# Spectrographs

ATI 2016 - Lecture 11 Matthew Kenworthy // Leiden Observatory

### **The Solar Spectrum**



# **Design drivers for spectrographs**

What spectral resolution do you need?

Spectral resolution 
$$R = \frac{\lambda}{\Delta \lambda}$$

What bandwidth (wavelength range) do you need?

Spectrograph is sensitive from 
$$\lambda_{blue}~$$
 to  $~\lambda_{red}~$ 

Maximising throughput for best efficiency

Etendue, limiting magnitude, throughput, multiplexing

# Science drivers for spectrographs



# **Basic spectroscopy: colour filters**



# Take multiple images with different bandpass filters

Johnson system designed to measure properties of stars

Thuan-Gunn filters for faint galaxy observations

Stromgren has better sensitivity to stellar properties (metallicity, temperature, surface gravity)

Sloan Digital Sky Survey (SDSS) for faint galaxy classification

### **Basic spectroscopy: colour filters**



**UBVRI Filter Characteristics** 





VRIJKLMNQ by Johnson (1960)

**UBV** by Johnson and Morgan (1953)

**Classifying stars with photomultipliers** 

Zero points of (B-V) and (U-B) color indices defined to be zero for A0 V stars

### **Slitless spectrographs**

#### Put a dispersing element in front of the telescope aperture



http://www.lpl.arizona.edu/~rhill/

# Slitless spectrographs



Slitless K-band Spectrum of Saturn

- 2.0µ

- 2.3µ

# **Slitless spectrographs**



Dispersed

R. Pogge (OSU) with NOAO 2.1m Telescope

#### **Slitless spectrographs**

The solar corona (solar disk is blocked by a coronagraph)



Wavelength

http://www.astro.virginia.edu/class/majewski/astr313/lectures/spectroscopy/spectrographs.html

# Layout of a spectrograph



**IMPORTANT!**  $d_{coll}$  and  $d_{cam}$  may not be the same!

# Layout of a spectrograph



#### Anamorphic magnification

#### **Resolution Element**

The resolution element is the minimum resolution of the spectrograph. This will depend of the spectral size of the image, which is a factor of image size, spectral magnification and the linear dispersion

Typically the central wavelength

$$R = \frac{\lambda}{\Delta \lambda}$$
 Resolution element

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# The Slit

We cannot record three dimensions of data (x,y, wavelength) onto a two dimensional detector, so we need to choose how we fill up our detector area:



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## Setting the slit width

For a seeing limited object, such as a star, varying the slit width is a balance between spectral resolution and throughput



Slit too wide, spectral resolution goes down

Slit too narrow, flux from seeing limited object is lost

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# The Slit

Spectrographic slits are given in terms of their angular size on the sky, either in arc seconds or in radians.

$$\phi = w/f$$

where f is the focal length of the telescope and w is the size of the slit in mm. The angle  $\phi$  is given in radians.



### Two types of magnification

Anamorphic magnification arises because the diffracting element may send light off at a large angle from the camera normal, and is defined as r.

$$r = \frac{d_{coll}}{d_{cam}} = \frac{d\beta}{d\alpha}$$

$$w' = rw\frac{f_{cam}}{f_{coll}}$$

Spatial (de)magnification occurs because of the different focal lengths of the camera and collimator so that detector pixels are Nyquist sampled

#### Two types of magnification

The size of the slit that the detector sees for the slit is therefore given by:

$$w' = rw\frac{f_{cam}}{f_{coll}} = r\phi f \frac{f_{cam}}{f_{coll}}$$

# **Definition of Dispersion**



#### **Dispersion of Glass Prisms**

Prisms are used near minimum deviations so that rays inside the prism are parallel to the base. The input and output beams are the same size.



$$A = \frac{d\beta}{d\lambda} = \frac{B}{d_{cam}} \frac{dn}{d\lambda}$$

Angular dispersion changes with wavelength For k identical prisms in a row, dispersion is multiplied by k

### **Dispersion of Glass Prisms**

Dispersion is not constant with wavelength, and very high resolution is not possible.



Can be transmissive or reflective, and consist of thousands of periodic features on an optically flat surface.

Manufactured using ruling engines in temperature controlled rooms



Made by David Rittenhouse in 1785 Reinvented by Frauenhofer in 1821





Frauenhofer gratings resolved Solar absorption spectrum, and labelled the absorption lines with letters (A,B,C,D...)

#### **HARPS** grating





Direction of constructive interference is wavelength dependent

# **Dispersion of Diffraction Gratings**

From diffraction theory, the grating equation relates the order m, the groove spacing  $\sigma$  (the number of mm between each ruled line)

$$m\lambda = \sigma(\sin\alpha \pm \sin\beta)$$
  
... where the sign is positive for reflection,  
negative for transmission  
Angular dispersion 
$$A = \frac{d\beta}{d\alpha} = \frac{m}{\sigma\cos\beta}$$

Typically 600 lines per mm and used at 60 degrees incidence

#### **Increasing spectral resolution**

Increasing  $\sigma$  is difficult, and  $\cos\beta$  cannot be greater than unity

Angular dispersion 
$$A = \frac{d\beta}{d\alpha} = \frac{m}{\sigma \cos \beta}$$

Look at large values of  $\,m\,$  to get high spectral resolution

$$R = nm$$

where n is the total number of illuminated grooves

# **Higher spectral orders**

Higher order dispersion from the grating will result in overlapping spectra:



Figure 4.1.10. A portion of the image structure for a single bichromatic point source viewed through several apertures.

The free spectral range of a spectrograph is given by:

$$\lambda' - \lambda = \lambda/m$$
$$m\lambda' = (m+1)\lambda$$

We can either use an ORDER BLOCKING FILTER or a CROSS disperser to split out the different spectral orders

# **Higher spectral orders**

**CROSS** disperser to split out the different spectral orders



# **Higher spectral orders**

#### **CROSS** disperser to split out the different spectral orders



Trispec

# **Diffraction grating efficiency**

150 /mm



# **Optimising the grating efficiency**



$$\theta_B = \frac{\alpha + \beta}{2} \lambda_B = \frac{2}{nm} \sin \theta_B \cos(\alpha - \theta_B)$$

#### Peak efficiencies at blaze wavelengths

150 /mm


## **Common spectrograph configurations**



## The Littrow spectrograph



Incident angle equals diffracted angle:

 $\alpha = \beta$ 

So for Littrow:

$$\lambda = \frac{2\sigma \sin \alpha}{m}$$

Simplifies the grating design, setting the blaze angle so that optimum efficiency is for  $\,\alpha\,$ 

# Detector

The smallest resolution for the spectrograph should be sampled at the minimum of the Nyquist frequency, which is 2 pixels per resolution element.



#### Spectral dispersion per pixel is:

$$\mu \frac{d\lambda}{dl}$$

where  $\mu$  is the pixel size in mm.



A Michelson interferometer with one moving arm

## Consider a monochromatic wave with:

$$k = 2\pi/\lambda$$

Electric field is then:  $e^{i(\omega t - kx)}$ 



At output of interferometer, the amplitude A is:

$$A = \frac{1}{2}e^{i\omega t}(e^{-ikx_1} + e^{-ikx_2})$$

Intensity output is: 
$$AA^* = \frac{1}{2}(1 + \cos k(x_2 - x_1))$$

Adding up all the incoherent intensities from a star with spectral distribution B(k) and taking  $x = x_2 - x_1$  and  $I_0$  as a constant, you can rewrite it as:

$$I(x) = I_0 + \frac{1}{2} \int_0^\infty B(k) \cos kx \, \mathrm{d}k$$

$$I(x) = I_0 + \frac{1}{2} \int_0^\infty B(k) \cos kx \, \mathrm{d}k$$

You can measure I(x) and get the spectral distribution back with a cosine fourier transform of  $I(x) - I_0$ 

Spectral resolution is given by largest path length difference L:

$$\Delta k = 2\pi/L$$

$$\lambda/\delta\lambda = 2 \times 10^6$$

PROS: Simple, compact, absolute calibration of spectral lines possible

CONS: very susceptible to any change in background flux



#### 1m Kitt Peak FTS - Eglin, Hanna, NOAO/AURA/NSF

# **Multi-Object Spectrographs**





#### Multi Object Spectrographs - drilled spectro slits

**DEIMOS slit masks milled with 0.015 inch diameter bits** 

![](_page_46_Picture_2.jpeg)

![](_page_46_Picture_3.jpeg)

#### http://www.ucolick.org/~phillips/deimos\_ref/masks.html

#### **Multi Object Spectrographs - laser cut slits**

IMACS on the Magellan 6.5m telescope

First spectrum with 240 slits

http://www.lco.cl/telescopes-information/magellan/instruments/imacs/

![](_page_47_Picture_4.jpeg)

#### Multi Object Spectrographs - laser cut slits

**VIMOS on the VLT telescopes** 

You decide where to put the slits on the science field

Can take up to two weeks to manufacture

![](_page_48_Picture_4.jpeg)

#### Multi Object Spectrographs - laser cut slits

VIMOS on the VLT telescopes

Number of spectra limited by sky coverage

http://www.eso.org/public/news/eso0209/

![](_page_49_Picture_4.jpeg)

### Night sky emission lines in NIR

VIMOS on the VLT telescopes

![](_page_50_Picture_2.jpeg)

## **Configurable slits on MOSFIRE (Keck)**

NIR multi-object spectrograph

![](_page_51_Figure_2.jpeg)

## **Configurable Slit Unit (CSU)**

Cryogenic slits can be reconfigured in cold and in vacuum dewar!

![](_page_52_Picture_2.jpeg)

![](_page_52_Picture_3.jpeg)

## **Configurable slits on MOSFIRE (Keck)**

Adjustable mechanical slits allow for much faster configuration

![](_page_53_Picture_2.jpeg)

McLean 2012

Figure 7. On the left is the layout of the MOSFIRE field on the sky with a 58s J-band image of The Antennae galaxies. The middle image is of a slit mask and the right image is the night sky emission with this mask in H-band.

# **Fibre Optics**

![](_page_54_Picture_1.jpeg)

# **Fibre Optics**

![](_page_55_Figure_1.jpeg)

#### Structure of an optical fibre

Cladding has higher refractive index than core material

![](_page_56_Figure_2.jpeg)

Largest acceptance angle dependent on core and cladding refractive indices

#### Everything is big when you are 100 microns in size

![](_page_57_Figure_1.jpeg)

**Figure 2-14** *Examining the fibre faces*. On the left the fibre face is checked for micro-pits - several can be clearly seen. On the right the back-illuminated fibre shows a clean ring of light across the face of the fibre.

#### **Optical Fibre - azimuthal scrambling**

![](_page_58_Figure_1.jpeg)

#### Thousands of internal reflections from curved interface

#### **Optical Fibre - azimuthal scrambling**

![](_page_59_Picture_1.jpeg)

#### Looking down on the end of the polished fibre end

#### **Optical Fibre - Focal Ratio Degradation**

![](_page_60_Figure_1.jpeg)

Deformation and stress causes light to 'spread' in output angle cone

# Loss of flux from FRD if you don't make the output optics bigger in size

![](_page_61_Figure_1.jpeg)

**Figure 2-19** *FRD test results for seven SPIRAL fibres*. This is for a 16m length of Polymicro 50/70/90/110µm fibre measured at 600nm.

#### Plug plates drilled manually to match target fields

![](_page_62_Figure_1.jpeg)

#### Hill 1988 ASPC

Figure 2. Schematic drawing of the aperture plate nucleus of the MEDUSA spectrograph.

#### **Optical Fibre Spectrograph**

![](_page_63_Figure_1.jpeg)

#### Gluing optical fibres onto a glass plate(!)

The fine plate-scale (67 arcsec/mm) meant that drilling holes in brass plates was not an option for fibre positioning, due to thermal and other considerations. The required positioning accuracy for the fibres was 10  $\mu$ m over the whole field (think of sticking a pin in a cricket pitch with a precision of 1 mm). It was one of the editors of these proceedings who suggested a viable alternative. Tacking the fibres directly onto transparent star and galaxy images on a positive copy of the target field using UV-curing cement seemed like a blindingly obvious solution to David Malin, with his background in photography and polymer chemistry.

Unlikely as it sounds, this technique worked rather well when it was tried out late in 1983. It required a special plate-holder to support the glass positive plate and bend it to the focal curvature. This had the same dimensions as the photographic plate-holders, so it could be loaded via the existing elevator, and was built for the project by UKST technicians Eric Coyte and Magnus Paterson (Fig. 3). It was another nine months before the necessary components for a fibre acquisition system had been built, but by October 1984, sets of stars spread over the full 6.5 degrees square field of the telescope were being simultaneously acquired. By then, too, the system had a name – FLAIR, for Fibre-Linked Array Image Reformatter. What else?

![](_page_64_Picture_3.jpeg)

**Fred Watson** 

### Sits at prime focus of 4m Anglo-Australian Telescope

400 fibres positioned whilst other 400 are observing!

![](_page_65_Picture_3.jpeg)

#### Diameter of 140 microns (2.1 arcsec on the sky)

![](_page_66_Figure_1.jpeg)

![](_page_67_Figure_1.jpeg)

Lewis (2002)

![](_page_68_Picture_1.jpeg)

![](_page_69_Picture_0.jpeg)

![](_page_70_Figure_1.jpeg)

optical fibres are backlit with a red LED to help positioning - typically 11 microns rms precision
## **3D Spectroscopy**

#### **3D spectroscopy**



#### **Optical Fibre Image Reformatter**

Use the flexibility of fibres to reformat the 2D sky into a 1D entrance slit



Fibers take light from telescope focus to the spectrograph

Spectrograph disperses the light

# Spectrograph can sit on floor of the observatory instead of at the focus





#### **Optical Fibre**



**Figure 6-1** *A raw IFS data frame*. In this data frame from SPIRAL the dispersion axis is across the page and the 37 separate fibre tracks can be seen. This is a twilight sky exposure, clearly showing absorption features in the atmosphere and the variation in throughput between fibres.

#### Hexagonal lenslets on the sky



**Figure 9-1** *Image reconstruction using the LDISPLAY software*. The image of the left is a raw image from the COHSI spectrograph. By knowing the relation between fibres on the sky and fibres in the slit an image can be reconstructed (right-hand panel).





**Disperse these images into short spectra -> limited wavelength range** 

#### **MUSE on VLT - 24 Integral Field Spectrographs**



#### **WOSE VEIDCITY HEIDS**



### Complete wavelength coverage gives both abundances, velocity fields of different species of atomic transition







#### **MUSE on VLT - 24 Integral Field Spectrographs**

