•  $\vec{E}_0$ ,  $\vec{H}_0$ , and  $\vec{k}$  orthogonal to each other, right-handed vector-triple

• conductive medium  $\Rightarrow$  complex  $\tilde{n}$ ,  $\vec{E}_0$  and  $\vec{H}_0$  out of phase

•  $\vec{E}_0$  and  $\vec{H}_0$  have constant relationship  $\Rightarrow$  consider only  $\vec{E}$ 



plane-wave solution must also fulfill Maxwell's equations

$$ec{E}_0 \cdot ec{k} = 0, \ ec{H}_0 \cdot ec{k} = 0, \ ec{H}_0 = ec{n} rac{ec{k}}{ec{k} ec{k}} imes ec{E}_0$$

 isotropic media: electric, magnetic field vectors normal to wave vector ⇒ transverse waves

•  $\vec{E}_0$ ,  $\vec{H}_0$ , and  $\vec{k}$  orthogonal to each other, right-handed vector-triple

- conductive medium  $\Rightarrow$  complex  $\tilde{n}$ ,  $\vec{E}_0$  and  $\vec{H}_0$  out of phase
- $\vec{E}_0$  and  $\vec{H}_0$  have constant relationship  $\Rightarrow$  consider only  $\vec{E}$

Poynting vector

$$ec{S}=rac{c}{4\pi}\left(ec{E} imesec{H}
ight)$$

•  $|\vec{S}|$ : energy through unit area perpendicular to  $\vec{S}$  per unit time

- direction of S is direction of energy flow
- time-averaged Poynting vector given by

$$\left\langle ec{S} \right
angle = rac{c}{8\pi} \mathrm{Re} \left( ec{E}_0 imes ec{H}_0^* 
ight)$$

Re real part of complex expression

- \* complex conjugate
- (.) time average

• energy flow parallel to wave vector (in isotropic media)

$$\left\langle ec{S} \right
angle = rac{c}{8\pi} \left| ec{n} \right| \left| E_0 \right|^2 rac{ec{k}}{\left| ec{k} \right|}$$

Poynting vector

$$ec{S} = rac{m{c}}{m{4}\pi} \left( ec{m{E}} imes ec{m{H}} 
ight)$$

- $|\vec{S}|$ : energy through unit area perpendicular to  $\vec{S}$  per unit time
- direction of  $\vec{S}$  is direction of energy flow
- time-averaged Poynting vector given by

$$\left\langle ec{S} \right
angle = rac{c}{8\pi} \mathrm{Re} \left( ec{E}_0 imes ec{H}_0^* 
ight)$$

- Re real part of complex expression
  - \* complex conjugate
  - (.) time average
- energy flow parallel to wave vector (in isotropic media)

$$\left\langle ec{S} \right
angle = rac{c}{8\pi} \left| ec{n} \right| \left| E_0 \right|^2 rac{ec{k}}{\left| ec{k} \right|}$$

Poynting vector

$$ec{S}=rac{c}{4\pi}\left(ec{E} imesec{H}
ight)$$

- $|\vec{S}|$ : energy through unit area perpendicular to  $\vec{S}$  per unit time • direction of  $\vec{S}$  is direction of energy flow
- time-averaged Poynting vector given by

$$\left\langle ec{S} 
ight
angle = rac{c}{8\pi} {
m Re} \left( ec{E}_0 imes ec{H}_0^st 
ight)$$

- Re real part of complex expression
  - \* complex conjugate
- $\langle . \rangle$  time average

energy flow parallel to wave vector (in isotropic media)

$$\left\langle ec{S} 
ight
angle = rac{c}{8\pi} \left| ec{n} 
ight| \left| E_0 
ight|^2 rac{ec{k}}{\left| ec{k} 
ight|}$$

Poynting vector

# $\vec{S} = rac{c}{4\pi} \left( \vec{E} imes \vec{H} ight)$

- $|\vec{S}|$ : energy through unit area perpendicular to  $\vec{S}$  per unit time • direction of  $\vec{S}$  is direction of energy flow
- time-averaged Poynting vector given by

$$\left< ec{S} \right> = rac{c}{8\pi} {
m Re} \left( ec{E}_0 imes ec{H}_0^* 
ight)$$

Re real part of complex expression

- \* complex conjugate
- $\langle . \rangle \,$  time average

energy flow parallel to wave vector (in isotropic media)

$$\left\langle ec{S} 
ight
angle = rac{c}{8\pi} \left| ec{n} 
ight| \left| E_0 
ight|^2 rac{ec{k}}{\left| ec{k} 
ight|}$$

- monochromatic light: purely theoretical concept
- monochromatic light wave always fully polarized
- real life: light includes range of wavelengths ⇒ quasi-monochromatic light
- quasi-monochromatic: superposition of mutually incoherent monochromatic light beams whose wavelengths vary in narrow range  $\delta\lambda$  around central wavelength  $\lambda_0$

$$\frac{\delta\lambda}{\lambda} \ll 1$$

 measurement of quasi-monochromatic light: integral over measurement time t<sub>m</sub>

• amplitude, phase (slow) functions of time for given spatial location

• *slow*: variations occur on time scales much longer than the mean period of the wave

- monochromatic light: purely theoretical concept
- monochromatic light wave always fully polarized
- real life: light includes range of wavelengths ⇒ quasi-monochromatic light
- quasi-monochromatic: superposition of mutually incoherent monochromatic light beams whose wavelengths vary in narrow range  $\delta\lambda$  around central wavelength  $\lambda_0$

$$\frac{\delta\lambda}{\lambda} \ll 1$$

- measurement of quasi-monochromatic light: integral over measurement time t<sub>m</sub>
- amplitude, phase (slow) functions of time for given spatial location
- *slow*: variations occur on time scales much longer than the mean period of the wave

- monochromatic light: purely theoretical conceptmonochromatic light wave always fully polarized
- real life: light includes range of wavelengths ⇒ quasi-monochromatic light
- quasi-monochromatic: superposition of mutually incoherent monochromatic light beams whose wavelengths vary in narrow range  $\delta\lambda$  around central wavelength  $\lambda_0$

$$\frac{\delta\lambda}{\lambda} \ll 1$$

 measurement of quasi-monochromatic light: integral over measurement time t<sub>m</sub>

• amplitude, phase (slow) functions of time for given spatial location

• *slow*: variations occur on time scales much longer than the mean period of the wave

- monochromatic light: purely theoretical concept
- monochromatic light wave always fully polarized
- real life: light includes range of wavelengths ⇒ quasi-monochromatic light
- quasi-monochromatic: superposition of mutually incoherent monochromatic light beams whose wavelengths vary in narrow range  $\delta\lambda$  around central wavelength  $\lambda_0$

$$\frac{\delta\lambda}{\lambda}\ll$$
 1

- measurement of quasi-monochromatic light: integral over measurement time t<sub>m</sub>
- amplitude, phase (slow) functions of time for given spatial location
- slow: variations occur on time scales much longer than the mean period of the wave

## Polarization of Quasi-Monochromatic Light

 electric field vector for quasi-monochromatic plane wave is sum of electric field vectors of all monochromatic beams

$$\vec{E}(t) = \vec{E}_0(t) e^{i\left(\vec{k}\cdot\vec{x}-\omega t\right)}$$

• can write this way because  $\delta\lambda\ll\lambda_0$ 

measured intensity of quasi-monochromatic beam

$$\left\langle \vec{E}_{x}\vec{E}_{x}^{*}
ight
angle +\left\langle \vec{E}_{y}\vec{E}_{y}^{*}
ight
angle =\lim_{t_{m}\rightarrow\infty}rac{1}{t_{m}}\int_{-t_{m/2}}^{t_{m/2}}\vec{E}_{x}(t)\vec{E}_{x}^{*}(t)+\vec{E}_{y}(t)\vec{E}_{y}^{*}(t)dt$$

 $\langle \cdots \rangle$ : averaging over measurement time  $t_m$ 

- measured intensity independent of time
- quasi-monochromatic: frequency-dependent material properties (e.g. index of refraction) are constant within  $\Delta \lambda$

## Polarization of Quasi-Monochromatic Light

 electric field vector for quasi-monochromatic plane wave is sum of electric field vectors of all monochromatic beams

$$\vec{E}(t) = \vec{E}_0(t) e^{i\left(\vec{k}\cdot\vec{x}-\omega t\right)}$$

- can write this way because  $\delta\lambda \ll \lambda_0$
- measured intensity of quasi-monochromatic beam

$$\left\langle ec{E}_xec{E}_x^*
ight
angle + \left\langle ec{E}_yec{E}_y^*
ight
angle = \lim_{t_m \to \infty} rac{1}{t_m}\int_{-t_m/2}^{t_m/2}ec{E}_x(t)ec{E}_x^*(t) + ec{E}_y(t)ec{E}_y^*(t)dt$$

 $\langle \cdots \rangle$ : averaging over measurement time  $t_m$ 

- measured intensity independent of time
- quasi-monochromatic: frequency-dependent material properties (e.g. index of refraction) are constant within Δλ

## Polychromatic Light or White Light

- wavelength range comparable wavelength  $(\frac{\delta\lambda}{\lambda} \sim 1)$
- incoherent sum of quasi-monochromatic beams that have large variations in wavelength
- cannot write electric field vector in a plane-wave form
- must take into account frequency-dependent material characteristics
- intensity of polychromatic light is given by sum of intensities of constituting quasi-monochromatic beams

## Polychromatic Light or White Light

- wavelength range comparable wavelength  $(\frac{\delta\lambda}{\lambda} \sim 1)$
- incoherent sum of quasi-monochromatic beams that have large variations in wavelength
- cannot write electric field vector in a plane-wave form
- must take into account frequency-dependent material characteristics
- intensity of polychromatic light is given by sum of intensities of constituting quasi-monochromatic beams

#### Polychromatic Light or White Light

- wavelength range comparable wavelength  $(\frac{\delta\lambda}{\lambda} \sim 1)$
- incoherent sum of quasi-monochromatic beams that have large variations in wavelength
- cannot write electric field vector in a plane-wave form
- must take into account frequency-dependent material characteristics
- intensity of polychromatic light is given by sum of intensities of constituting quasi-monochromatic beams

## Polychromatic Light or White Light

- wavelength range comparable wavelength  $(\frac{\delta\lambda}{\lambda} \sim 1)$
- incoherent sum of quasi-monochromatic beams that have large variations in wavelength
- cannot write electric field vector in a plane-wave form
- must take into account frequency-dependent material characteristics
- intensity of polychromatic light is given by sum of intensities of constituting quasi-monochromatic beams

# Electromagnetic Waves Across Interfaces

#### Fields at Interfaces

- classical optics due to interfaces between 2 different media
- from Maxwell's equations in integral form at interface from medium 1 to medium 2

$$\begin{pmatrix} \vec{D}_2 - \vec{D}_1 \end{pmatrix} \cdot \vec{n} = 0 \begin{pmatrix} \vec{B}_2 - \vec{B}_1 \end{pmatrix} \cdot \vec{n} = 0 \begin{pmatrix} \vec{E}_2 - \vec{E}_1 \end{pmatrix} \times \vec{n} = 0 \begin{pmatrix} \vec{H}_2 - \vec{H}_1 \end{pmatrix} \times \vec{n} = 0$$

normal on interface, points from medium 1 to medium 2
 normal components of *D* and *B* are continuous across interface
 tangential components of *E* and *H* are continuous across interface

# Electromagnetic Waves Across Interfaces

#### Fields at Interfaces

- classical optics due to interfaces between 2 different media
- from Maxwell's equations in integral form at interface from medium 1 to medium 2

$$\begin{pmatrix} \vec{D}_2 - \vec{D}_1 \end{pmatrix} \cdot \vec{n} = 0 \begin{pmatrix} \vec{B}_2 - \vec{B}_1 \end{pmatrix} \cdot \vec{n} = 0 \begin{pmatrix} \vec{E}_2 - \vec{E}_1 \end{pmatrix} \times \vec{n} = 0 \begin{pmatrix} \vec{H}_2 - \vec{H}_1 \end{pmatrix} \times \vec{n} = 0$$

normal on interface, points from medium 1 to medium 2
 normal components of D
 and B
 are continuous across interface

 tangential components of E
 and H
 are continuous across interface

# Electromagnetic Waves Across Interfaces

#### Fields at Interfaces

- classical optics due to interfaces between 2 different media
- from Maxwell's equations in integral form at interface from medium 1 to medium 2

$$\begin{pmatrix} \vec{D}_2 - \vec{D}_1 \end{pmatrix} \cdot \vec{n} = 0 \begin{pmatrix} \vec{B}_2 - \vec{B}_1 \end{pmatrix} \cdot \vec{n} = 0 \begin{pmatrix} \vec{E}_2 - \vec{E}_1 \end{pmatrix} \times \vec{n} = 0 \begin{pmatrix} \vec{H}_2 - \vec{H}_1 \end{pmatrix} \times \vec{n} = 0$$

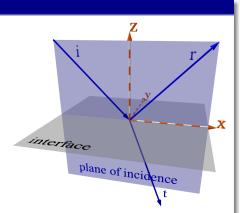
- $\vec{n}$  normal on interface, points from medium 1 to medium 2
- normal components of  $\vec{D}$  and  $\vec{B}$  are continuous across interface
- tangential components of  $\vec{E}$  and  $\vec{H}$  are continuous across interface

#### Plane of Incidence

- plane wave onto interface
- incident (<sup>i</sup>), reflected (<sup>r</sup>), and transmitted (<sup>t</sup>) waves

$$\vec{E}^{i,r,t} = \vec{E}_0^{i,r,t} e^{i(\vec{k}^{i,r,t}\cdot\vec{x}-\omega t)}$$
$$\vec{H}^{i,r,t} = \frac{c}{\omega}\vec{k}^{i,r,t}\times\vec{E}^{i,r,t}$$

• interface normal  $\vec{n} \parallel z$ -axis



spatial, temporal behavior at interface the same for all 3 waves

$$(\vec{k}^i \cdot \vec{x})_{z=0} = (\vec{k}^r \cdot \vec{x})_{z=0} = (\vec{k}^t \cdot \vec{x})_{z=0}$$

• valid for all  $\vec{x}$  in interface  $\Rightarrow$  all 3 wave vectors in one plane, *plane* of incidence

#### Plane of Incidence

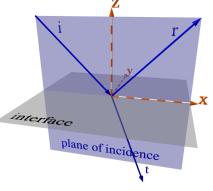
- plane wave onto interface
- incident (<sup>i</sup>), reflected (<sup>r</sup>), and transmitted (<sup>t</sup>) waves

$$\vec{E}^{i,r,t} = \vec{E}_0^{i,r,t} e^{i(\vec{k}^{i,r,t}\cdot\vec{x}-\omega t)}$$
$$\vec{H}^{i,r,t} = \frac{c}{\omega}\vec{k}^{i,r,t} \times \vec{E}^{i,r,t}$$

- interface normal  $\vec{n} \parallel z$ -axis
- spatial, temporal behavior at interface the same for all 3 waves

$$(ec{k}^i \cdot ec{x})_{z=0} = (ec{k}^r \cdot ec{x})_{z=0} = (ec{k}^t \cdot ec{x})_{z=0}$$

• valid for all  $\vec{x}$  in interface  $\Rightarrow$  all 3 wave vectors in one plane, *plane* of incidence



#### Plane of Incidence

- plane wave onto interface
- incident (<sup>i</sup>), reflected (<sup>r</sup>), and transmitted (<sup>t</sup>) waves

$$\vec{E}^{i,r,t} = \vec{E}_0^{i,r,t} e^{i(\vec{k}^{i,r,t}\cdot\vec{x}-\omega t)}$$
$$\vec{H}^{i,r,t} = \frac{c}{\omega}\vec{k}^{i,r,t} \times \vec{E}^{i,r,t}$$

- interface normal  $\vec{n} \parallel z$ -axis
- spatial, temporal behavior at interface the same for all 3 waves

$$(ec{k}^i \cdot ec{x})_{z=0} = (ec{k}^r \cdot ec{x})_{z=0} = (ec{k}^t \cdot ec{x})_{z=0}$$

• valid for all  $\vec{x}$  in interface  $\Rightarrow$  all 3 wave vectors in one plane, *plane* of incidence

interface

plane of incidence

## Snell's Law

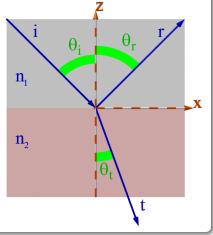
 spatial, temporal behavior the same for all three waves

$$(\vec{k}^i \cdot \vec{x})_{z=0} = (\vec{k}^r \cdot \vec{x})_{z=0} = (\vec{k}^t \cdot \vec{x})_{z=0}$$

• 
$$\left| \vec{k} \right| = \frac{\omega}{c} \tilde{n}$$

ω, c the same for all 3 waves
Snell's law

$$\tilde{n}_1 \sin \theta_i = \tilde{n}_1 \sin \theta_r = \tilde{n}_2 \sin \theta_t$$



## Snell's Law

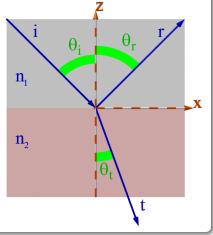
• spatial, temporal behavior the same for all three waves

$$(\vec{k}^i \cdot \vec{x})_{z=0} = (\vec{k}^r \cdot \vec{x})_{z=0} = (\vec{k}^t \cdot \vec{x})_{z=0}$$

• 
$$\left|\vec{k}\right| = \frac{\omega}{c}\tilde{n}$$

ω, c the same for all 3 waves
 Spall's law

$$\tilde{n}_1 \sin \theta_i = \tilde{n}_1 \sin \theta_r = \tilde{n}_2 \sin \theta_t$$



## Snell's Law

• spatial, temporal behavior the same for all three waves

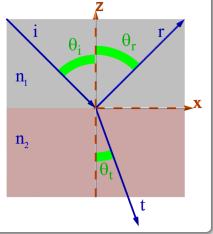
$$(\vec{k}^i \cdot \vec{x})_{z=0} = (\vec{k}^r \cdot \vec{x})_{z=0} = (\vec{k}^t \cdot \vec{x})_{z=0}$$

• 
$$\left|\vec{k}\right| = \frac{\omega}{c}\tilde{n}$$

•  $\omega$ , *c* the same for all 3 waves

Snell's law

$$\tilde{n}_1 \sin \theta_i = \tilde{n}_1 \sin \theta_r = \tilde{n}_2 \sin \theta_t$$



#### Monochromatic Wave at Interface

# $\vec{H}_0^{i,r,t} = \frac{c}{\omega} \vec{k}^{i,r,t} \times \vec{E}_0^{i,r,t}, \quad \vec{B}_0^{i,r,t} = \frac{c}{\omega} \vec{k}^{i,r,t} \times \vec{E}_0^{i,r,t}$

• boundary conditions for monochromatic plane wave:

• 4 equations are not independent

• only need to consider last two equations (tangential components of  $\vec{E}_0$  and  $\vec{H}_0$  are continuous)

#### Monochromatic Wave at Interface

 $\vec{H}_0^{i,r,t} = \frac{c}{\omega} \vec{k}^{i,r,t} \times \vec{E}_0^{i,r,t}, \quad \vec{B}_0^{i,r,t} = \frac{c}{\omega} \vec{k}^{i,r,t} \times \vec{E}_0^{i,r,t}$ 

• boundary conditions for monochromatic plane wave:

$$\begin{pmatrix} \tilde{n}_1^2 \vec{E}_0^i + \tilde{n}_1^2 \vec{E}_0^r - \tilde{n}_2^2 \vec{E}_0^t \end{pmatrix} \cdot \vec{n} = 0 \\ \begin{pmatrix} \vec{k}^i \times \vec{E}_0^i + \vec{k}^r \times \vec{E}_0^r - \vec{k}^t \times \vec{E}_0^t \end{pmatrix} \cdot \vec{n} = 0 \\ \begin{pmatrix} \vec{E}_0^i + \vec{E}_0^r - \vec{E}_0^t \end{pmatrix} \times \vec{n} = 0 \\ \begin{pmatrix} \vec{k}^i \times \vec{E}_0^i + \vec{k}^r \times \vec{E}_0^r - \vec{k}^t \times \vec{E}_0^t \end{pmatrix} \times \vec{n} = 0$$

4 equations are not independent

• only need to consider last two equations (tangential components of  $\vec{E}_0$  and  $\vec{H}_0$  are continuous)

#### Monochromatic Wave at Interface

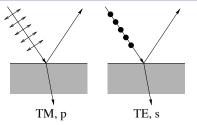
 $\vec{H}_0^{i,r,t} = \frac{c}{\omega} \vec{k}^{i,r,t} \times \vec{E}_0^{i,r,t}, \quad \vec{B}_0^{i,r,t} = \frac{c}{\omega} \vec{k}^{i,r,t} \times \vec{E}_0^{i,r,t}$ 

• boundary conditions for monochromatic plane wave:

$$\left(\vec{E}_0^i + \vec{E}_0^r - \vec{E}_0^t\right) \times \vec{n} = 0$$
$$\left(\vec{k}^i \times \vec{E}_0^i + \vec{k}^r \times \vec{E}_0^r - \vec{k}^t \times \vec{E}_0^t\right) \times \vec{n} = 0$$

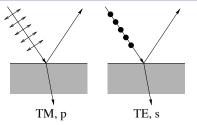
- 4 equations are not independent
- only need to consider last two equations (tangential components of *E*<sub>0</sub> and *H*<sub>0</sub> are continuous)

#### **Two Special Cases**



- electric field parallel to plane of incidence ⇒ magnetic field is transverse to plane of incidence (TM)
- electric field particular (German: senkrecht) or transverse to plane of incidence (TE)
  - general solution as (coherent) superposition of two cases
    choose direction of magnetic field vector such that Poynting vector parallel, same direction as corresponding wave vector

### **Two Special Cases**



- electric field parallel to plane of incidence ⇒ magnetic field is transverse to plane of incidence (TM)
- electric field particular (German: senkrecht) or transverse to plane of incidence (TE)
  - general solution as (coherent) superposition of two cases
  - choose direction of magnetic field vector such that Poynting vector parallel, same direction as corresponding wave vector