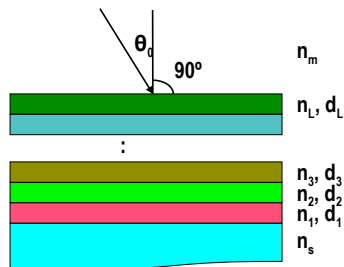


Outline

- 1 Thin Films
- 2 Calculating Thin Film Stack Properties
- 3 Fabry-Perot Tunable Filter
- 4 Anti-Reflection Coatings
- 5 Interference Filters

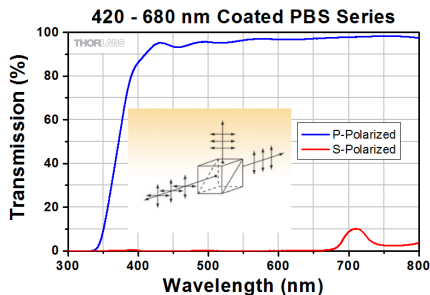
Introduction

- *thin film*:
 - layer with thickness $\lesssim \lambda$
 - extends in 2 other dimensions $\gg \lambda$
- reflection, refraction at all interfaces
- L layers of thin films: *thin film stack*
- layer thickness $d_i \lesssim \lambda \Rightarrow$ interference between reflected and refracted waves
- infinite number of multiple reflections and refractions must be considered
- assumption: *substrate* (index n_s), incident medium (index n_m) have infinite thickness



Thin Films and Polarization

- thin film coatings and polarization are closely coupled
- some polarizers (plate, cube) are based on thin-film coatings
- can dramatically reduce polarizing effects of optical components
- all aluminum mirrors have aluminum oxide thin film



www.thorlabs.com/newgroupage9.cfm?objectgroup_id=739

Calculating Thin-Film Stack Properties

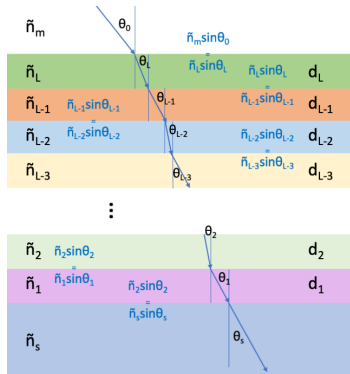
- many layers \Rightarrow consider all interferences between reflected and refracted rays of each interface
- complexity of calculation significantly reduced by matrix approach
- signs and conventions are not consistent in literature

Plane Waves and Thin-Film Stacks

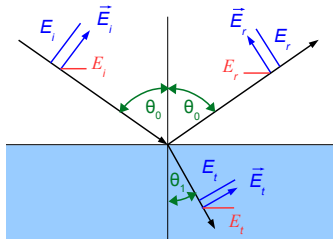
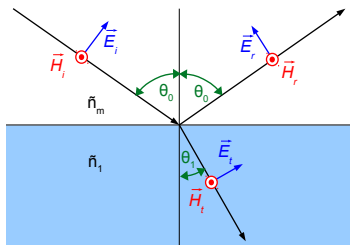
- plane wave $\vec{E} = \vec{E}_0 e^{i(\vec{k} \cdot \vec{x} - \omega t)}$
- layers numbered from 1 to L with complex index of refraction \tilde{n}_j , geometrical thickness d_j
- substrate has refractive index \tilde{n}_s
- incident medium has index \tilde{n}_m
- angle of incidence in incident medium: θ_0
- Snell's law:

$$\tilde{n}_m \sin \theta_0 = \tilde{n}_L \sin \theta_L = \dots = \tilde{n}_1 \sin \theta_1 = \tilde{n}_s \sin \theta_s$$

- $\Rightarrow \theta_j$ for every layer j



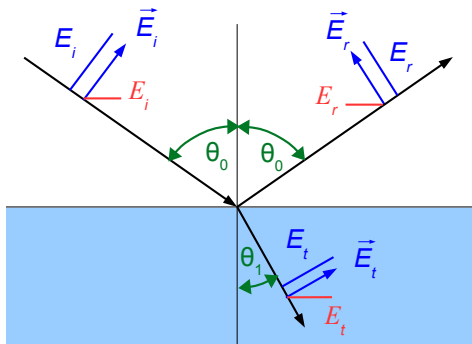
p/TM Wave: Electric Field at Interface



- p/TM wave: electric fields are in (parallel to) plane of incidence
- $E_{i,r,t}$: complex amplitudes of $\vec{E}_{i,r,t}$ of incident, reflected, transmitted electric fields
- $E_{i,r,t}$: components tangential to interface of complex amplitudes of $\vec{E}_{i,r,t}$ of incident, reflected, transmitted electric fields

$$E_i = E_i \cos \theta_0, \quad E_r = E_r \cos \theta_0, \quad E_t = E_t \cos \theta_1$$

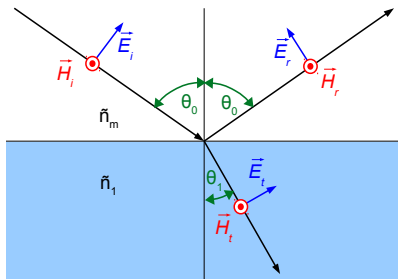
p/TM Wave: Electric Field at Interface (continued)



- continuous electric field tangential to interface:

$$E_i \cos \theta_0 - E_r \cos \theta_0 = E_t \cos \theta_1$$

p/TM Wave: Electric Field at Interface (continued)

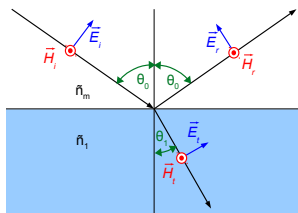


- thin-film formalism uses field components parallel to interface
- continuous electric field tangential to interface

$$E_i - E_r = E_t$$

- electric field direction fully determined by angle of incidence
- complex scalar quantities are sufficient since the electric field is perpendicular to wave vector and in plane of incidence

p/TM Wave: Magnetic Field at Interface

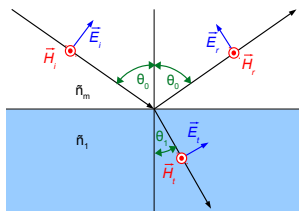


- non-magnetic material ($\mu = 1$): $\vec{H} = \tilde{n} \frac{\vec{k}}{|\vec{k}|} \times \vec{E}$
- \vec{k} , \vec{E} perpendicular, cross-product becomes multiplication
- (complex) magnitudes of $\vec{E}_{i,r,t}$ and $\vec{H}_{i,r,t}$ related by

$$H_{i,r} = \tilde{n}_m E_{i,r}, \quad H_t = \tilde{n}_1 E_t$$

- tangential component of \vec{H} continuous: $\mathbf{H}_i + \mathbf{H}_r = \mathbf{H}_t$
- $\vec{H}_{i,r,t}$ is already tangential to interface: $\vec{H}_{i,r,t} = \vec{H}_{i,r,t}$

p/TM Wave: Magnetic Field at Interface (continued)



- tangential component of \vec{H} continuous: $H_i + H_r = H_t$
- expressed in electric field

$$\tilde{n}_m E_i + \tilde{n}_m E_r = \tilde{n}_1 E_t$$

- expressed in terms of tangential electric field components

$$\frac{\tilde{n}_m}{\cos \theta_0} \mathbf{E}_i + \frac{\tilde{n}_m}{\cos \theta_0} \mathbf{E}_r = \frac{\tilde{n}_1}{\cos \theta_1} \mathbf{E}_t$$

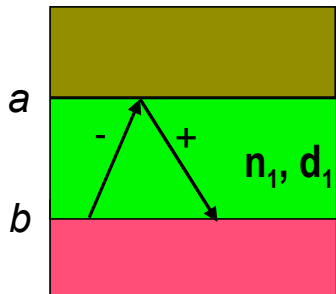
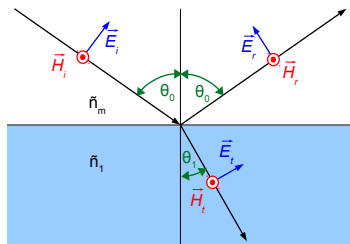
Matrix Formalism: Tangential Components in one Medium

- single interface in thin-film stack, combine all waves into
 - wave that travels towards substrate (+ superscript)
 - wave that travels away from substrate (– superscript)
- at interface a , *tangential components* of complex electric and magnetic field amplitudes in medium 1 given by

$$E_a = E_{1a}^+ - E_{1a}^-$$

$$H_a = \frac{\tilde{n}_1}{\cos \theta_1} (E_{1a}^+ + E_{1a}^-)$$

- negative sign for outwards traveling electric field component

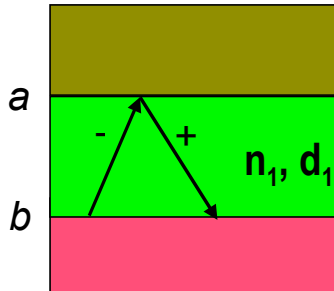


Matrix Formalism: Electric Field Propagation

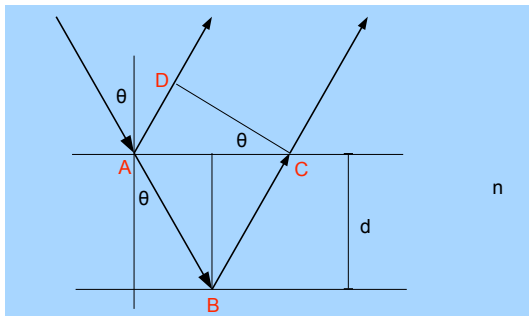
- field amplitudes in medium 1 at (other) interface b from wave propagation with common phase factor
- d_1 : geometrical thickness of layer
- phase factor for forward propagating wave:

$$\delta = \frac{2\pi}{\lambda} \tilde{n}_1 d_1 \cos \theta_1$$

- backwards propagating wave: same phase factor with negative sign



Plane Wave Path Length for Oblique Incidence



- consider theoretical reflections in single medium
- need to correct for plane wave propagation
- path length for “reflected light”: $\bar{AB} + \bar{BC} - \bar{AD}$

$$\frac{2d}{\cos \theta} - 2d \tan \theta \cdot \sin \theta = 2d \frac{1 - \sin^2 \theta}{\cos \theta} = 2d \cos \theta$$

Matrix Formalism

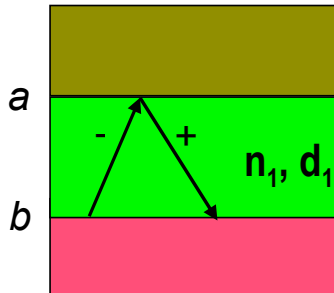
- at interface b in medium 1

$$E_{1b}^+ = E_{1a}^+(\cos \delta + i \sin \delta)$$

$$E_{1b}^- = E_{1a}^-(\cos \delta - i \sin \delta)$$

$$H_{1b}^+ = H_{1a}^+(\cos \delta + i \sin \delta)$$

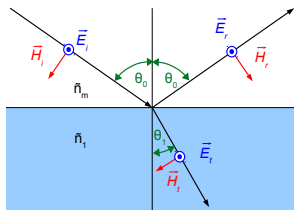
$$H_{1b}^- = H_{1a}^-(\cos \delta - i \sin \delta)$$



- from before $E_{a,b} = E_{1a,b}^+ - E_{1a,b}^-$, $H_{a,b} = \frac{\tilde{n}_1}{\cos \theta_1} (E_{1a,b}^+ + E_{1a,b}^-)$
- propagation of tangential components from a to b

$$\begin{pmatrix} E_b \\ H_b \end{pmatrix} = \begin{pmatrix} \cos \delta & \frac{\cos \theta_1}{\tilde{n}_1} i \sin \delta \\ i \frac{\tilde{n}_1}{\cos \theta_1} \sin \delta & \cos \delta \end{pmatrix} \begin{pmatrix} E_a \\ H_a \end{pmatrix}$$

s/TE waves: Electric and Magnetic Fields at Interface



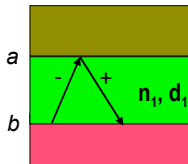
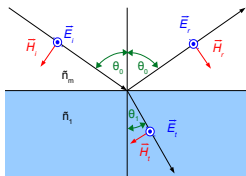
- tangential electric field component continuous: $E_i + E_r = E_t$
- tangential magnetic field component continuous

$$H_i \cos \theta_0 - H_r \cos \theta_0 = H_t \cos \theta_1$$

- and using relation between H and E

$$\tilde{n}_m \cos \theta_0 E_i - \tilde{n}_m \cos \theta_0 E_r = \tilde{n}_1 \cos \theta_1 E_t$$

s-Polarized Waves in one Medium



- at boundary a:

$$E_a = E_{1a}^+ + E_{1a}^-$$

$$H_a = (E_{1a}^+ - E_{1a}^-) \tilde{n}_1 \cos \theta_1$$

- propagation of tangential components from a to b

$$\begin{pmatrix} E_b \\ H_b \end{pmatrix} = \begin{pmatrix} \cos \delta & \frac{1}{\tilde{n}_1 \cos \theta_1} i \sin \delta \\ i \tilde{n}_1 \cos \theta_1 \sin \delta & \cos \delta \end{pmatrix} \begin{pmatrix} E_a \\ H_a \end{pmatrix}$$

Summary of Matrix Method

for each layer j calculate:

- θ_j using Snell's law: $n_m \sin \theta_0 = n_j \sin \theta_j$
- s-polarization: $\eta_j = n_j \cos \theta_j$
- p-polarization: $\eta_j = \frac{n_j}{\cos \theta_j}$
- phase delays: $\delta_j = \frac{2\pi}{\lambda} n_j d_j \cos \theta_j$
- characteristic matrix:

$$M_j = \begin{pmatrix} \cos \delta_j & \frac{i}{\eta_j} \sin \delta_j \\ i\eta_j \sin \delta_j & \cos \delta_j \end{pmatrix}$$

Summary of Matrix Method (continued)

- total characteristic matrix M is product of all characteristic matrices

$$M = M_L M_{L-1} \dots M_2 M_1$$

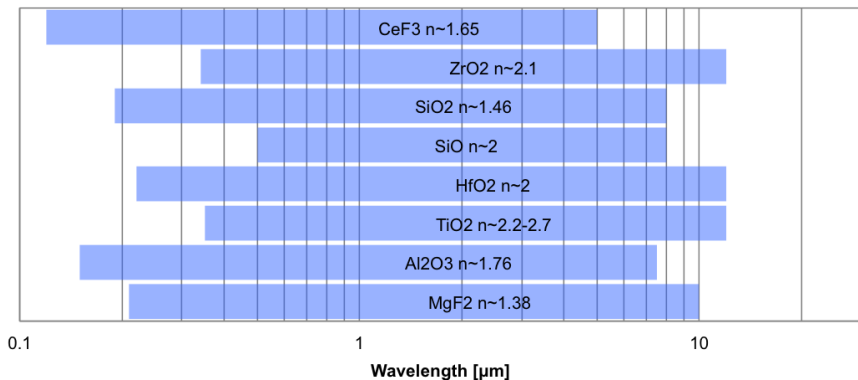
- fields in incident medium given by

$$\begin{pmatrix} E_m \\ H_m \end{pmatrix} = M \begin{pmatrix} 1 \\ \eta_s \end{pmatrix}$$

- complex reflection and transmission coefficients

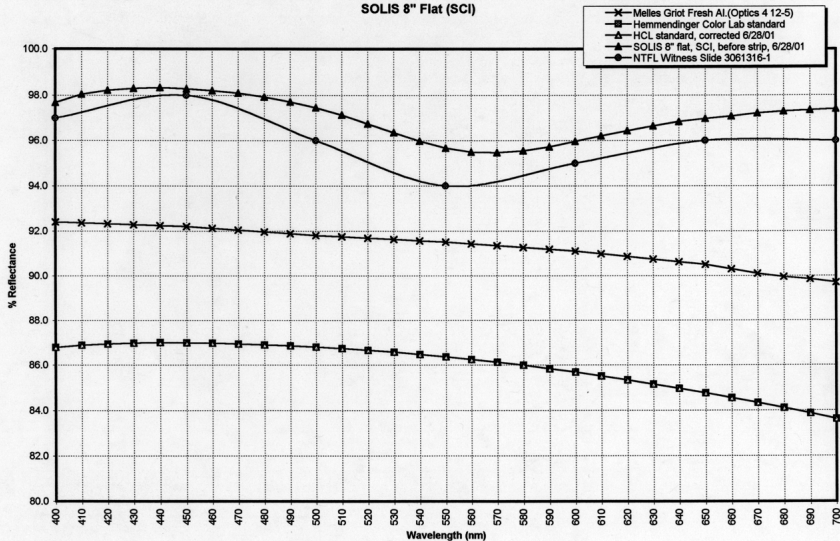
$$r = \frac{\eta_m E_m - H_m}{\eta_m E_m + H_m}, \quad t = \frac{2\eta_m}{\eta_m E_m + H_m}$$

Materials and Refractive Indices

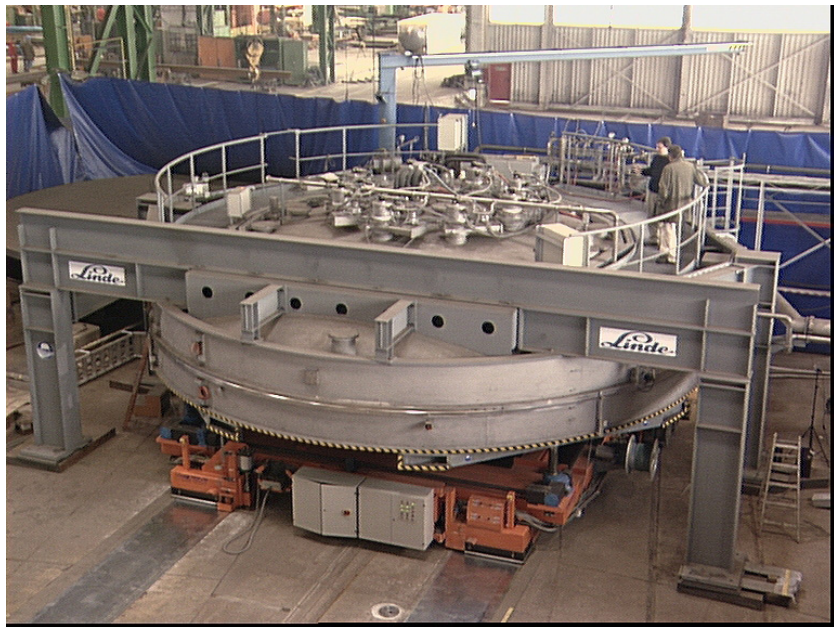


Mirror Coatings

SOLIS 8" Flat (SCI)



VLT Coating Chamber



VLT Coating Chamber with Magnetogron



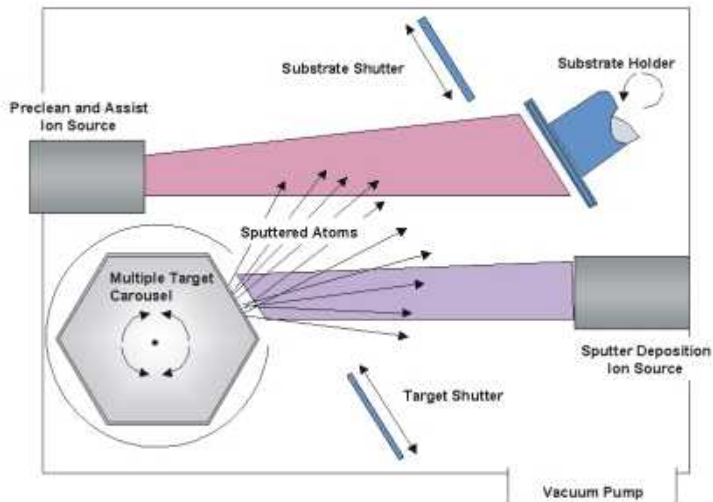
Evaporation

- evaporation of solid material through high electric current (e.g. classic Al mirror coatings)
- sputtering: plasma glow discharge ejects material from solid substance

Deposition

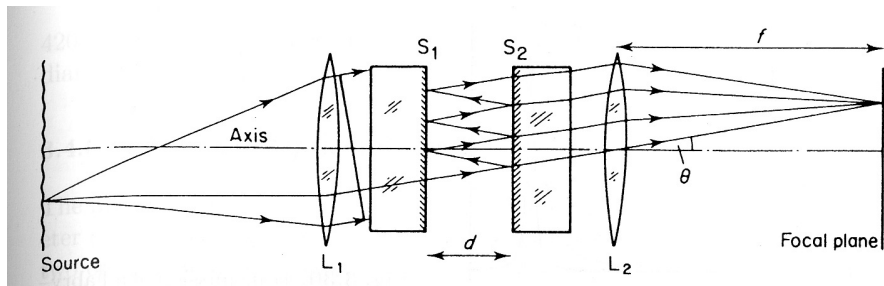
- uncontrolled ballistic flights, mechanical shields to homogenize beam
- ion-assisted deposition

Ion-Beam Deposition



www.4waveinc.com/ibd.html

Tunable Filter



- invented by Fabry and Perot in 1899
- interference between partially transmitting plates containing medium with index of refraction n
- angle of incidence in material θ , distance d
- path difference between successive beams: $\Delta = 2nd \cos \theta$

Fabry Perot continued

- path difference between successive beams: $\Delta = 2nd \cos \theta$
- phase difference: $\delta = 2\pi\Delta/\lambda = 4\pi nd \cos \theta/\lambda$
- incoming wave: $e^{i\omega t}$
- intensity transmission at surface: T
- intensity reflectivity at surface: R
- outgoing wave is the *coherent sum* of all beams

$$Ae^{i\omega t} = Te^{i\omega t} + TRe^{i(\omega t + \delta)} + TR^2e^{i(\omega t + 2\delta)} + \dots$$

- write this as

$$A = T(1 + Re^{i\delta} + R^2e^{i2\delta} + \dots) = \frac{T}{1 - Re^{i\delta}}$$

Fabry Perot continued

- emerging amplitude

$$A = \frac{T}{1 - Re^{i\delta}}$$

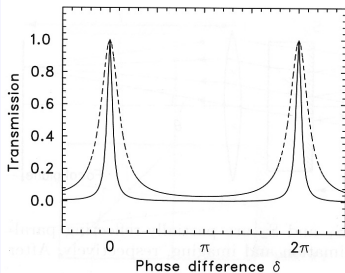
- emerging intensity is therefore

$$I = AA^* = \frac{T^2}{1 - 2R \cos \delta + R^2} = \frac{T^2}{(1 - R)^2 + 4R \sin^2 \frac{\delta}{2}}$$

- with $I_{\max} = T^2 / (1 - R)^2$

$$I = I_{\max} \frac{1}{1 + \frac{4R}{(1-R)^2} \sin^2 \frac{\delta}{2}}$$

Fabry Perot Properties



- transmission is periodic
- distance between transmission peaks, *free spectral range*

$$\text{FSR} = \frac{\lambda_0^2}{2nd \cos \theta}$$

- Full-Width at Half Maximum (FWHM) $\Delta\lambda$

$$\Delta\lambda = \text{FSR}/F$$

- *finesse* F

$$F = \frac{\pi\sqrt{R}}{1-R}$$

$$I = I_{\max} \frac{1}{1 + \frac{4R}{(1-R)^2} \sin^2 \frac{\delta}{2}}$$

Reflection Off Uncoated Substrate

- reflectivity of bare substrate:

$$R = \left(\frac{n_m - n_s}{n_m + n_s} \right)^2$$

- fused silica at 600 nm: $n_s = 1.46 \Rightarrow R = 0.035$
- extra-dense flint glass at 600 nm: $n_s = 1.75 \Rightarrow R = 0.074$
- silicon at 600 nm: $n_s = 3.96 \Rightarrow R = 0.6$
- loss in transparency
- ghost reflections

Analytical treatment of single layer

- assume single-layer coating
- determine optimum thickness and index of material
- amplitude reflection from coating/air interface

$$A_1 = \frac{n_1 - 1}{n_1 + 1}$$

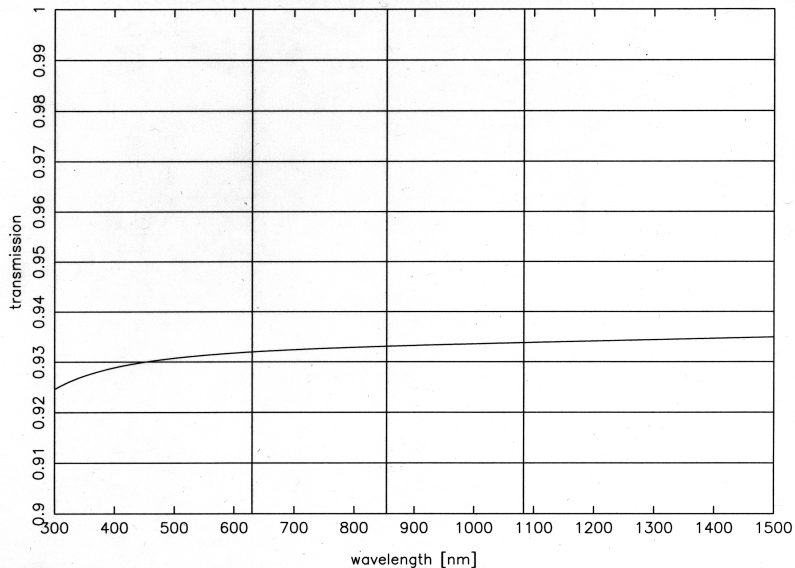
- amplitude reflection from coating/substrate interface

$$A_1 = \frac{n_s - n_1}{n_s + n_1}$$

- amplitudes subtract for 180 degree phase difference \Rightarrow coating should have optical path length $\lambda/4$ thick
- best cancellation for $n_1 = \sqrt{n_s}$

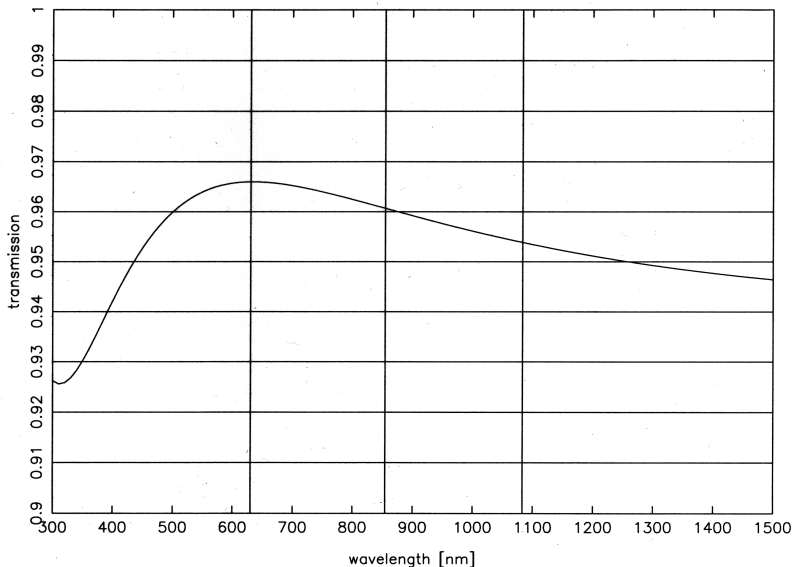
Uncoated Fused Silica

FS substrate, no coating



MgF₂ Coating on Fused Silica

FS substrate, 114.8 nm MgF₂



Multi-Layer Coatings on Fused Silica

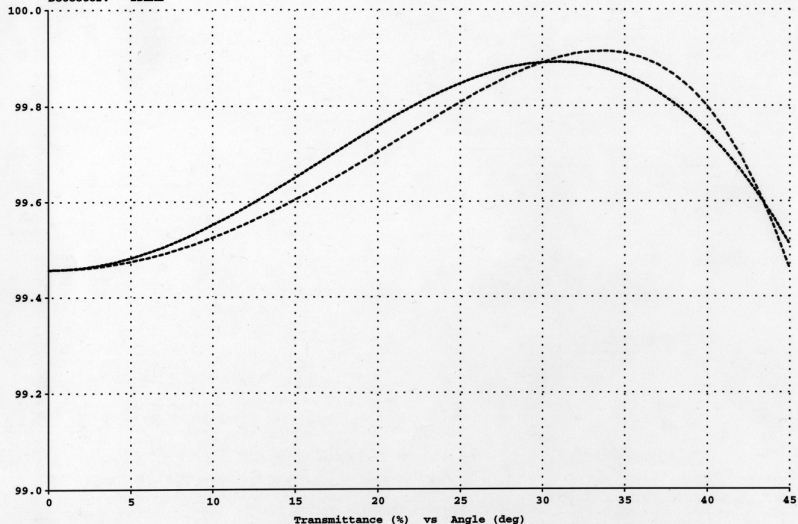
Software Spectra, Inc.

corrlens

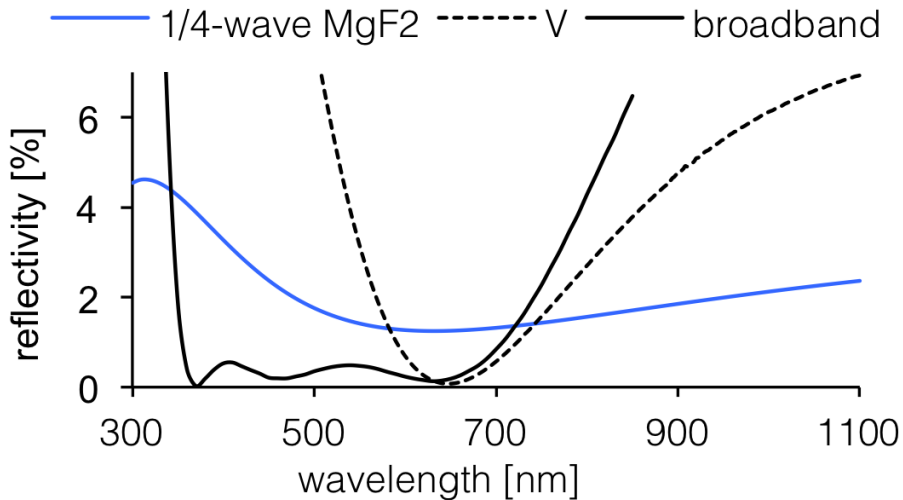
1/20/99 3:20:04 PM

Illuminant: WHITE
Medium: AIR
Substrate: GLASS
Exit: GLASS
Detector: IDEAL

Wavelength: 630.2 (nm)
Reference: 632.8 (nm)
Polarization: S --- P —

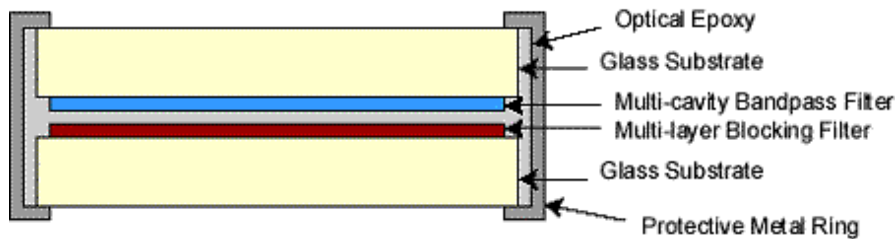


AR Coating Types



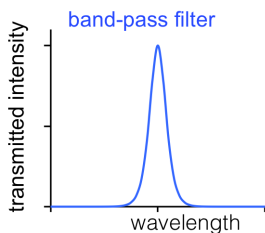
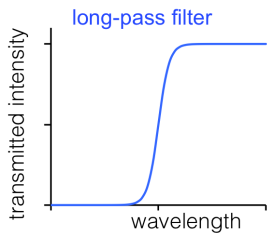
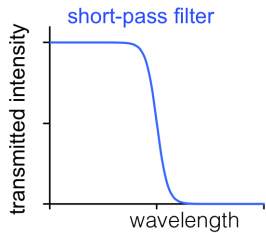
Overview

Complete Interference Filter

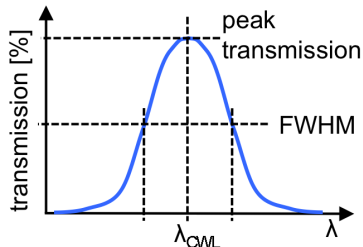


- many layers can achieve many things
- band-pass, long-pass, short-pass, dichroic filters
- colored glass substrates often used in addition to coatings
- sensitive to angle of incidence
- evaporated coatings are very temperature sensitive

Types of Interference Filters



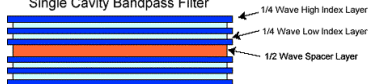
Bandpass Interference Filter Terminology



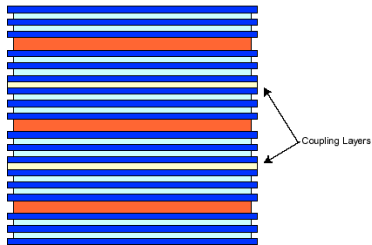
- **Bandpass:** range of wavelengths passed by filter
- **Blocking:** degree of attenuation outside of passband
- **Center Wavelength (CWL):** wavelength at midpoint of passband FWHM
- **Number of cavities:** Fabry-Perot-like thin-film arrangement
- **Full-width Half-Maximum (FWHM):** width of bandpass at one-half of peak transmission
- **Peak Transmission in %:** maximum transmission in passband

Cavities

Single Cavity Bandpass Filter



Three Cavity Bandpass Filter



- basic design element: Fabry-Perot cavity
- cavity: $\lambda/2$ spacer and reflective multi-layer coatings on either side
- stacks of cavities provide much better performance

www.fluorescence.com/tutorial/int-filt.htm

Multiple Cavity Performance

