

# **Astronomical Observing Techniques**

## **Lecture 4: Your Noise is My Signal**

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# Outline

1. Introduction
2. Statistics
3. Signal-to-Noise Ratio
4. Instrument Sensitivities

# Noise

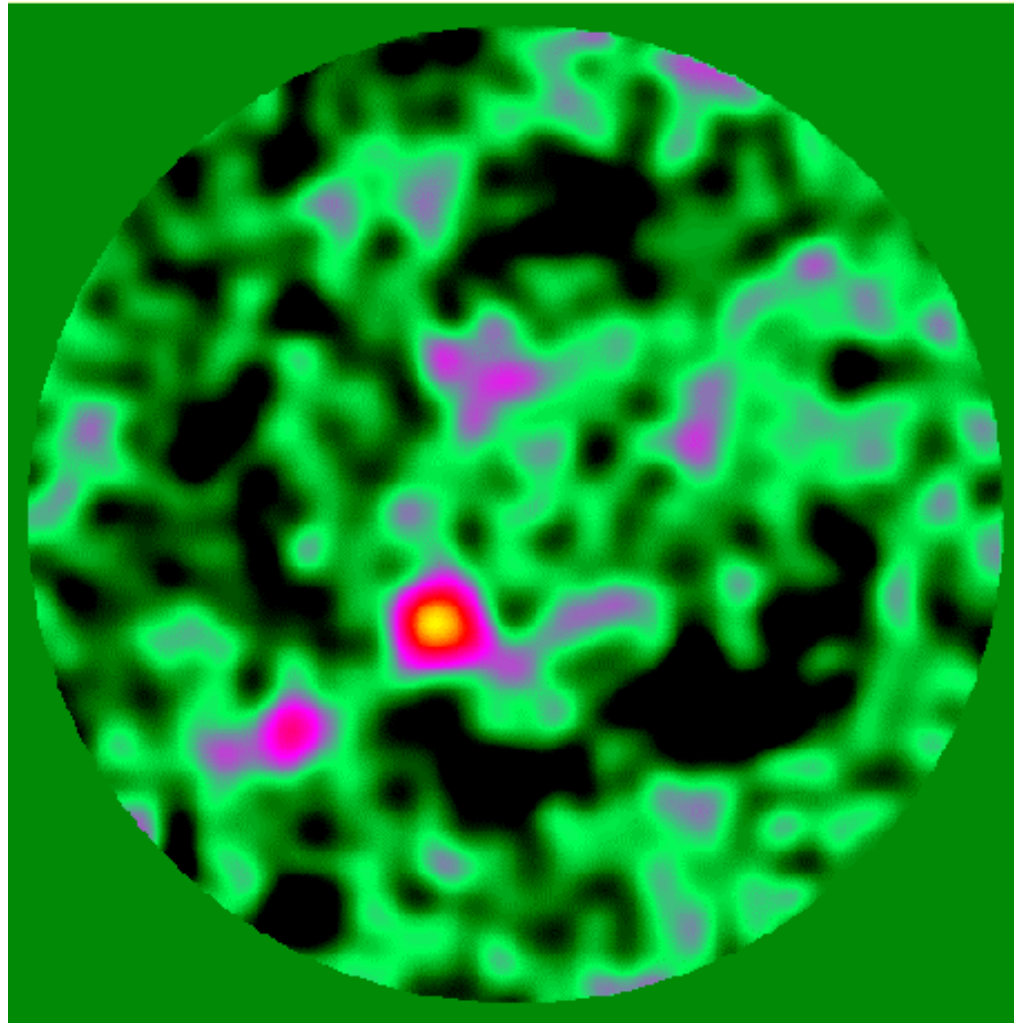
from Wikipedia:

- Common use: **unwanted sound**
- Signal processing: **random unwanted data without meaning**
- "Signal-to-noise ratio" is sometimes used to refer to the ratio of **useful to irrelevant information**



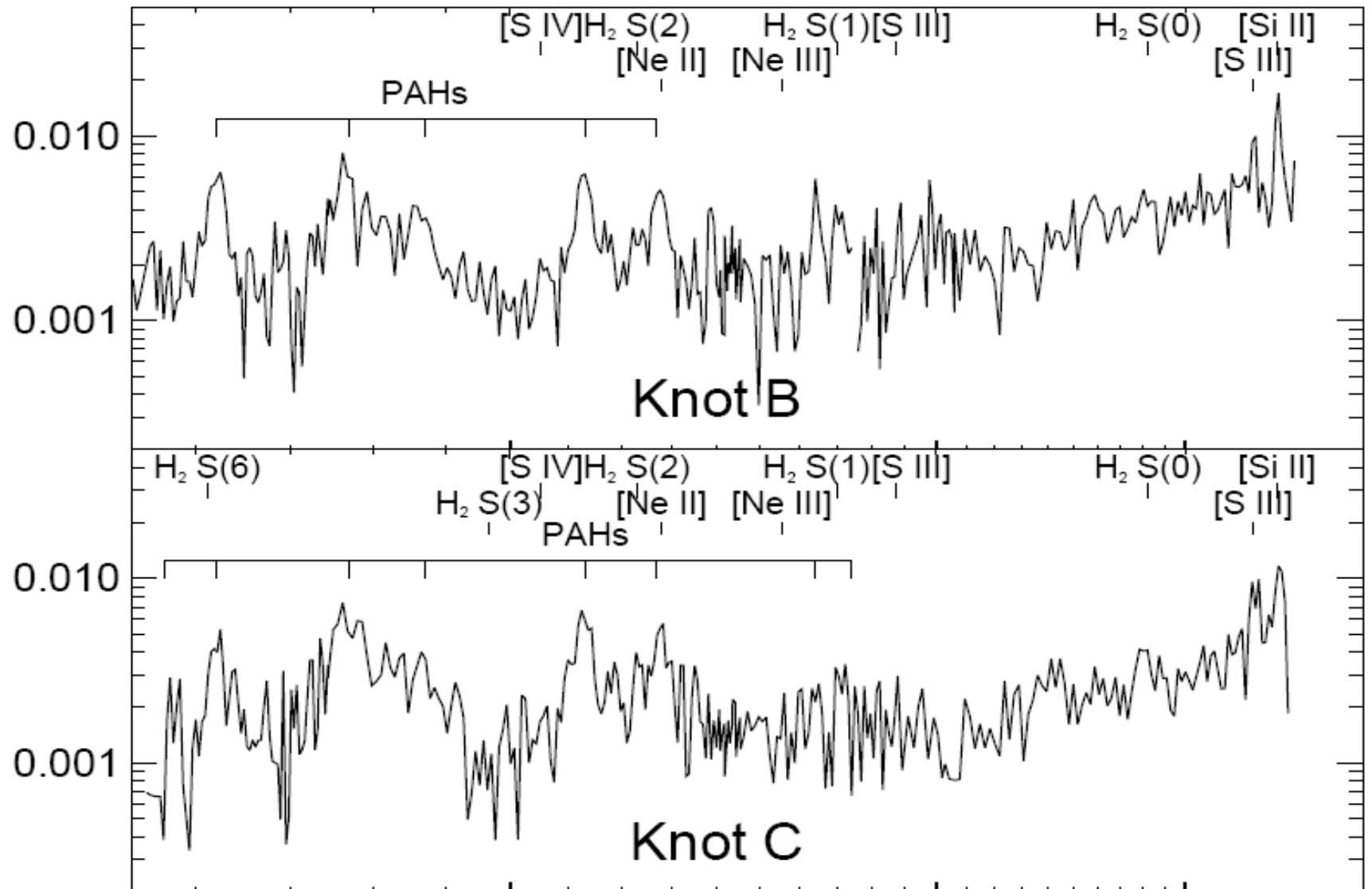
*NASA researchers at Glenn Research Center conducting tests on aircraft engine noise in 1967*

# Signal?



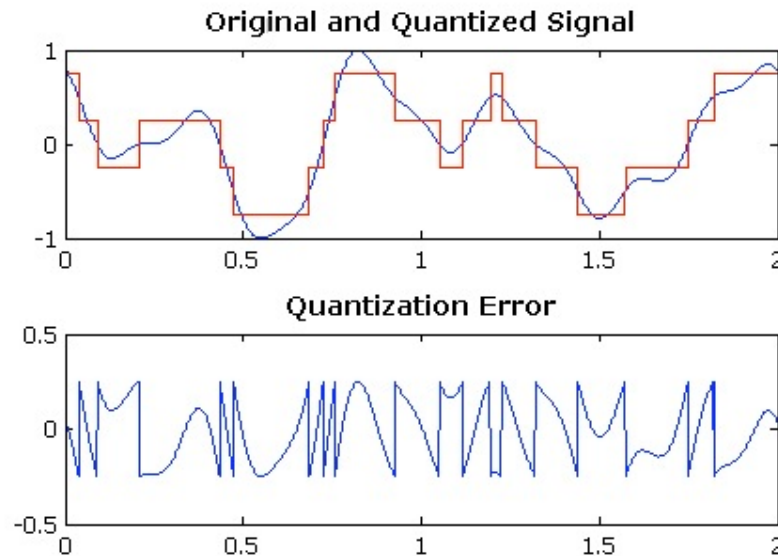
SCUBA 850 $\mu$ m map of the Hubble deep field

# Noise or Signal?

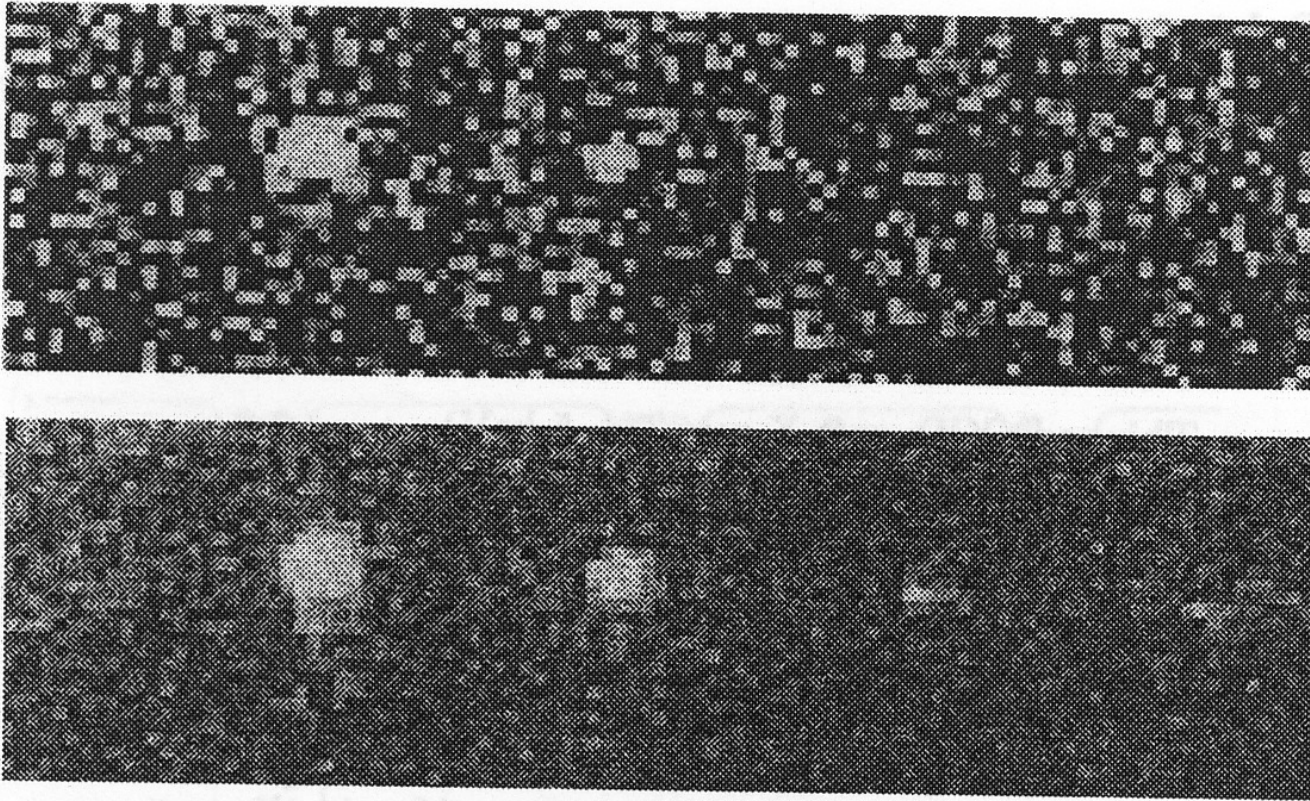


# Digitization/Quantization Noise

- Analog-to-Digital Signal Converter (ADC).
- Number of bits determines **dynamic range** of ADC
- Resolution: 12 bit  $2^{12} = 4096$  quantization levels  
16 bit  $2^{16} = 65536$  quantization levels
- Discrete, “artificial” steps in signal levels → noise



# Read Noise



**Fig. 7.28.** (a) Image of four source points, by a CCD with  $\sigma_R = 7.6 e^-$  rms. (b) The same image in multiple readout ( $N = 64$ ), where  $\sigma_R = 0.97 e^-$  rms. The faintest source corresponds to a signal of 3.5 photocharges. (After Janesik et al., in *The CCD in Astronomy*, ASP Conf. Ser. 8, 1989)

# Some Noise Sources in Astronomical Data

Noise type	Signal	Background
Photon shot noise	X	X
Scintillation	X	
Cosmic rays		X
Image stability	X	
Read noise	X	X
Dark current noise	X	X
CTE (CCDs)	X	X
Flat fielding (non-linearity)	X	X
Digitization noise	X	X
Other calibration errors	X	X
Image subtraction	X	X



# Distribution Functions

- for every  $t$ ,  $X(t)$  is distributed according to cumulative distribution function

$$F(x;t) = \mathbf{P}\{X(t) \leq x\}$$

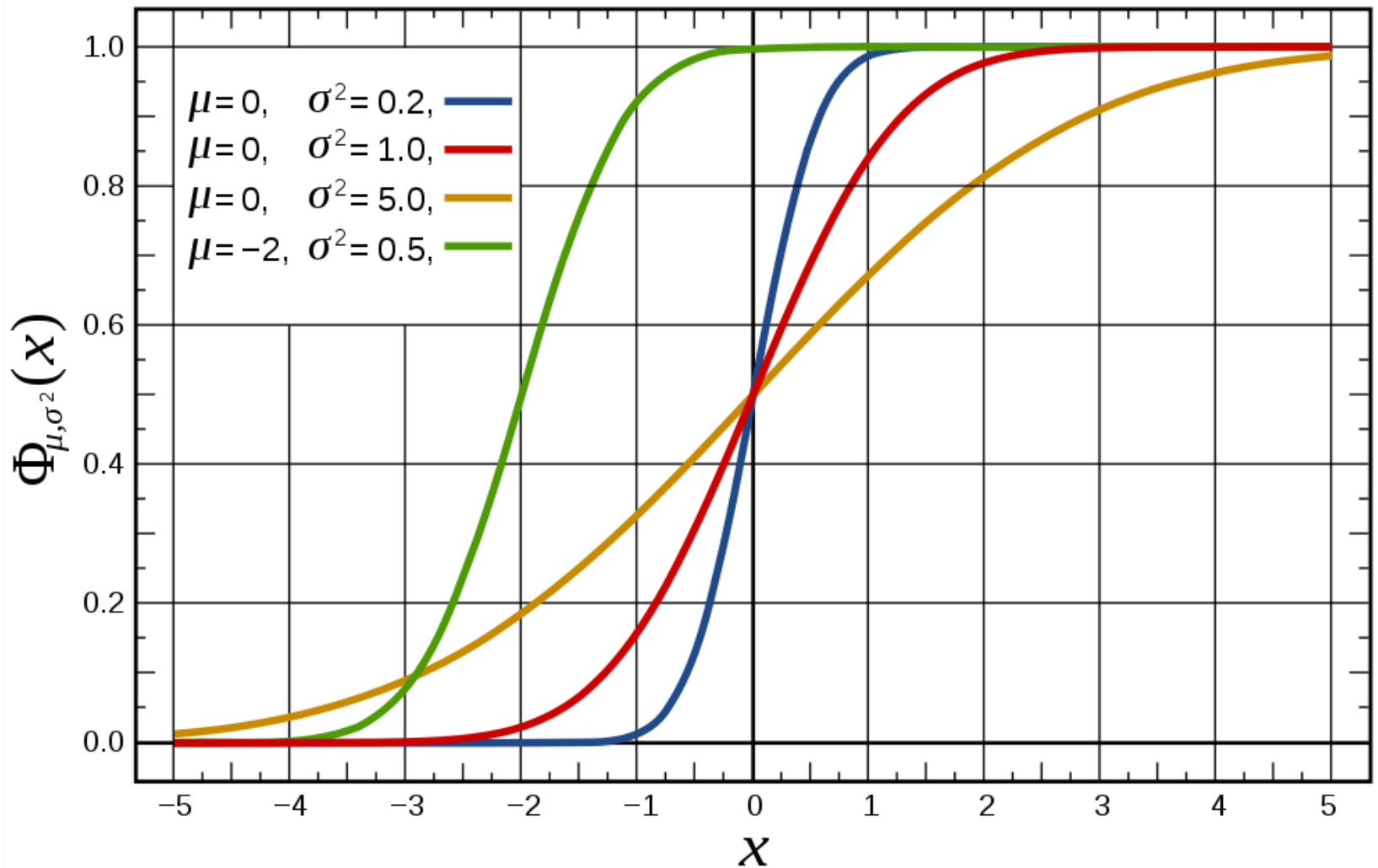
- indicates probability that outcome at  $t$  will not exceed  $x$
- probability density function (PDF) of  $X(t)$  defined by

$$f(x;t) \equiv \frac{\partial F(x;t)}{\partial x}$$

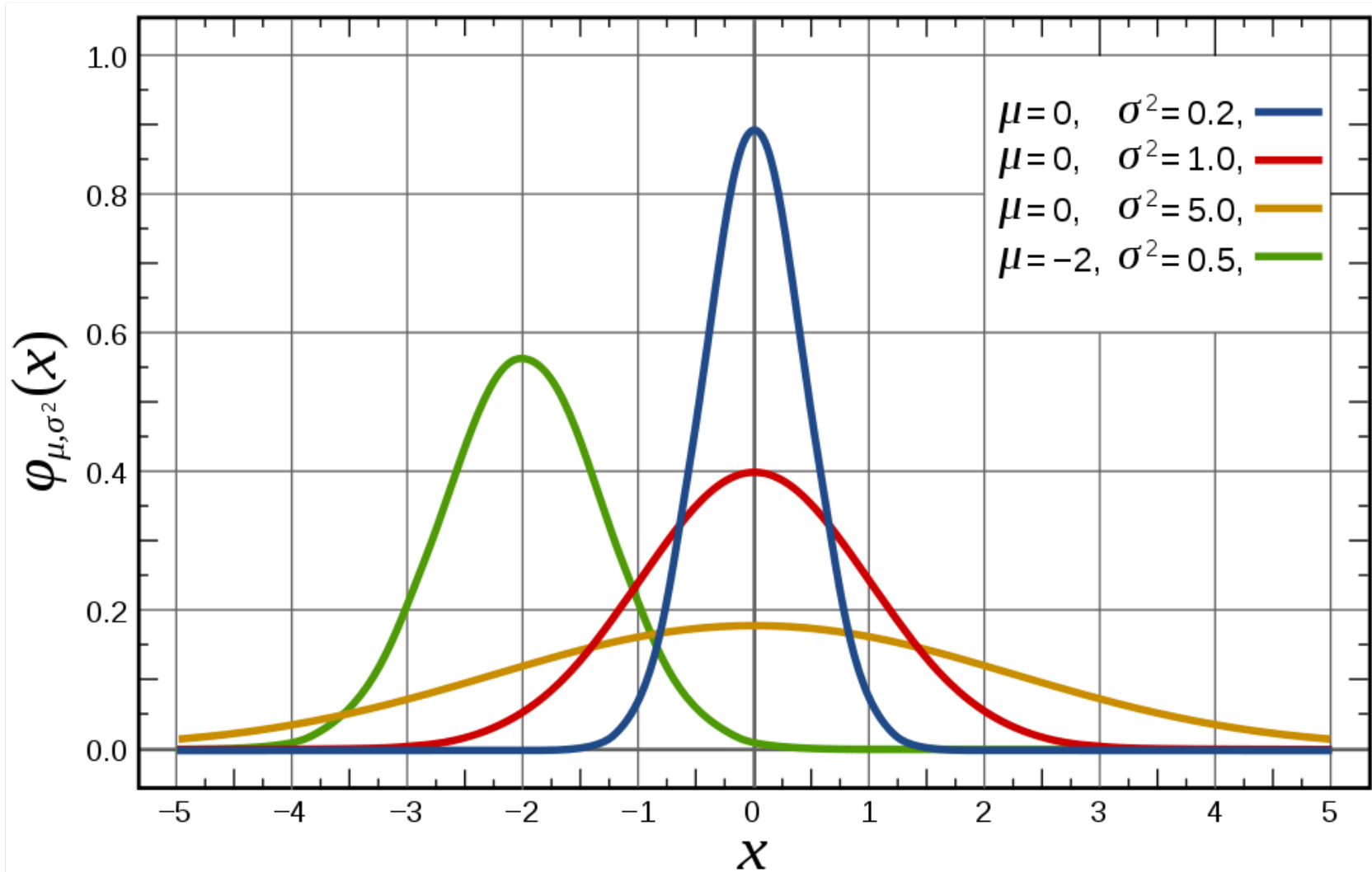
# Typical Probability Density Functions

- binomial:  $\binom{n}{k} p^k (1-p)^{n-k}$
- Poisson:  $f(k; \lambda) = \frac{\lambda^k e^{-\lambda}}{k!}$
- Gaussian:  
(normal)  $\frac{1}{\sqrt{2\pi}} e^{-\frac{(x-\mu)^2}{2\sigma^2}}$

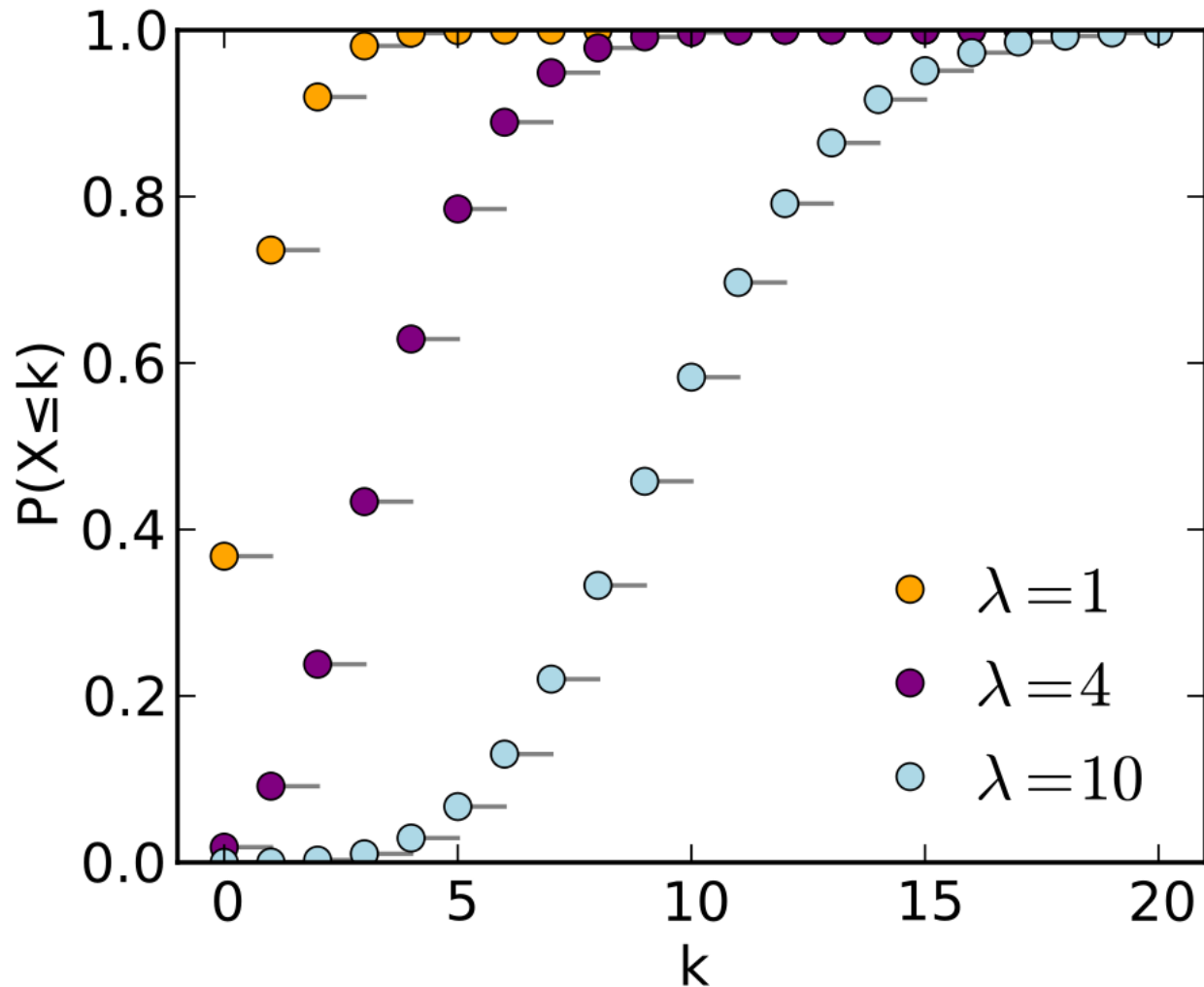
# Normal Cumulative Distribution



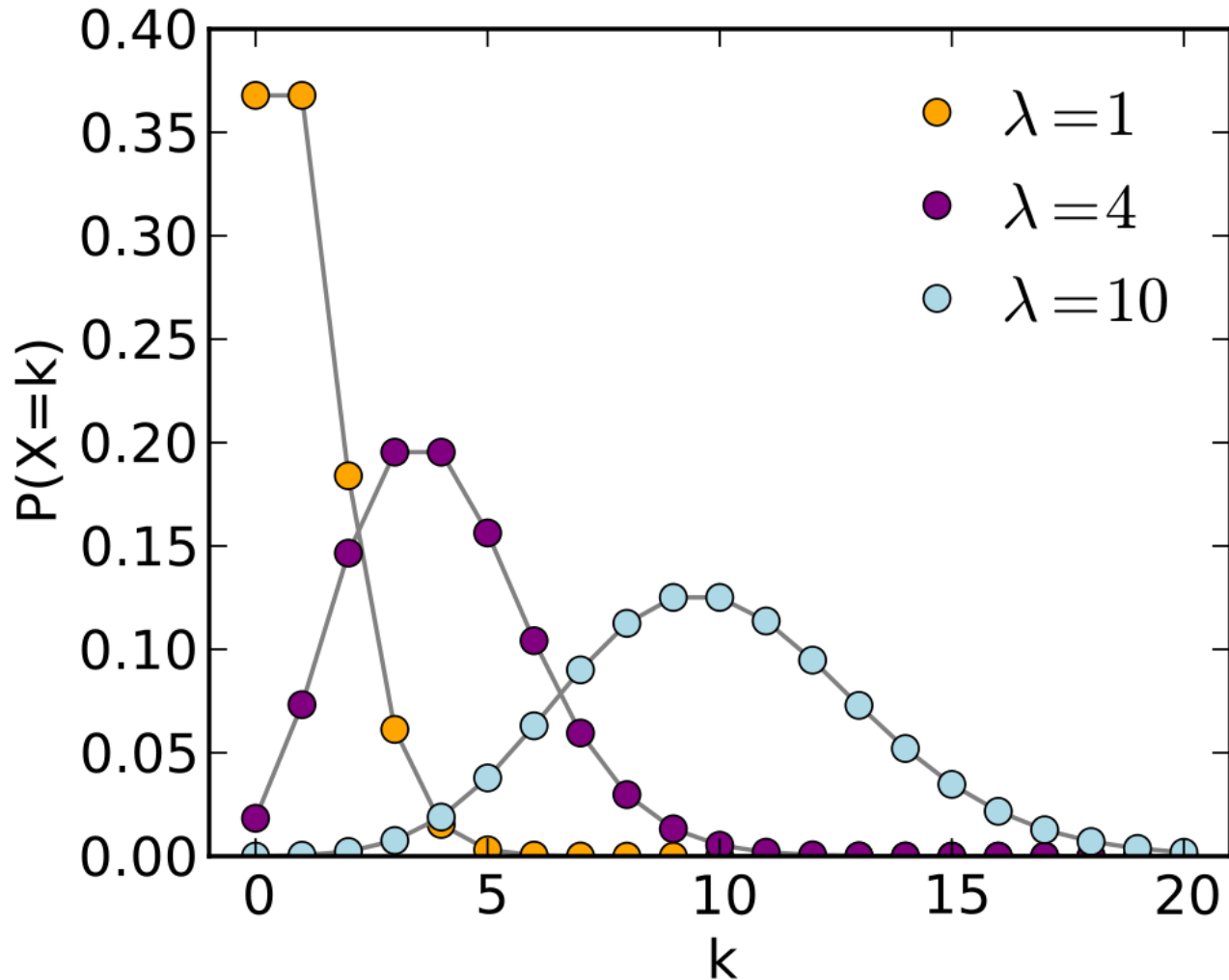
# Normal/Gaussian PDF



# Poisson Cumulative Distribution Function



# Poisson PDF



# Mean, Variance, RMS

- properties of distributions often described by a few parameters, often **moments** of distribution

- **mean** or **average**  $\mu(t)$  of  $X(t)$ : expected value of  $X(t)$

$$\mu(t) = \mathbf{E}\{X(t)\} = \int_{-\infty}^{+\infty} x f(x;t) dx$$

- **variance** of  $X(t)$ : expected value of the square of the difference of  $X(t)$  and  $\mu(t)$

$$\sigma^2(t) = \mathbf{E}\{(X(t) - \mu(t))^2\} = \mathbf{E}\{X^2(t)\} - \mu^2(t)$$

- variance is square of **standard deviation** or **root mean square (RMS)**

# Noise Propagation

- same as error propagation
- function  $f(u, v, \dots)$  depends on variables  $u, v, \dots$
- estimate variance of  $f$  knowing variances  $\sigma_u^2$ ,  $\sigma_v^2, \dots$  of variables  $u, v, \dots$

$$\sigma_f^2 \equiv \lim_{N \rightarrow \infty} \sum_{i=1}^N (f_i - \bar{f})^2$$

- make assumption / approximately that average of  $f$  is well approximated by value of  $f$  for averages of variables:

$$\bar{f} = f(\bar{u}, \bar{v}, \dots)$$



# Noise Propagation (cont.)

- Taylor expansion of  $f$  around average:

$$f_i - \bar{f} \approx (u_i - \bar{u}) \frac{\partial f}{\partial u} + (v_i - \bar{v}) \frac{\partial f}{\partial v} + \dots$$

- variance in  $f$ :

$$\begin{aligned} \sigma_f^2 &\approx \lim_{N \rightarrow \infty} \sum_{i=1}^N \left[ (u_i - \bar{u}) \frac{\partial f}{\partial u} + (v_i - \bar{v}) \frac{\partial f}{\partial v} + \dots \right]^2 \\ &= \lim_{N \rightarrow \infty} \sum_{i=1}^N \left[ (u_i - \bar{u})^2 \left( \frac{\partial f}{\partial u} \right)^2 + (v_i - \bar{v})^2 \left( \frac{\partial f}{\partial v} \right)^2 + 2(u_i - \bar{u})(v_i - \bar{v}) \frac{\partial f}{\partial u} \frac{\partial f}{\partial v} + \dots \right] \end{aligned}$$

# Noise Propagation (cont.)

- variances of  $u$  and  $v$

$$\sigma_u^2 \equiv \lim_{N \rightarrow \infty} \sum_{i=1}^N (u_i - \bar{u})^2; \quad \sigma_v^2 \equiv \lim_{N \rightarrow \infty} \sum_{i=1}^N (v_i - \bar{v})^2$$

- covariance of  $u$  and  $v$

$$\sigma_{uv}^2 \equiv \lim_{N \rightarrow \infty} \sum_{i=1}^N (u_i - \bar{u})(v_i - \bar{v})$$

- combine Taylor expansion and these definitions

$$\sigma_f^2 = \sigma_u^2 \left( \frac{\partial f}{\partial u} \right)^2 + \sigma_v^2 \left( \frac{\partial f}{\partial v} \right)^2 + 2\sigma_{uv}^2 \frac{\partial f}{\partial u} \frac{\partial f}{\partial v} + \dots$$

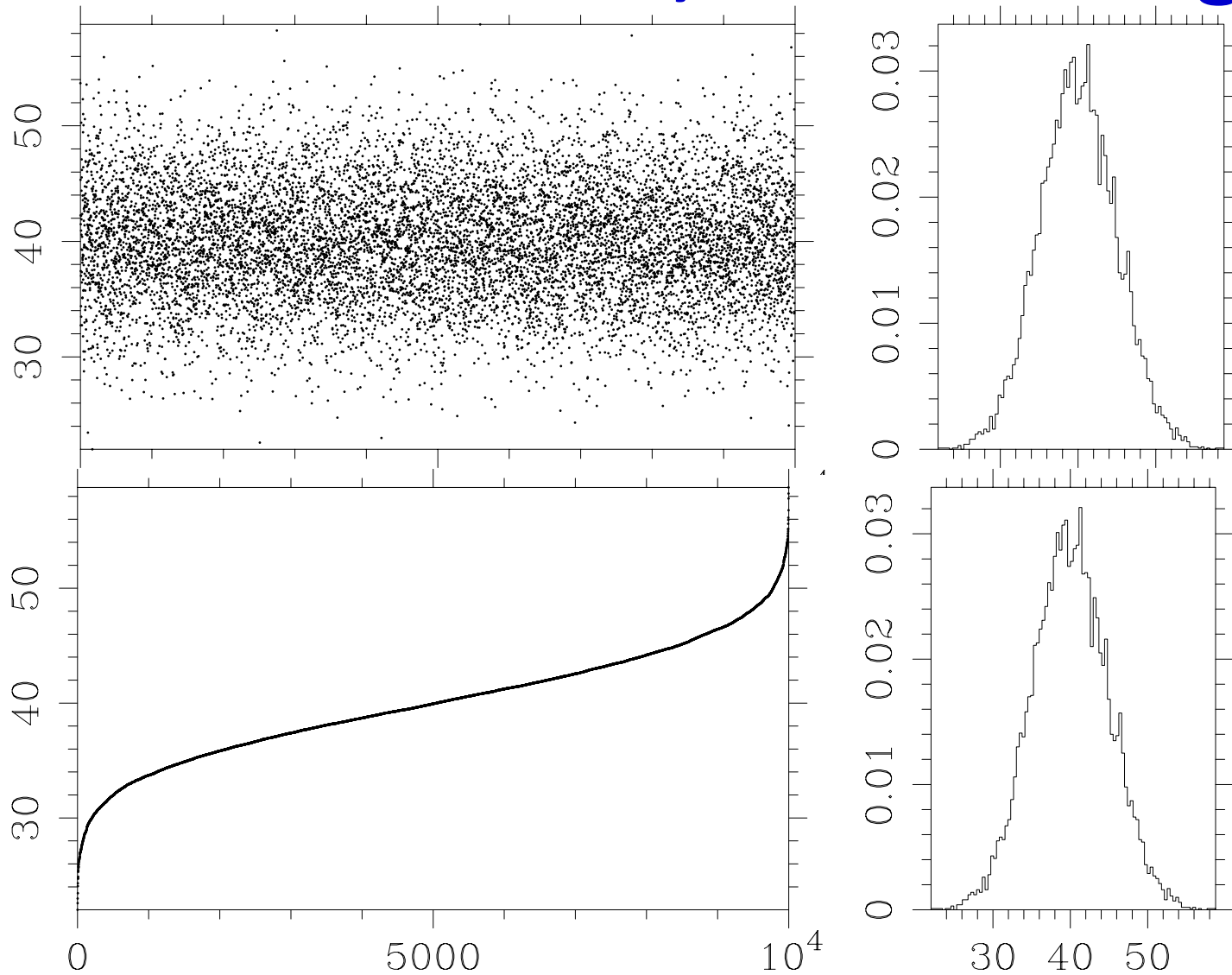
# Noise Propagation (cont.)

- from before

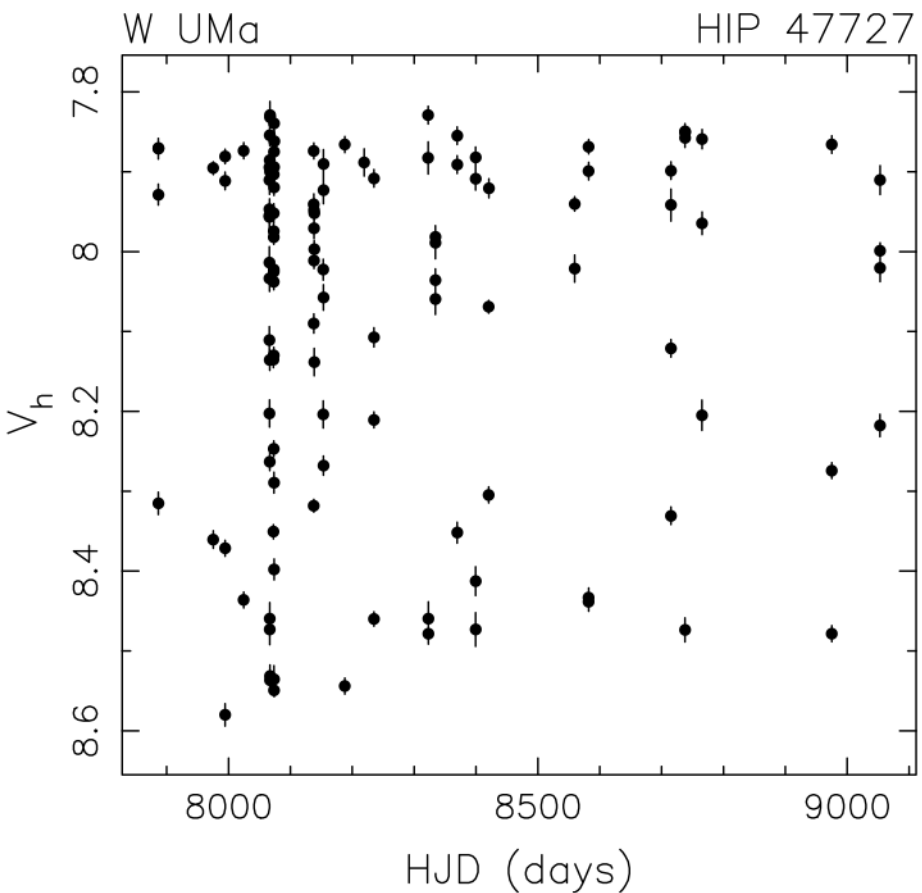
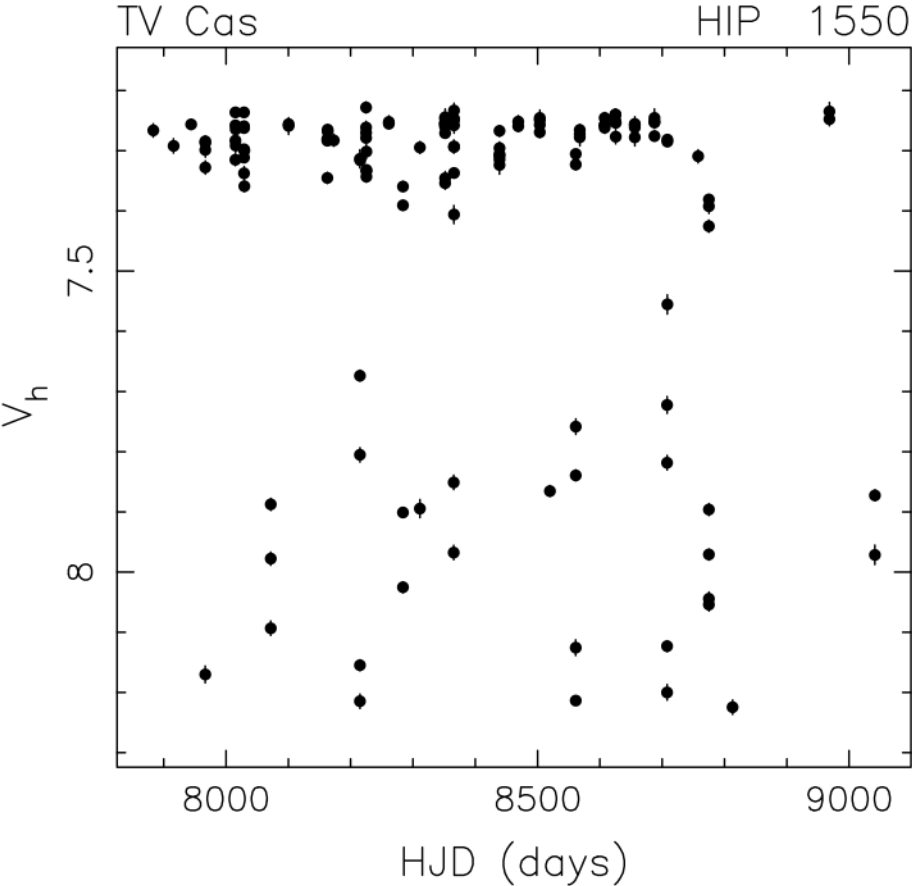
$$\sigma_f^2 = \sigma_u^2 \left( \frac{\partial f}{\partial u} \right)^2 + \sigma_v^2 \left( \frac{\partial f}{\partial v} \right)^2 + 2\sigma_{uv} \frac{\partial f}{\partial u} \frac{\partial f}{\partial v} + \dots$$

- if differences  $(u_i - \bar{u})$  and  $(v_i - \bar{v})$  not correlated  $\rightarrow$   
sign of product as often positive as negative  $\rightarrow$   
covariance small compared to other terms
- if differences are correlated  $\rightarrow$  most products  
 $(u_i - \bar{u})(v_i - \bar{v})$  have the same sign  $\rightarrow$  cross-  
correlation term can be large

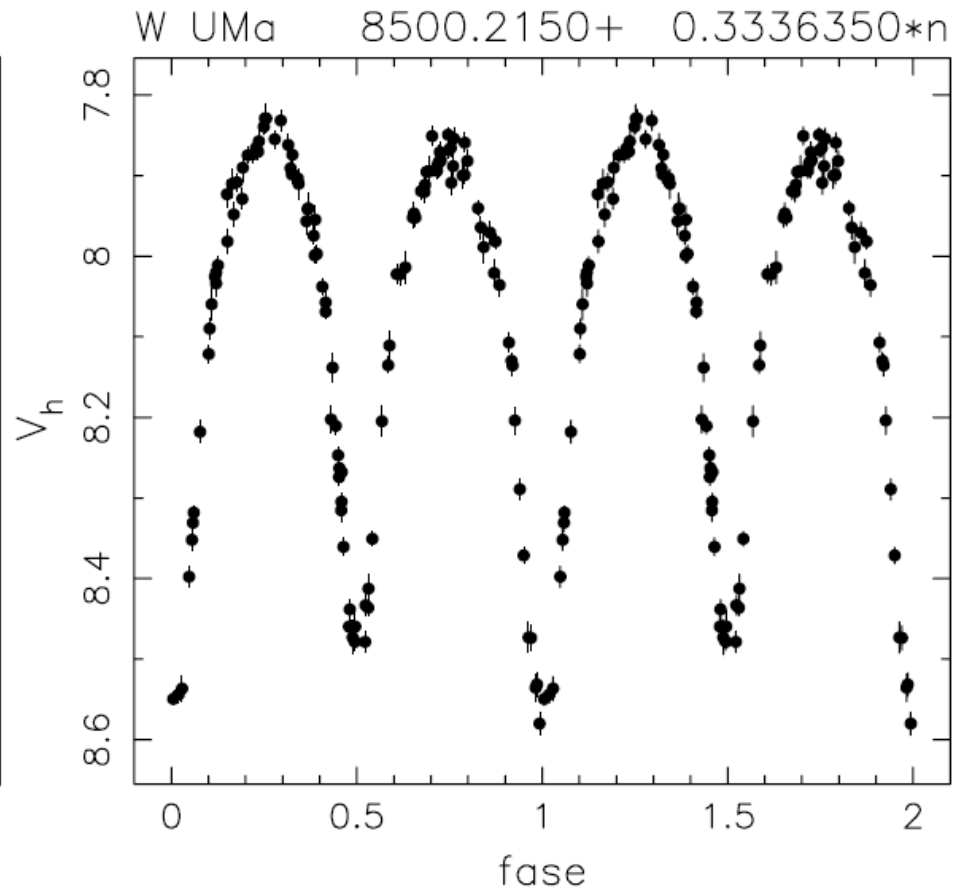
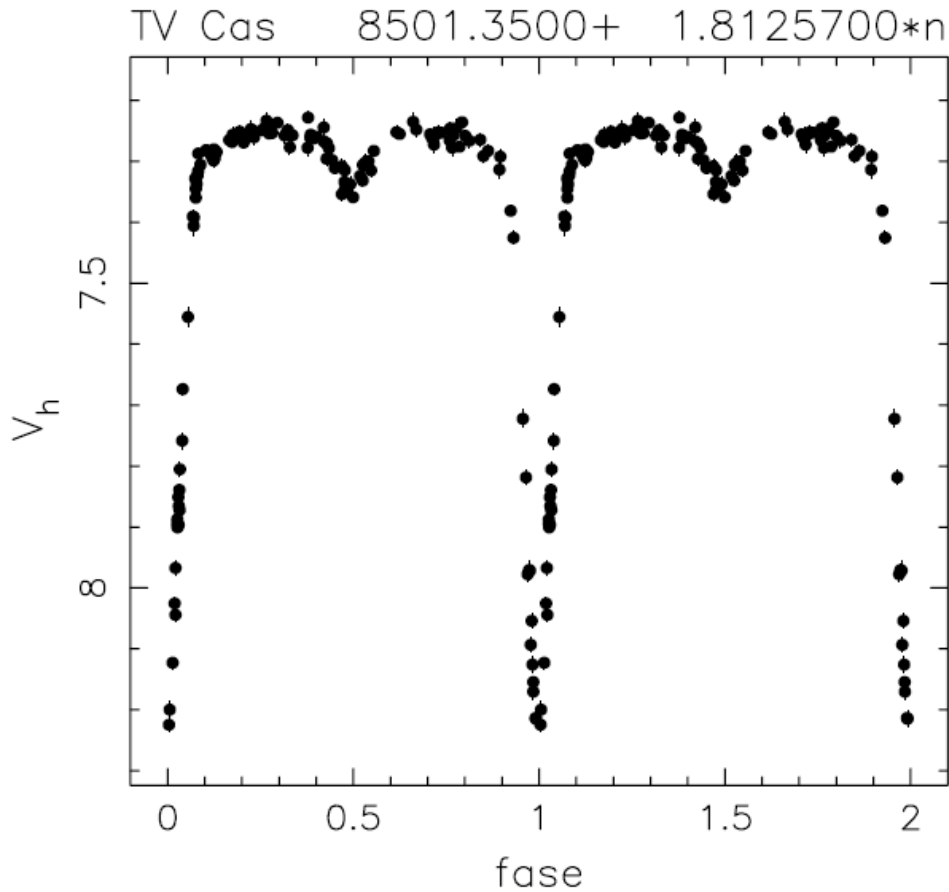
# Same Distribution, Different Signals



# Time Series of Two Stars



# Two Stars: Differently Sorted



# Gaussian Noise

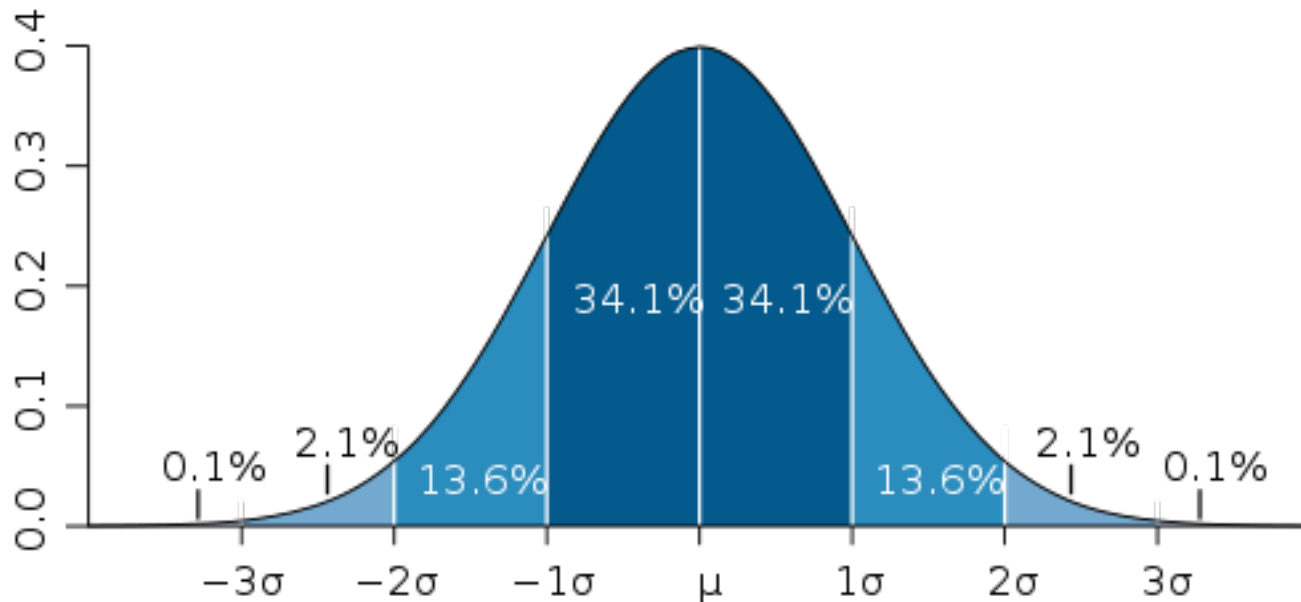
- Gaussian noise has Gaussian (normal) distribution
- Sometimes (incorrectly) called white noise (uncorrelated noise)

$$S = \frac{1}{\sqrt{2\pi}\sigma} \exp\left[-\frac{(x-\mu)^2}{2\sigma^2}\right]$$

$x$ : actual value

$\mu$ : mean of distribution

$\sigma$ : standard deviation of distribution



1- $\sigma$  ~ 68%

2- $\sigma$  ~ 95%

3- $\sigma$  ~ 99.7%

Astronomers usually consider  $S/N > 3\sigma$  as significant.

# Poisson Noise

- **Poisson noise** has Poisson distribution
- probability of number of events occurring in constant interval of time/space **if** events occur with known *average rate* and *independently* of each other

$$P(k, \lambda) = \frac{\lambda^k e^{-\lambda}}{k!}$$

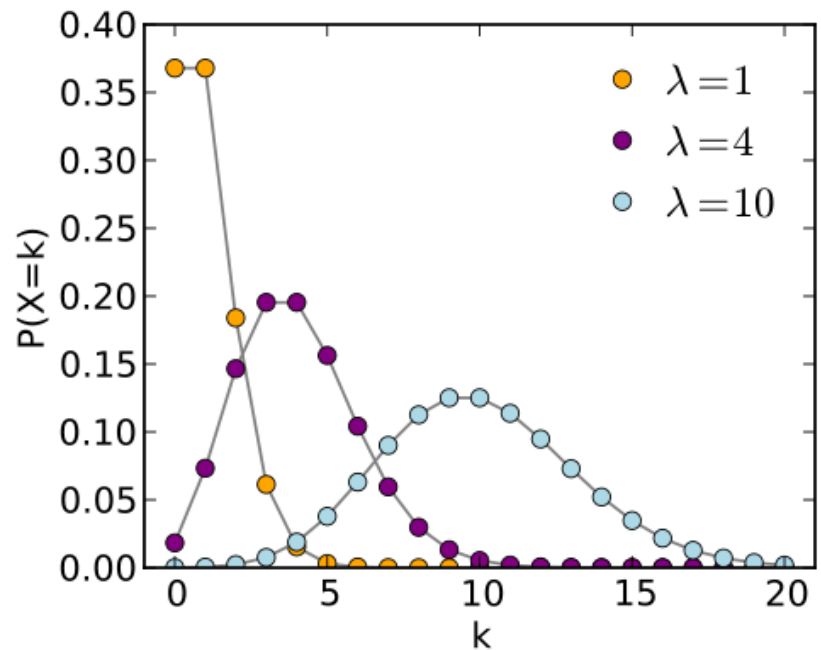
$k$ : number of occurrences of an event (probability)

$\lambda$ : *expected* (average) number of occurrences

- **mean** of  $P(k, \lambda)$  is  $\lambda$

- **standard deviation** of  $P(k, \lambda)$  is **square root of  $\lambda$**

- example: fluctuations in photon flux in finite time intervals  $\Delta t$ .  
Chance to detect  $k$  photons with average flux of  $\lambda$  photons



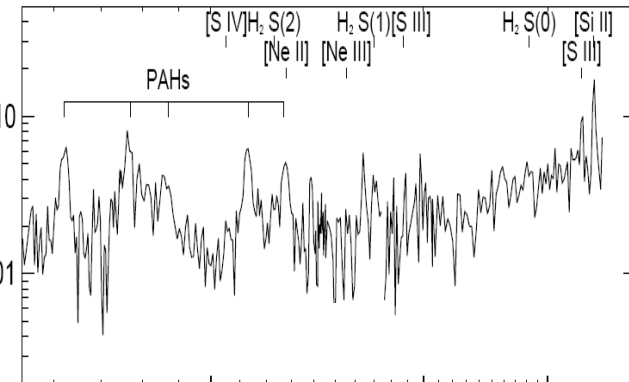
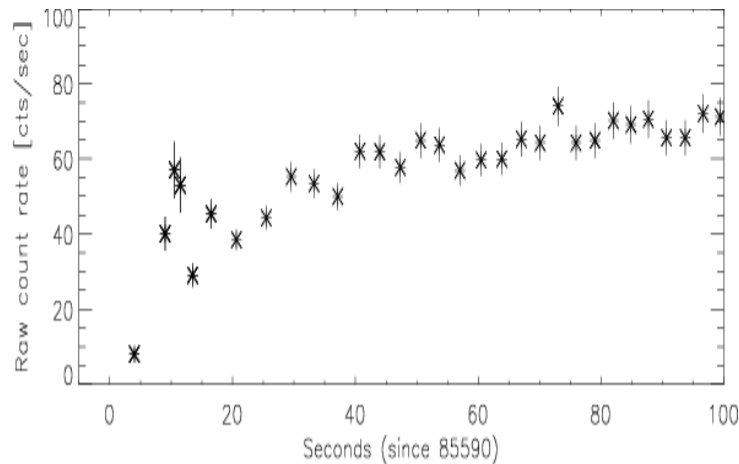
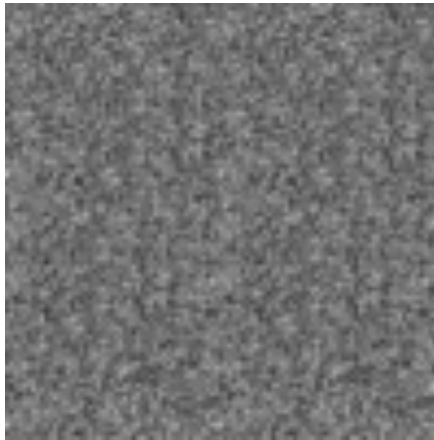


# Noise Measurement

If purely Gauss or Poisson noise distribution, no other systematic noise and no correlations,

*then the spatial distribution (neighbouring pixels) of the noise is equivalent to the temporal distribution (successive measurements with one pixel)*

*This is analogous to throwing 5 dices once versus throwing one dice*



Case 1: Spatial noise  
(detector pixels)

Case 2: Repeated measurements  
in time (time series)

Case 3: Spectrum  
(dispersed information)

# Poisson Noise and Integration Time

- Integrate light from uniform, extended source on CCD
- In finite time interval  $\Delta t$ , expect average of  $\lambda$  photons
- Statistical nature of photon arrival rate  $\rightarrow$  some pixels will detect more, some less than  $\lambda$  photons.
- Noise of average signal  $\lambda$  (i.e., between pixels) is  $\sqrt{\lambda}$
- Integrate for  $2 \times \Delta t \rightarrow$  expect average of  $2 \times \lambda$  photons
- Noise of that signal is now  $\sqrt{2 \times \lambda}$ , i.e., increased by  $\sqrt{2}$
- **With respect to integration time  $t$ , noise will only increase  $\sim \sqrt{t}$  while signal increases  $\sim t$**

# Signal-to-Noise Ratio

from Wikipedia:

*Signal-to-noise ratio (often abbreviated SNR or S/N) is a measure used in science and engineering that compares the level of a desired signal to the level of background noise.*

Signal = S; Background = B; Noise = N;

$$\sigma = \frac{\text{Signal}}{\text{Noise}}$$

← measured as  $(S+B) - \text{mean}\{B\}$

← total noise =  $\sqrt{\sum (N_i)^2}$  (if statistically independent)

Both  $S$  and  $N$  should be in units of events (photons, electrons, data numbers) per unit area (pixel, PSF size, arcsec<sup>2</sup>)

# S/N and Integration Time

Assuming the signal suffers from **Poisson shot noise**. Let's calculate the dependence on **integration time**  $t_{\text{int}}$ :

Integrating  $t_{\text{int}}$ : 
$$\sigma = \frac{S}{N}$$

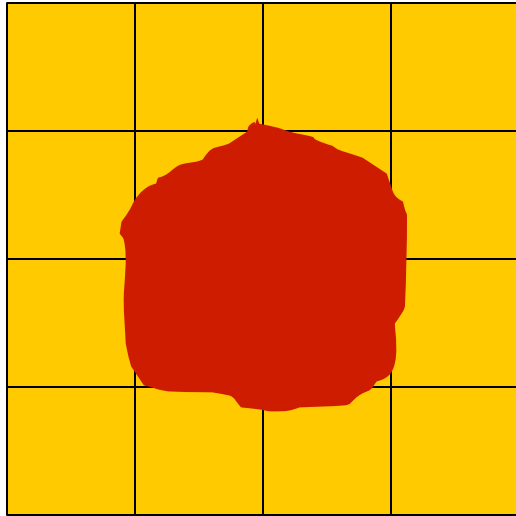
Integrating  $n \times t_{\text{int}}$ : 
$$\sigma = \frac{n \cdot S}{\sqrt{n \cdot B}} \stackrel{N=\sqrt{B}}{=} \sqrt{n} \frac{S}{N} \Rightarrow \frac{S}{N} \propto \sqrt{t_{\text{int}}}$$

*Need to integrate four times as long to get twice the S/N*

# 3 Cases

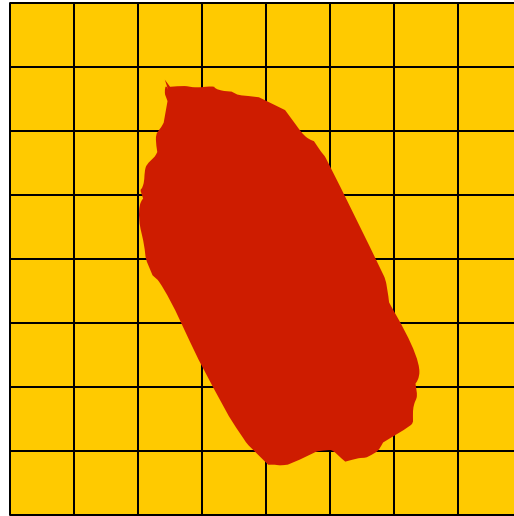
Background (=noise)

Target



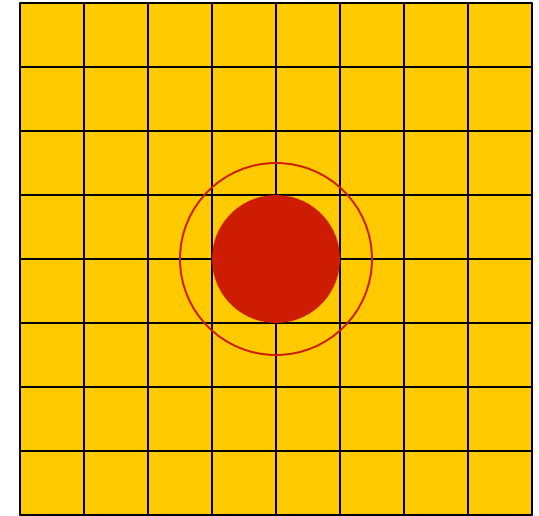
Seeing-limited  
point source

- pixel size  $\sim$  seeing
- PSF  $\neq$   $f(D)$



Diffraction-limited,  
extended source

- pixel size  $\sim$  diff.lim
- PSF =  $f(D)$
- target  $\gg$  PSF



Diffraction-limited,  
point source

- pixel size  $\sim$  diff.lim
- PSF =  $f(D)$
- target  $\ll$  PSF

# Case 1: Seeing-limited “Point Source”

Signal = S; Background = B; Noise = N; Telescope diameter = D

$$\theta_{\text{seeing}} \sim \text{const}$$

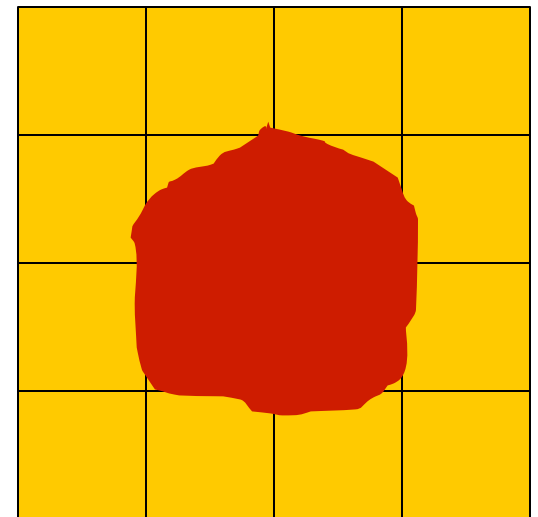
If detector is Nyquist-sampled to  $\theta_{\text{seeing}}$ :

$$S \sim D^2 \text{ (area)}$$

$$B \sim D^2 \rightarrow N \sim D \text{ (Poisson std.dev)}$$

$$\rightarrow S/N \sim D$$

$$\rightarrow t_{\text{int}} \sim D^{-2}$$



# Case 2: Diffraction-limited extended Source

Signal = S; Background = B; Noise = N; Telescope diameter = D

“Diameter” of PSF  $\sim$  const

If detector Nyquist sampled to  $\theta_{\text{diff}}$ : pixel  $\sim D^{-2}$  but  $S \sim D^2$

$D^2$  (telescope size) and  $D^{-2}$  (pixel FOV) cancel each other  $\rightarrow$  no change in signal

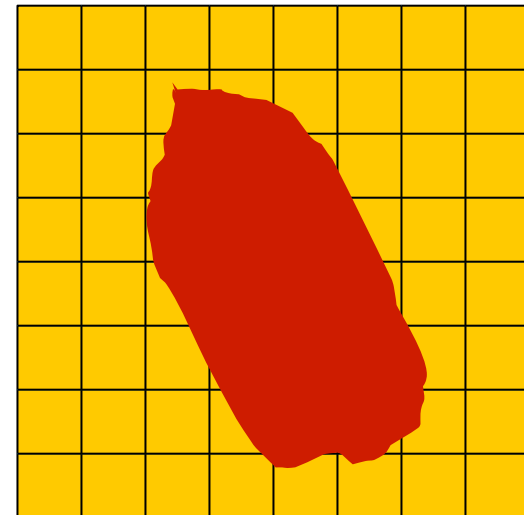
same for the background flux

$\rightarrow S/N \sim \text{const} \rightarrow t_{\text{int}} \sim \text{const}$

$\rightarrow$  no gain for larger telescopes!

Case 2B: offline re-sampling by a factor x  
(makes  $\theta_{\text{diff}}$  x-times larger)

since  $S/N \sim \sqrt{n_{\text{pix}}} \rightarrow S/N \sim \sqrt{x^2} = x \rightarrow t_{\text{int}} \sim x^{-2}$



# Case 3: Diffraction-limited “Point Source”

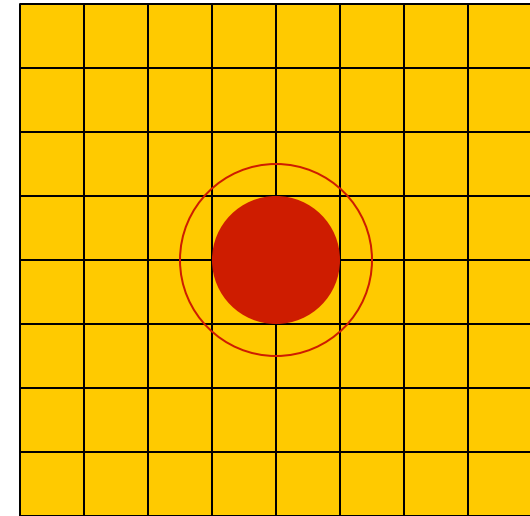
Signal = S; Background = B; Noise = N; Telescope diameter = D

“ $S/N = (S/N)_{\text{light bucket}} \cdot (S/N)_{\text{pixel scale}}$ ”

(i) Effect of telescope aperture:

$$\text{Signal} \quad S \sim D^2 \quad \rightarrow \quad S/N \sim D$$

$$\text{Background} \quad B \sim D^2 \rightarrow N \sim D$$



(ii) Effect of pixel FOV (if Nyquist sampled to  $\theta_{\text{diff}}$ ):

$$S \sim \text{const} \text{ (pixel samples PSF = all source flux)}$$

$$B \sim D^{-2} \rightarrow N \sim D^{-1} \quad \rightarrow \quad S/N \sim D$$

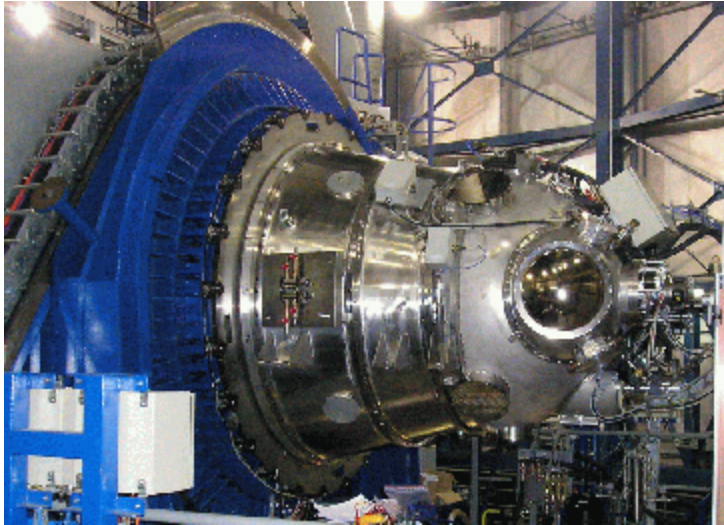
(i) and (ii) combined  $S/N \sim D^2 \rightarrow t_{\text{int}} \sim D^{-4}$

*→ huge gain: 1hr ELT = 3 months VLT*



# Instrument Sensitivity Example: HAWK-I

<http://www.eso.org/observing/etc/bin/ut4/hawki/script/hawkisimu>



## Input Flux Distribution

**Uniform (constant with wavelength)**  
NOTE: Please use the "Uniform" template spectrum instead of this option.

**Template Spectrum:** AOV (Pickles) (9480 K)  
Redshift z = 0.00

**Blackbody:** Temperature : 15000.00 K

**Single Line :** Lambda: 1250.000 nm  
Flux: 50.000  $10^{-16}$  ergs/s/cm<sup>2</sup> (per arcsec<sup>2</sup> for extended sources)  
FWHM: 1.000 nm

Target Magnitude and Mag.System:  
K = 20.00  Vega  AB  
Magnitudes are given per arcsec<sup>2</sup> for extended sources.

## Spatial Distribution:

- Point Source**
- Extended Source** diameter: 1.00 arcsec
- Extended Source (per pixel)** The Magnitude (or flux) is given per arcsec<sup>2</sup> for extended sources.

## Sky Conditions

**Airmass:** 1.20

**Seeing:** 0.80 arcsec (FWHM in V band)

## Instrument Setup

**Filter:** K

**Detector mode:** Non-destructive Read-out (NDR)

## Results

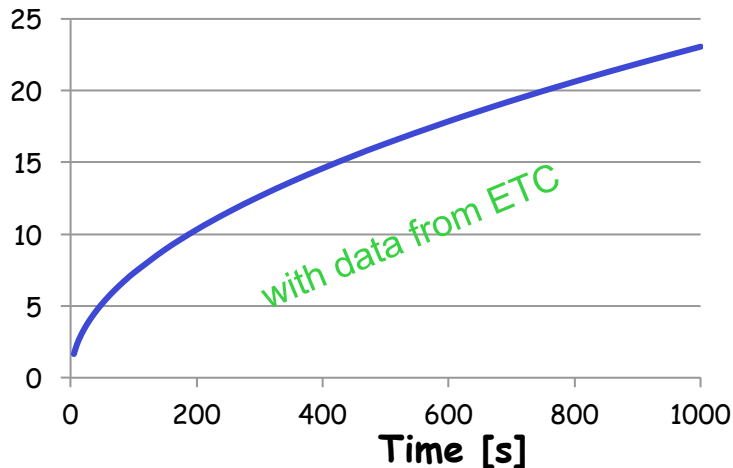
**S/N ratio:** S/N = 100.000

**Exposure Time:** NDIT = 100

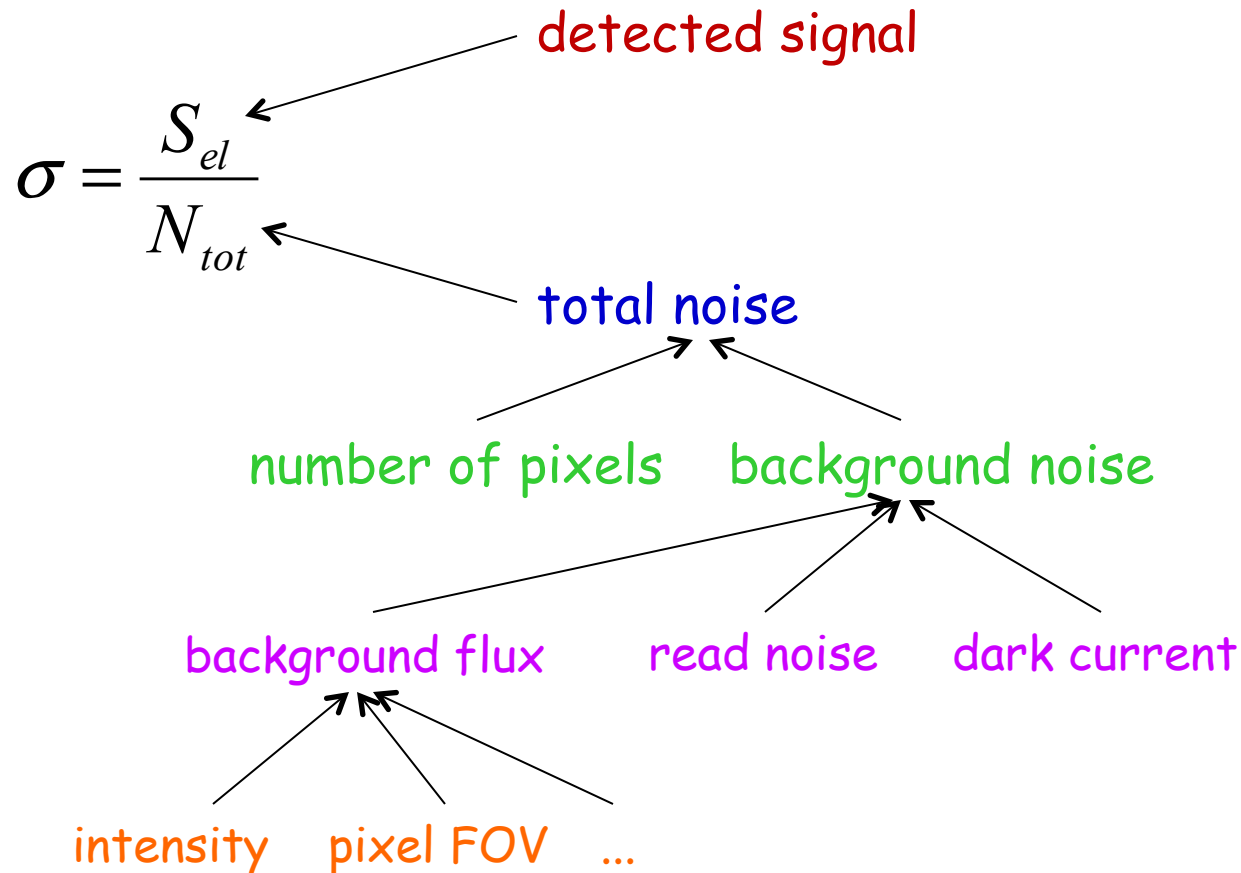
DIT = 60.000 sec

Operating temperature	75K, controlled to 1mK
Dark current [e-/s] (at 75K)	between 0.10 & 0.15
Read noise* (DCR)	~ 12 e-
Read noise* (NDR)	~ 5 e-
Linear range (1%)	60.000e- (~30.000 ADUs)
Saturation level	between 40.000 & 50.000 ADUs

S/N



# Instrument Sensitivity: Example



# Detected Signal

Detected signal  $S_{el}$  depends on:

- source flux density  $S_{src}$  [photons  $s^{-1} cm^{-2} \mu m^{-1}$ ]
- integration time  $t_{int}$  [s]
- telescope aperture  $A_{tel}$  [ $m^2$ ]
- transmission of the atmosphere  $\eta_{atm}$
- total throughput of the system  $\eta_{tot}$ , which includes:
  - reflectivity of all telescope mirrors
  - reflectivity (or transmission) of all instrument components, such as mirrors, lenses, filters, beam splitters, grating efficiencies, slit losses, etc.
- Strehl ratio  $SR$  (ratio of actual to theoretical maximum intensity)
- detector responsivity  $\eta_D G$
- spectral bandwidth  $\Delta\lambda$  [ $\mu m$ ]

$$S_{el} = S_{src} \cdot SR \cdot \Delta\lambda \cdot A_{tel} \cdot \eta_D G \cdot \eta_{atm} \cdot \eta_{tot} \cdot t_{int}$$

# Total Noise

Total noise  $N_{tot}$  depends on:

- number of pixels  $n_{pix}$  of one resolution element
- background noise per pixel  $N_{back}$

$$N_{tot} = N_{back} \sqrt{n_{pix}}$$

total background noise  $N_{back}$  depends on:

- background flux density  $S_{back}$
- integration time  $t_{int}$
- detector dark current  $I_d$
- pixel read noise ( $N$ ) and detector frames ( $n$ )

$$N_{back} = \sqrt{S_{back} \cdot t_{int} + I_d \cdot t_{int} + N_{read}^2 \cdot n}$$

# Background Flux

Background flux density  $S_{back}$  depends on:

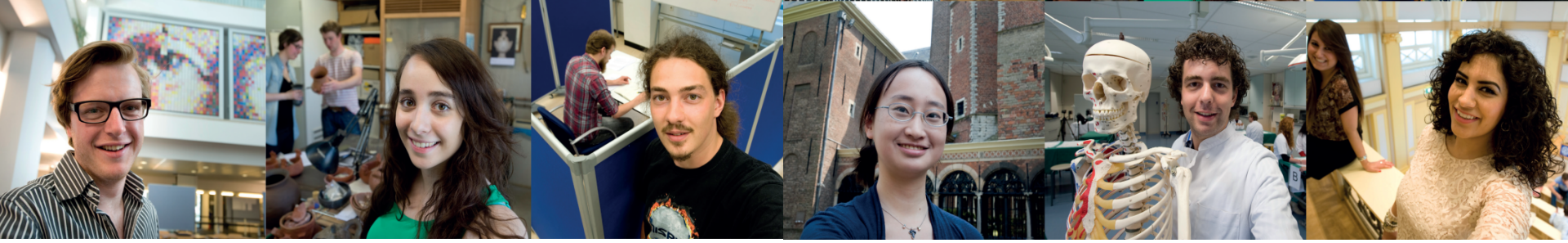
- total background intensity  $B_{tot} = (B_T + B_A) \cdot \eta_{tot}$   
 $B_T, B_A$  are thermal emissions from telescope, atmosphere (~black body)
- spectral bandwidth  $\Delta\lambda$
- pixel field of view  $A \times \Omega = 2\pi \left( 1 - \cos \left( \arctan \left( \frac{1}{2F\#} \right) \right) \right) D^2_{pix}$
- detector responsivity  $\eta_D G$
- photon energy  $hc/\lambda$

$$S_{back} = B_{tot} \cdot A \times \Omega \cdot \frac{\eta_D G \lambda}{hc} \cdot \Delta\lambda$$

# Instrument Sensitivity

Putting it all together, the minimum detectable source signal is:

$$\sigma = \frac{S_{el}}{N_{tot}} = \frac{S_{src} \cdot SR \cdot \Delta\lambda \cdot A_{tel} \cdot \eta_D G \cdot \eta_{atm} \cdot \eta_{tot} \cdot t_{int}}{N_{back} \sqrt{n_{pix}}}$$
$$\Rightarrow S_{src} = \frac{\sigma \cdot \sqrt{S_{back} \cdot t_{int} + I_d \cdot t_{int} + N_{read}^2 \cdot n \cdot \sqrt{n_{pix}}}}{SR \cdot \Delta\lambda \cdot A_{tel} \cdot \eta_D G \cdot \eta_{atm} \cdot \eta_{tot} \cdot t_{int}}$$



# Masterdag vrijdag 11 maart

Weet jij al wat je gaat doen na je bachelor?

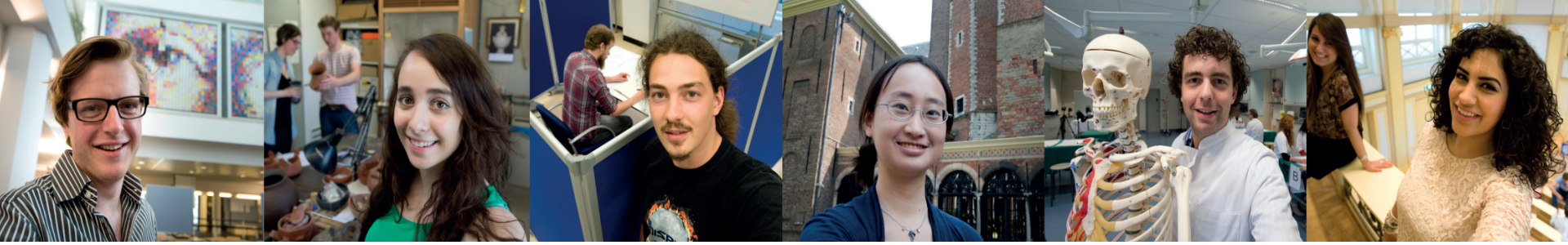
Kom langs op de Masterdag en ontdek jouw master. Meld je nu aan!

[unileidenmasters.nl](http://unileidenmasters.nl)



Universiteit  
Leiden

Bij ons leer je de wereld kennen



# Masterdag & minorenmarkt vrijdag 11 maart

Ben je aan het oriënteren op een science minor? Je minor kan van essentieel belang zijn voor jouw masterkeuze. Bezoek de minorenmarkt tijdens de Masterdag.

[unileidenmasters.nl](http://unileidenmasters.nl)



Universiteit  
Leiden

Bij ons leer je de wereld kennen