# **Astronomical Observing Techniques**

# **Lecture 4: Signal & Noise**

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# **Outline**

- 1. Introduction
- 2. Statistics
- 3. Signal-to-Noise Ratio
- 4. Instrument Sensitivities

### **Noise**

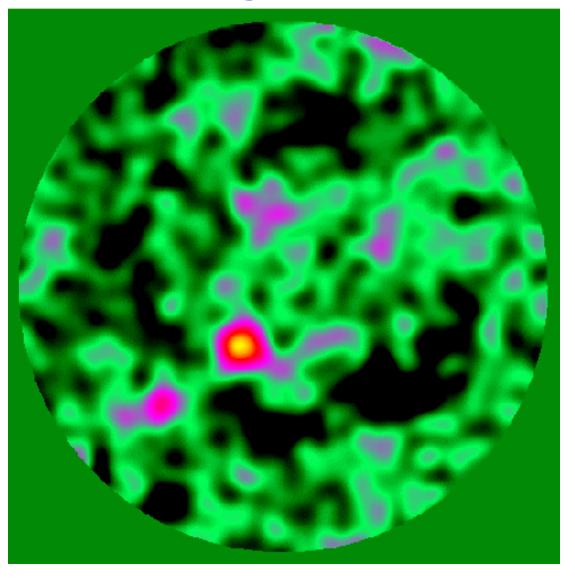
### from Wikipedia:

- Common use: unwanted sound
- Signal processing: random unwanted data without meaning
- "Signal-to-noise ratio" is sometimes used to refer to the ratio of useful to irrelevant information



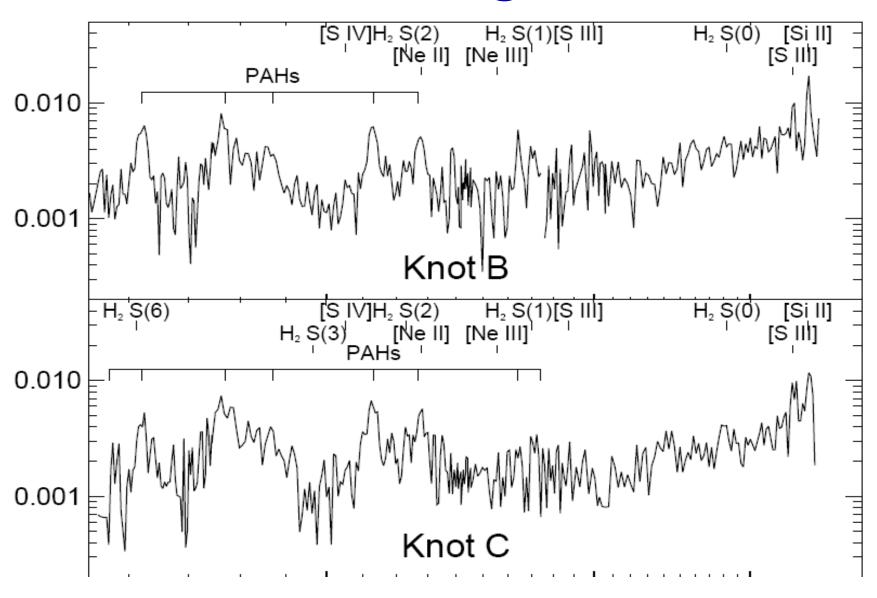
NASA researchers at Glenn Research Center conducting tests on aircraft engine noise in 1967

# Signal?



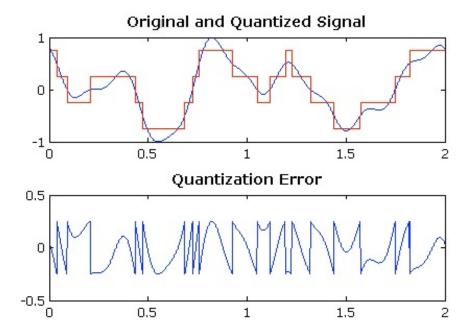
SCUBA 850µm map of the Hubble deep field

# **Noise or Signal?**



# Digitization/Quantization Noise

- Analog-to-Digital Signal Converter (ADC).
- Number of bits determines dynamic range of ADC
- Resolution: 12 bit  $2^{12} = 4096$  quantization levels 16 bit  $2^{16} = 65636$  quantization levels
- Discrete, "artificial" steps in signal levels → noise



### **Some Noise Sources in Astronomical Data**

Noise type	Signal	Background
Photon shot noise	X	X
Scintillation	X	
Cosmic rays		X
Image stability	X	
Read noise	X	X
Dark current noise	X	X
CTE (CCDs)	X	X
Flat fielding (non-linearity)	X	X
Digitization noise	X	X
Other calibration errors	X	X
Image subtraction	X	X

### **Distribution Functions**

for every t, X(t) is distributed according to cumulative distribution function

$$F(x;t) = \mathbf{P}\{X(t) \le x\}$$

- indicates probability that outcome at t will not exceed x
- probability density function (PDF) of X(t) defined by

$$f(x;t) \equiv \frac{\partial F(x;t)}{\partial x}$$

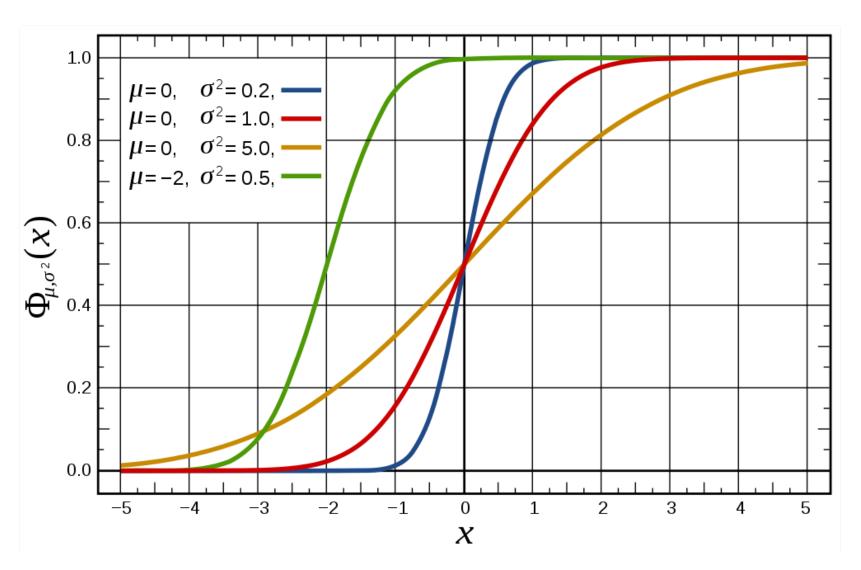
probability density function often

binomial: 
$$\binom{n}{k} p^k (1-p)^{n-k}$$

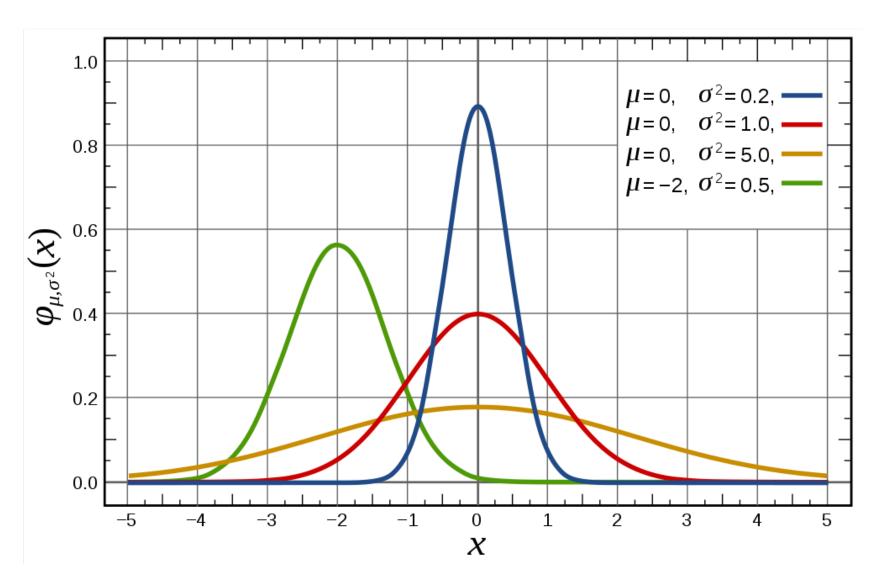
Gaussian (normal): 
$$\frac{1}{\sqrt{2\pi}}e^{-\frac{(x-\mu)^2}{2\sigma^2}}$$

Poisson: 
$$f(k;\lambda) = \frac{\lambda^k e^{-\lambda}}{k!}$$

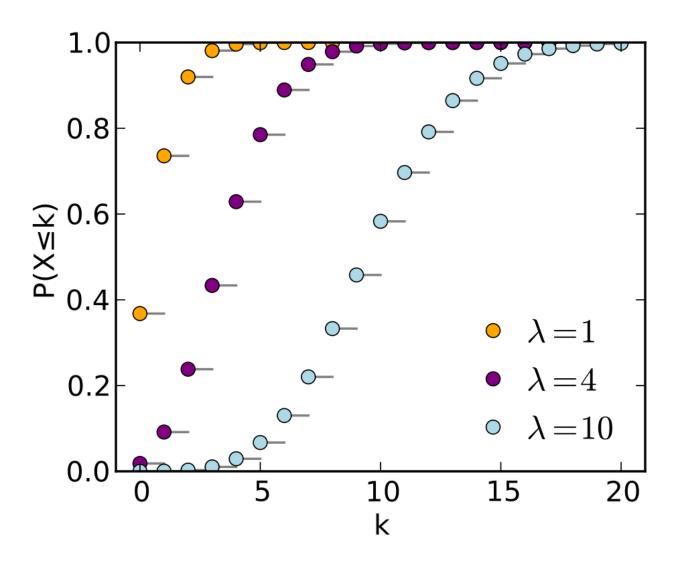
# **Normal Cumulative Distribution**



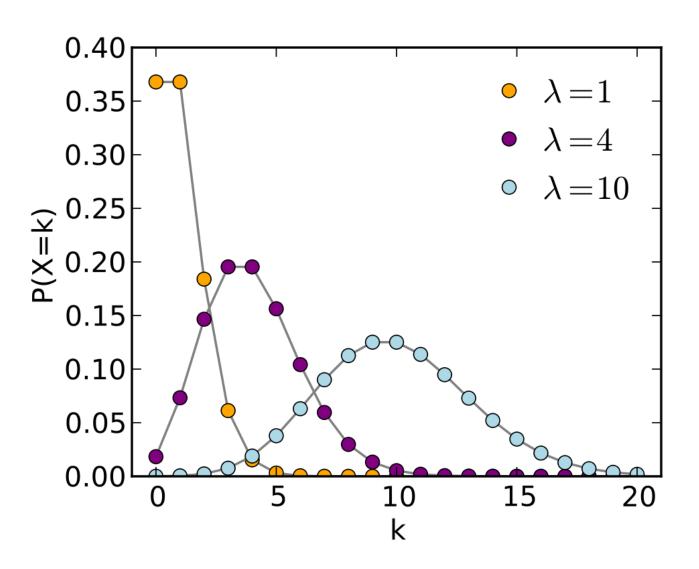
# **Normal PDF**



### **Poisson Cumulative Distribution Function**



## **Poisson PDF**



# Mean, Variance, RMS

- properties of distributions often described by a few parameters, often moments of distribution
- mean or average  $\mu(t)$  of X(t): expected value of X(t)  $\mu(t) = \mathbf{E}\{X(t)\} = \int_{-\infty}^{+\infty} x f(x;t) dx$

• variance of X(t): expected value of the square of the difference of X(t) and  $\mu(t)$ 

$$\sigma^{2}(t) = \mathbf{E}\{(X(t) - \mu(t))^{2}\} = \mathbf{E}\{X^{2}(t)\} - \mu^{2}(t)$$

 variance is square of standard deviation or root mean square (RMS)

# **Noise Propagation**

- same as error propagation
- function f(u,v,...) depends on variables u,v,...
- estimate variance of f knowing variances  $\sigma_u^2$ ,  $\sigma_v^2$ ,... of variables u,v,...  $\sigma_f^2 \equiv \lim_{N \to \infty} \sum_{i=1}^{N} (f_i \overline{f})^2$

• make assumption / approximately that average of f is well approximated by value of f for averages of variables:  $\overline{f} = f(\overline{u}, \overline{v}, ...)$ 

# **Noise Propagation (cont.)**

Taylor expansion of f around average:

$$f_i - \overline{f} \simeq (u_i - \overline{u}) \frac{\partial f}{\partial u} + (v_i - \overline{v}) \frac{\partial f}{\partial v} + \dots$$

variance in f:

$$\sigma_f^2 \simeq \lim_{N \to \infty} \sum_{i=1}^N \left[ (u_i - \overline{u}) \frac{\partial f}{\partial u} + (v_i - \overline{v}) \frac{\partial f}{\partial v} + \dots \right]^2$$

$$= \lim_{N \to \infty} \sum_{i=1}^N \left[ (u_i - \overline{u})^2 \left( \frac{\partial f}{\partial u} \right)^2 + (v_i - \overline{v})^2 \left( \frac{\partial f}{\partial v} \right)^2 + 2(u_i - \overline{u})(v_i - \overline{v}) \frac{\partial f}{\partial u} \frac{\partial f}{\partial v} + \dots \right]$$

# **Noise Propagation (cont.)**

variances of u and v

$$\sigma_u^2 \equiv \lim_{N \to \infty} \sum_{i=1}^N (u_i - \overline{u})^2; \qquad \sigma_v^2 \equiv \lim_{N \to \infty} \sum_{i=1}^N (v_i - \overline{v})^2$$

covariance of u and v

$$\sigma_{uv}^{2} \equiv \lim_{N \to \infty} \sum_{i=1}^{N} (u_{i} - \overline{u})(v_{i} - \overline{v})$$

combine Taylor expansion and these definitions

$$\sigma_f^2 = \sigma_u^2 \left(\frac{\partial f}{\partial u}\right)^2 + \sigma_v^2 \left(\frac{\partial f}{\partial v}\right)^2 + 2\sigma_{uv}^2 \frac{\partial f}{\partial u} \frac{\partial f}{\partial v} + \dots$$

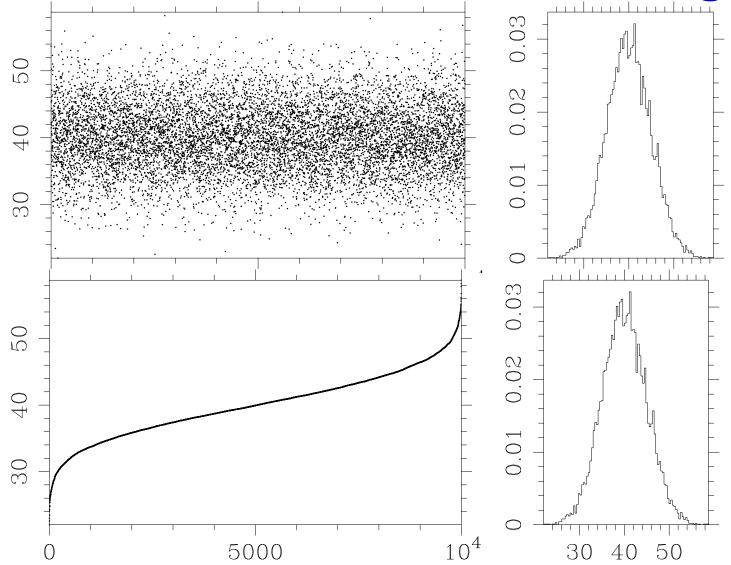
# **Noise Propagation (cont.)**

from before

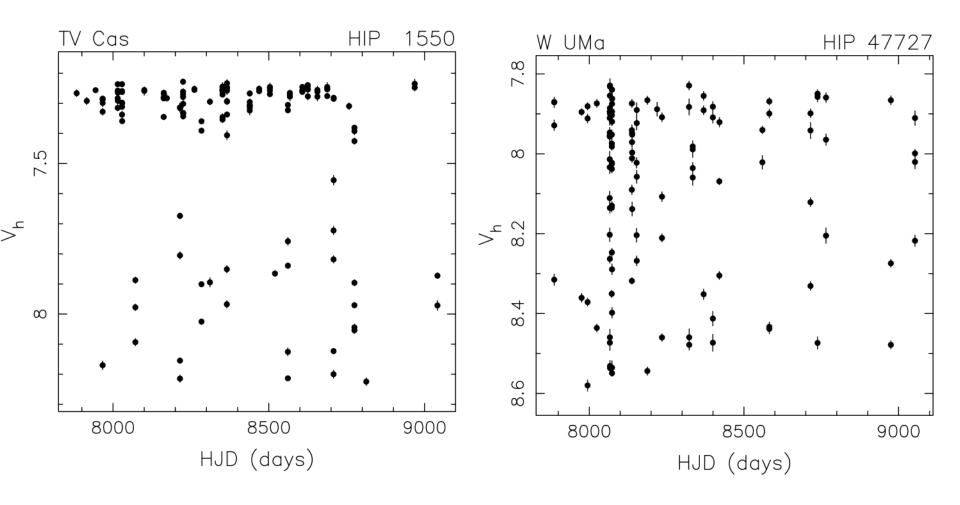
$$\sigma_f^2 = \sigma_u^2 \left(\frac{\partial f}{\partial u}\right)^2 + \sigma_v^2 \left(\frac{\partial f}{\partial v}\right)^2 + 2\sigma_{uv}^2 \frac{\partial f}{\partial u} \frac{\partial f}{\partial v} + \dots$$

- if differences  $(u_i \overline{u})$  and  $(v_i \overline{v})$  not correlated  $\rightarrow$  sign of product as often positive as negative  $\rightarrow$  covariance small compared to other terms
- if differences are correlated  $\rightarrow$  most products  $(u_i \overline{u})(v_i \overline{v})$  have the same sign  $\rightarrow$  cross-correlation term can be large

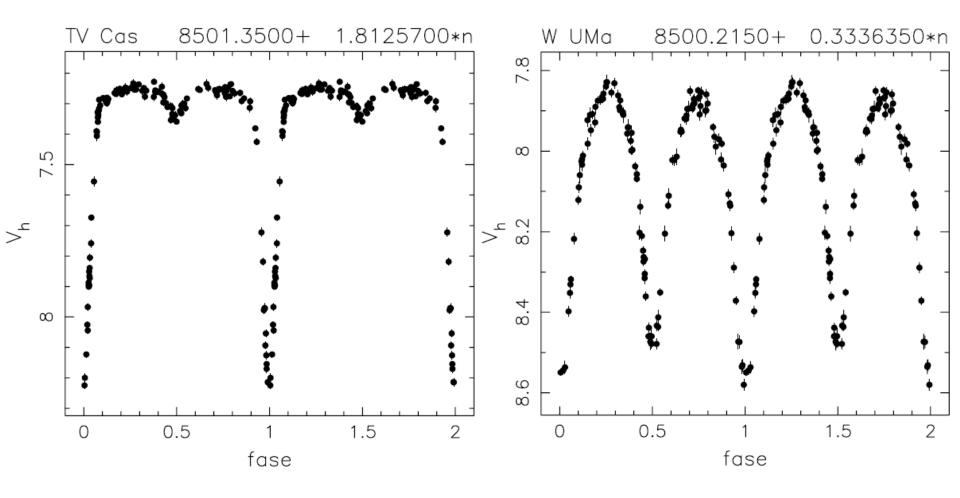
Same Distribution, Different Signals



## **Time Series of Two Stars**



# **Two Stars: Differently Sorted**



#### **Gaussian Noise**

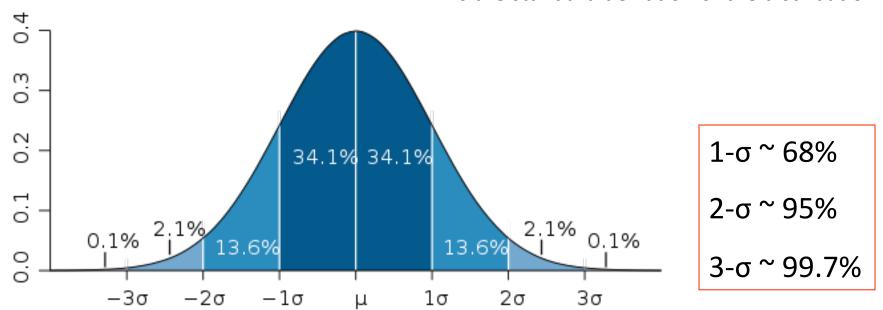
- Gaussian noise has Gaussian (normal) distribution
- Sometimes (incorrectly) called white noise (uncorrelated noise)

$$S = \frac{1}{\sqrt{2\pi}\sigma} \exp \left[ -\frac{(x-\mu)^2}{2\sigma^2} \right]$$

x is the actual value

 $\mu$  is the mean of the distribution

s is the standard deviation of the distribution



Astronomers usually consider  $S/N > 3\sigma$  as significant.

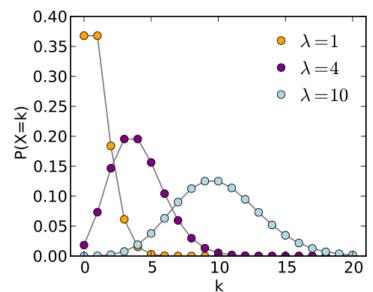
#### **Poisson Noise**

- Poisson noise has Poisson distribution
- probability of number of events occurring in constant interval of time/space if events occur with known average rate and

$$P(k,\lambda) = \frac{\lambda^k e^{-\lambda}}{k!}$$

*k* is the number of occurrences of an event (probability)

 $\lambda$  is the *expected* number of occurrences



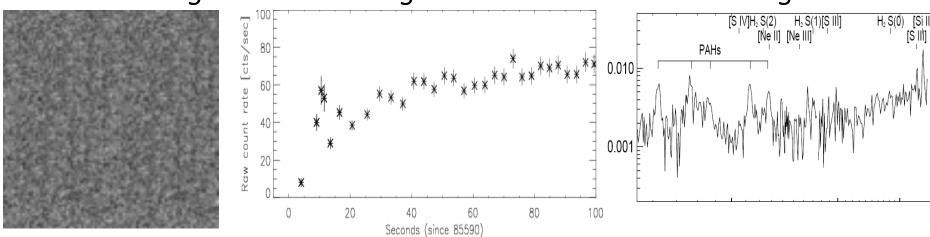
- mean of  $P(k,\lambda)$  is  $\lambda$
- standard deviation of  $P(k,\lambda)$  is square root of  $\lambda$
- example: fluctuations in photon flux in finite time intervals  $\Delta t$ . Chance to detect k photons with average flux of  $\lambda$  photons

### **Noise Measurement**

Assume purely Gaussian or Poissonian noise distribution, no other systematic noise, no correlations

Then the spatial distribution (neighbouring pixels) of the noise is equivalent to the temporal distribution (successive measurements with one pixel)

This is analogous to throwing 5 dices once versus throwing one dice



Case 1: Spatial noise (detector pixels)

Case 2: Repeated measurements Case 3: Spectrum in time (time series)

(dispersed information)

# **Poisson Noise and Integration Time**

- Integrate light from uniformly extended source on CCD
- In finite time interval  $\Delta t$ , expect average of  $\lambda$  photons
- Statistical nature of photon arrival rate  $\rightarrow$  some pixels will detect more, some less than  $\lambda$  photons.
- Noise of average signal  $\lambda$  (i.e., between pixels) is  $\sqrt{\lambda}$
- Integrate for  $2\times\Delta t \rightarrow$  expect average of  $2\times\lambda$  photons
- Noise of that signal is now  $V(2\times\lambda)$ , i.e., increased by  $V(2\times\lambda)$
- With respect to integration time t, noise will only increase
   ~Vt while signal increases ~t

## **Signal-to-Noise Ratio**

#### from Wikipedia:

Signal-to-noise ratio (often abbreviated SNR or S/N) is a measure used in science and engineering that compares the level of a desired signal to the level of background noise.

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Signal = S; Background = B; Noise = N;
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 $\leftarrow$  measured as (S+B) - mean $\{B\}$ 

 $\leftarrow$  total noise =  $\sqrt{\sum (N_i)^2}$  (if statistically independent)

Both S and N should be in units of events (photons, electrons, data numbers) per unit area (pixel, PSF size, arcsec<sup>2</sup>)

# S/N and Integration Time

Assuming the signal suffers from Poisson shot noise. Let's calculate the dependence on integration time t<sub>int</sub>:

Integrating 
$$t_{int}$$
:  $\sigma = \frac{S}{N}$ 

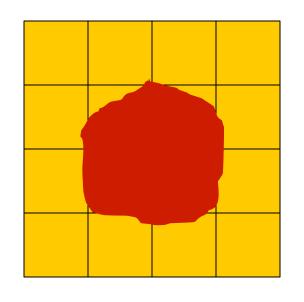
Integrating 
$$n \times t_{int}$$
:  $\sigma = \frac{n \cdot S}{\sqrt{n \cdot B}} = \sqrt{n} \frac{S}{N} \implies \frac{S}{N} \propto \sqrt{t_{int}}$ 

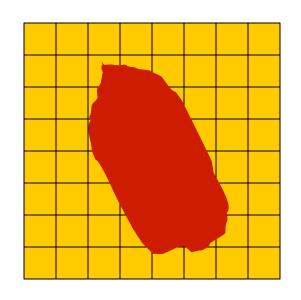
Need to integrate four times as long to get twice the S/N

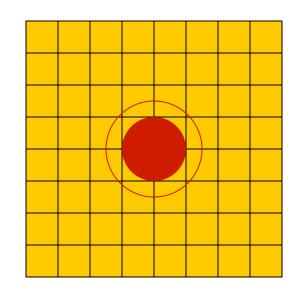
### **3 Cases**

#### Background (=noise)

#### Target







Seeing-limited point source

- pixel size ~ seeing
- PSF ≠ f(D)

Diffraction-limited, extended source

- pixel size ~ diff.lim
- PSF = f(D)
- target >> PSF

Diffraction-limited, point source

- pixel size ~ diff.lim
- PSF = f(D)
- target << PSF</li>

## Case 1: Seeing-limited "Point Source"

Signal = S; Background = B; Noise = N; Telescope diameter = D

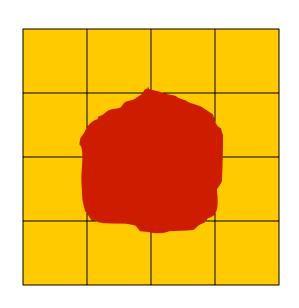
$$\theta_{\text{seeing}} \sim \text{const}$$

If detector is Nyquist-sampled to  $\theta_{\text{seeing}}$ :

$$S \sim D^2$$
 (area)

$$B \sim D^2 \rightarrow N \sim D$$
 (Poisson std.dev)

- → S/N ~ D
- → t<sub>int</sub> ~ D<sup>-2</sup>



### **Case 2: Diffraction-limited extended Source**

Signal = S; Background = B; Noise = N; Telescope diameter = D

"Diameter" of PSF ~ const

If detector Nyquist sampled to  $\theta_{diff}$ : pixel  $\sim D^{-2}$  but  $S \sim D^2$ 

 $D^2$  (telescope size) and  $D^{-2}$  (pixel FOV) cancel each other  $\rightarrow$  no change in signal

same for the background flux

- $\rightarrow$  S/N ~ const  $\rightarrow$  t<sub>int</sub> ~ const
- → no gain for larger telescopes!

Case 2B: offline re-sampling by a factor x (makes  $\theta_{diff}$  x-times larger)

since S/N 
$$\sim Vn_{pix} \rightarrow S/N \sim Vx^2 = x \rightarrow t_{int} \sim x^{-2}$$

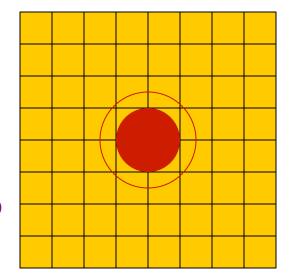
#### **Case 3: Diffraction-limited "Point Source"**

Signal = S; Background = B; Noise = N; Telescope diameter = D

"
$$S/N = (S/N)_{light bucket} \cdot (S/N)_{pixel scale}$$
"

(i) Effect of telescope aperture:

Signal 
$$S \sim D^2$$
  $\rightarrow S/N \sim D$   
Background  $B \sim D^2 \rightarrow N \sim D$ 



- (ii) Effect of pixel FOV (if Nyquist sampled to  $\theta_{diff}$ ):
  - S ~ const (pixel samples PSF = all source flux)

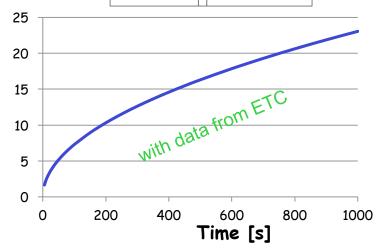
$$B \sim D^{-2} \rightarrow N \sim D^{-1} \rightarrow S/N \sim D$$

- (i) and (ii) combined  $S/N \sim D^2 \rightarrow t_{int} \sim D^{-4}$
- $\rightarrow$  huge gain: 1hr ELT = 3 months VLT

## **Instrument Sensitivity Example: HAWK-I**



Operating temperature	75K, controlled to 1mK
Dark current [e-/s] (at 75K)	between 0.10 & 0.15
Read noise* (DCR)	~ 12 e-
Read noise* (NDR)	~ 5 e-
Linear range (1%)	60.000e- (~30.000 ADUs)
Saturation level	between 40.000 & 50.000 ADUs



S/N

http://www.eso.org/observing/etc/bin/ut4/hawki/script/hawkisimu

#### Input Flux Distribution

NOTE: Please use the "Unifor	wavelength) m" template spectrum instead of this option.	Target Magnitude and Mag.System  © Vega
Template Spectrum:	A0V (Pickles) (9480 K) ▼	K ▼ = 20.00
Tell Coll	Redshift z = 0.00	Magnitudes are given per arcsec <sup>2</sup>
Blackbody:	Temperature: 15000.00 K	for extended sources.
V III	Lambda: 1250.000 nm	
O Single Line :	Flux: 50.000 10 <sup>-16</sup> ergs/s/cm <sup>2</sup> (per	arcsec <sup>2</sup> for extended sources)
	FWHM: 1.000 nm	

The Magnitude (or flux) is given per arcsec<sup>2</sup> for extended sources.

#### **Spatial Distribution:**

- Point Source
- Extended Source diameter: 1.00 arcsec
- Extended Source (per pixel)

#### **Sky Conditions**

Airmass: 1.20

Seeing: 0.80 arcsec (FWHM in V band)

#### **Instrument Setup**

Filter:

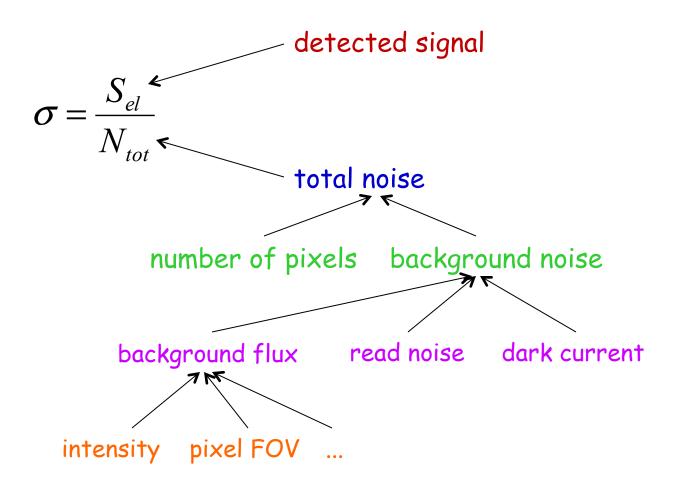
K ▼

**Detector mode:** Non-destructive Read-out (NDR)

#### Results

<u>S/N ratio:</u> <u>S/N</u> = 100.000 <u>DIT</u> = 60.000 sec <u>DIT</u> = 60.000 sec

### **Instrument Sensitivity: Example**



### **Detected Signal**

#### Detected signal S<sub>el</sub> depends on:

- source flux density S<sub>src</sub> [photons s<sup>-1</sup> cm<sup>-2</sup> μm<sup>-1</sup>]
- integration time t<sub>int</sub> [s]
- telescope aperture A<sub>tel</sub> [m<sup>2</sup>]
- transmission of the atmosphere  $\eta_{atm}$
- total throughput of the system  $\eta_{tot}$ , which includes:
  - reflectivity of all telescope mirrors
  - reflectivity (or transmission) of all instrument components, such as mirrors, lenses, filters, beam splitters, grating efficiencies, slit losses, etc.
- Strehl ratio SR (ratio of actual to theoretical maximum intensity)
- detector responsivity η<sub>D</sub>G
- spectral bandwidth Δλ [μm]

$$S_{el} = S_{src} \cdot SR \cdot \Delta \lambda \cdot A_{tel} \cdot \eta_D G \cdot \eta_{atm} \cdot \eta_{tot} \cdot t_{int}$$

### **Total Noise**

#### Total noise N<sub>tot</sub> depends on:

- number of pixels n<sub>pix</sub> of one resolution element
- background noise per pixel N<sub>back</sub>

$$N_{tot} = N_{back} \sqrt{n_{pix}}$$

total background noise N<sub>back</sub> depends on:

- background flux density S<sub>back</sub>
- integration time t<sub>int</sub>
- detector dark current I<sub>d</sub>
- pixel read noise (N) and detector frames (n)

$$N_{back} = \sqrt{S_{back} \cdot t_{int} + I_d \cdot t_{int} + N_{read}^2 \cdot n}$$

### **Background Flux**

#### Background flux density S<sub>back</sub> depends on:

- the total background intensity  $B_{tot} = (B_T + B_A) \cdot \eta_{tot}$  where  $B_T$  and  $B_A$  are the thermal emissions from telescope and atmosphere, approximated by  $B_{T,A} = \frac{2hc^2}{\lambda^5} \left[ \frac{\varepsilon}{\exp\left[\frac{hc}{kT\lambda}\right] 1} \right]$
- the spectral bandwidth  $\Delta \lambda$
- the pixel field of view  $A \times \Omega = 2\pi \left(1 \cos\left(\arctan\left(\frac{1}{2F^{\#}}\right)\right)\right) D^{2}_{pix}$
- the detector responsivity  $\eta_D G$ , and
- the photon energy hc/l

$$S_{back} = B_{tot} \cdot A \times \Omega \cdot \frac{\eta_D G \lambda}{hc} \cdot \Delta \lambda$$

### **Instrument Sensitivity**

Putting it all together, the minimum detectable source signal is:

$$\sigma = \frac{S_{el}}{N_{tot}} = \frac{S_{src} \cdot SR \cdot \Delta\lambda \cdot A_{tel} \cdot \eta_D G \cdot \eta_{atm} \cdot \eta_{tot} \cdot t_{int}}{N_{back} \sqrt{n_{pix}}}$$

$$\Rightarrow S_{src} = \frac{\sigma \cdot \sqrt{S_{back} \cdot t_{int} + I_d \cdot t_{int} + N_{read}^2 \cdot n \cdot \sqrt{n_{pix}}}}{SR \cdot \Delta\lambda \cdot A_{tel} \cdot \eta_D G \cdot \eta_{atm} \cdot \eta_{tot} \cdot t_{int}}$$