

Detection of Light 2020

Homework 3

Detector Noise and Extrinsic Photoconductors

17 marks total

28 February 2020

Due: Before the start of class Fri 06 Mar 2020
Questions, please e-mail to: dorval@strw.leidenuniv.nl

1. Detector Noise [4]

- (a) [2 points] The Noise-Equivalent Power (NEP) is defined as “The signal power yielding an RMS signal-to-noise ratio of unity for a frequency bandwidth of 1Hz”.
- i) From this, derive an expression for NEP as a function of **only** the RMS noise current $\sqrt{\langle I_N^2 \rangle}$ (in Amps), sensitivity S and frequency bandwidth Δf .
- ii) Hence show that for a system where 1/f noise sources may be neglected, the total system NEP is given by

$$NEP = \frac{2hc}{\lambda} \sqrt{\frac{\phi}{\eta} + \frac{kT}{(q\eta G)^2 R}}$$

- (b) [2 points] Consider now our intrinsic detector from Ex.2 part 2, operated under the same conditions (i.e. $\phi = 10^5 \gamma/s$, $t_{exp} = 0.5s$, take $\eta = 0.44$, $G = 0.68$, $R = 7.9 \times 10^8 \Omega$). Calculate:
- i) The critical temperature beyond which the detector must be cooled, in order to obtain photon-limited performance, and
- ii) The NEP in this limit and the maximum signal-to-noise ratio attainable with this detector for the source and exposure time in question. Take $\lambda = 0.7$ microns

2. Thermal Excitation of Extrinsic Photoconductors [3]

By analogy with intrinsic semiconductors, it can be shown that the concentration of negative carriers in a doped n-type semiconductor is given by:

$$n = \left(\frac{N_D - N_A}{N_A + n} \right) \left(\frac{2}{\delta} \right) \left(\frac{2\pi m_n^* kT}{h^2} \right)^{3/2} e^{-E_i/kT}$$

Impurity	Ge		Si	
	Cutoff wavelength	Photoionization cross section	Cutoff wavelength	Photoionization cross section
	λ_c (μm)	σ_i (cm^2)	λ_c (μm)	σ_i (cm^2)
Al			18.5	8×10^{-16}
B	119	1.0×10^{-14}	28	1.4×10^{-15}
Be	52		8.3	5×10^{-18}
Ga	115	1.0×10^{-14}	17.2	5×10^{-16}
In	111		7.9	3.3×10^{-17}
As	98	1.1×10^{-14}	23	2.2×10^{-15}
Cu	31	1.0×10^{-15}	5.2	5×10^{-18}
P	103	1.5×10^{-14}	27	1.7×10^{-15}
Sb	129	1.6×10^{-14}	29	6.2×10^{-15}

Figure 1: Properties of some extrinsic photoconductors.

where N_D and N_A are the concentrations of donors and acceptor respectively, δ is the ground state degeneracy of the donor impurity (very close to unity) and E_i is in this case not the gap energy but the ionization energy for the n-type impurity. Similarly, for p-type materials:

$$p = \left(\frac{N_A - N_D}{N_D + p} \right) \left(\frac{2}{\delta} \right) \left(\frac{2\pi m_p^* kT}{h^2} \right)^{3/2} e^{-E_i/kT}$$

where it has been assumed that $E_i \gg kT$.

- [2 points] Consider a Si:B material. Is this a n-type, or p-type material? Use the value of the cutoff wavelength λ_c from Fig. 1 to calculate the ionization energy of the impurity atoms (in meV). Is the stated assumption valid at room temperature ($T = 300\text{K}$)? Is it valid at 77K , the temperature of liquid nitrogen?
- [1 points] Use Fig. 1 to calculate the linear absorption coefficient of a silicon crystal doped with $1.3 \times 10^{13} \text{cm}^{-3}$ boron atoms (Assume that the photoionization cross-section does not significantly change with energy). How does this compare to the absorption coefficient of intrinsic Si from Ex.2? What would be a natural solution to improve the quantum efficiency of the extrinsic photoconductor?

3. BIB Detectors [10]

Consider the Blocked Impurity Band (BIB) photoconductor shown in Fig. 2, consisting of a highly-doped Si:As (n-type) infrared-active layer with a small but non-negligible fraction of p-type impurities, which is overlaid with a high-purity Si blocking layer to which a positive bias voltage V_b is applied. A depletion region of width w is formed in which photons may be absorbed and subsequently detected with high efficiency.

For this question take the blocking and IR-active layer widths to be $t_B = 5.0 \mu\text{m}$ and $t_A = 15.0 \mu\text{m}$ respectively, and the densities of n-type dopants and p-type impurities are $N_D = 5 \times 10^{17} \text{cm}^{-3}$, $N_A = 10^{13} \text{cm}^{-3}$. The dielectric constant $\kappa_0 = 11.8$ for Si at room temperature, and the permittivity of free space $\epsilon_0 = 8.854 \times 10^{-12} \text{F m}^{-1}$.

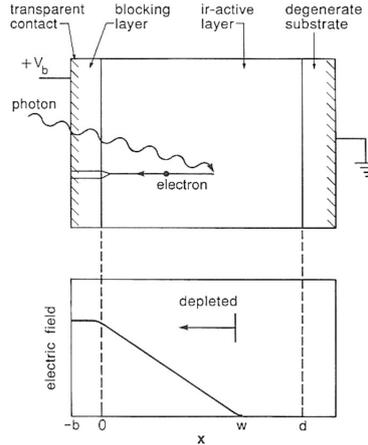


Figure 2: Diagram of a BIB depletion region.

- (a) [3 points] Sketch an energy-level diagram for this BIB detector as a function of position x along the detector. With reference to this explain:
- What is meant by “hopping” in over-doped extrinsic semiconductors, and how does the BIB band structure help overcome this problem?
 - What is meant by the depletion region of the detector, and why does it have a finite extent within the IR-active layer?
- (b) [4 points] i) By solving the one-dimensional Poisson equation

$$\frac{dE}{dx} = \frac{\rho}{\epsilon_0 \kappa_0},$$

obtain an expression for the electric field E as a function of distance x from the inner edge of the blocking layer, depletion layer width w and acceptor density N_A .

ii) Hence show that the width of the depletion region as a function of the bias voltage V_b and thickness of the blocking layer t_B is given by:

$$w = \left(\frac{2\kappa_0\epsilon_0}{qN_A} |V_b| + t_B^2 \right)^{\frac{1}{2}} - t_B$$

[Hint: Use the boundary condition $V(x = -t_B) = |V_b|$ at the outer edge of the blocking layer.]

- (c) [1 points] What is the value of the critical bias voltage in this detector for which the thickness of the depletion region is the same as the thickness of the IR-active layer, t_A ?
- (d) [2 points] Consider photo-electrons accelerated near the boundary of the blocking layer, where the mean-free path length is $\langle l \rangle \sim 0.3 \mu\text{m}$.
- On average how much energy will these electrons gain in between collisions, when the detector is operated under the critical bias voltage?
 - Given that the ionisation energy of arsenic in Si:As is $E_i = 54 \text{ meV}$, explain briefly how gains of $G > 1$ may be obtained in BIB detectors. In what part of the detector does this occur, and why?